# On the Hunter-Saxton equation on $\text{Diff}_{\mathcal{A}}(\mathbb{R})$

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# Groups related to $\text{Diff}_c(\mathbb{R})$

The reflexive nuclear (LF) space  $C_c^{\infty}(\mathbb{R})$  of smooth functions with compact support leads to the well-known regular Lie group  $\text{Diff}_{c}(\mathbb{R})$ . We will now define an extension of this group which will play a major role in the later parts of this article. Define  $C_{c,2}^{\infty}(\mathbb{R}) = \{f : f' \in C_c^{\infty}(\mathbb{R})\}$  to be the space of antiderivatives of smooth functions with compact support. It is a reflexive nuclear (LF) space. We also define the space  $C^\infty_{c,1}(\mathbb{R}) = \left\{ f \in C^\infty_{c,2}(\mathbb{R}) \ : \ f(-\infty) = 0 
ight\}$  of antiderivatives of the form  $x \mapsto \int_{-\infty}^{x} g \, dy$  with  $g \in C_{c}^{\infty}(\mathbb{R})$ .  $\operatorname{Diff}_{c,2}(\mathbb{R}) = \{ \varphi = \operatorname{Id} + f : f \in C^{\infty}_{c,2}(\mathbb{R}), f' > -1 \}$  is the corresponding group. Define the two functionals  $\text{Shift}_{\ell}$ ,  $\text{Shift}_{r}$ :  $\text{Diff}_{c,2}(\mathbb{R}) \to \mathbb{R}$  by

Shift<sub>$$\ell$$</sub>( $\varphi$ ) = ev <sub>$-\infty$</sub> ( $f$ ) =  $\lim_{x \to -\infty} f(x)$ , Shift <sub>$r$</sub> ( $\varphi$ ) = ev <sub>$\infty$</sub> ( $f$ ) =  $\lim_{x \to \infty} f(x)$   
for  $\varphi(x) = x + f(x)$ .

Then the short exact sequence of smooth homomorphisms of Lie groups

$$\operatorname{Diff}_{c}(\mathbb{R}) \longrightarrow \operatorname{Diff}_{c,2}(\mathbb{R}) \xrightarrow{(\operatorname{Shift}_{\ell},\operatorname{Shift}_{r})} \gg (\mathbb{R}^{2},+)$$

describes a semidirect product, where a smooth homomorphic section  $s : \mathbb{R}^2 \to \text{Diff}_{c,2}(\mathbb{R})$  is given by the composition of flows  $s(a,b) = \text{Fl}_a^{X_\ell} \circ \text{Fl}_b^{X_r}$  for the vectorfields  $X_\ell = f_\ell \partial_x$ ,  $X_r = f_r \partial_x$  with  $[X_\ell, X_r] = 0$  where  $f_\ell, f_r \in C^\infty(\mathbb{R}, [0, 1])$  satisfy

$$f_{\ell}(x) = \begin{cases} 1 & \text{for } x \leq -1 \\ 0 & \text{for } x \geq 0, \end{cases} \qquad f_r(x) = \begin{cases} 0 & \text{for } x \leq 0 \\ 1 & \text{for } x \geq 1. \end{cases}$$
(1)

The normal subgroup  $\text{Diff}_{c,1}(\mathbb{R}) = \text{ker}(\text{Shift}_{\ell}) = \{\varphi = \text{Id} + f : f \in C^{\infty}_{c,1}(\mathbb{R}), f' > -1\}$  of diffeomorphisms which have no shift at  $-\infty$  will play an important role later on.

# Some diffeomorphism groups on $\ensuremath{\mathbb{R}}$

We have the following smooth injective group homomorphisms:



Each group is a normal subgroup in any other in which it is contained, in particular in  $\text{Diff}_{\mathcal{B}}(\mathbb{R})$ .

For S this works the same as for  $C_c^{\infty}$ . For  $H^{\infty}$  it is slightly mote subtle.

### The setting

We will denote by  $\mathcal{A}(\mathbb{R})$  any of the spaces  $C_c^{\infty}(\mathbb{R})$ ,  $\mathcal{S}(\mathbb{R})$  or  $H^{\infty}(\mathbb{R})$  and by  $\text{Diff}_{\mathcal{A}}(\mathbb{R})$  the corresponding groups  $\text{Diff}_c(\mathbb{R})$ ,  $\text{Diff}_{\mathcal{S}}(\mathbb{R})$  or  $\text{Diff}_{H^{\infty}}(\mathbb{R})$  as defined in Sections ??, ?? and ??.

Similarly  $\mathcal{A}_1(\mathbb{R})$  will denote any of the spaces  $C_{c,1}^{\infty}(\mathbb{R})$ ,  $\mathcal{S}_1(\mathbb{R})$  or  $H_1^{\infty}(\mathbb{R})$  and  $\text{Diff}_{\mathcal{A}|1}(\mathbb{R})$  the corresponding groups  $\text{Diff}_{c,1}(\mathbb{R})$ ,  $\text{Diff}_{\mathcal{S}_1}(\mathbb{R})$  or  $\text{Diff}_{\mathcal{H}_1^{\infty}}(\mathbb{R})$  as defined in Sections ??, ?? and ??.

The  $\dot{H}^1$ -metric. For  $\text{Diff}_{\mathcal{A}}(\mathbb{R})$  and  $\text{Diff}_{\mathcal{A}_1}(\mathbb{R})$  the homogeneous  $H^1$ -metric is given by

$$G_{\varphi}(X \circ \varphi, Y \circ \varphi) = G_{\mathsf{Id}}(X, Y) = \int_{\mathbb{R}} X'(x) Y'(x) \ dx \ ,$$

where X, Y are elements of the Lie algebra  $\mathcal{A}(\mathbb{R})$  or  $\mathcal{A}_1(\mathbb{R})$ . We shall also use the notation

$$\langle \cdot, \cdot \rangle_{\dot{H}^1} := G(\cdot, \cdot)$$

#### Theorem

On  $\text{Diff}_{\mathcal{A} 1}(\mathbb{R})$  the geodesic equation is the Hunter-Saxton equation

$$(\varphi_t)\circ \varphi^{-1}=u \qquad u_t=-uu_x+\frac{1}{2}\int_{-\infty}^x (u_x(z))^2 dz ,$$

and the induced geodesic distance is positive.

On the other hand the geodesic equation does not exist on the subgroups  $\text{Diff}_{\mathcal{A}}(\mathbb{R})$ , since the adjoint  $\operatorname{ad}(X)^*\check{G}_{\text{Id}}(X)$  does not lie in  $\check{G}_{\text{Id}}(\mathcal{A}(\mathbb{R}))$  for all  $X \in \mathcal{A}(\mathbb{R})$ .

One obtains the classical form of the Hunter-Saxton equation by differentiating:

$$u_{tx}=-uu_{xx}-\frac{1}{2}u_x^2,$$

Note that  $\text{Diff}_{\mathcal{A}}(\mathbb{R})$  is a natural example of a non-robust Riemannian manifold.

#### Proof

Note that  $\check{G}_{\mathsf{Id}} : \mathcal{A}_1(\mathbb{R}) \to \mathcal{A}_1(\mathbb{R})^*$  is given by  $\check{G}_{\mathsf{Id}}(X) = -X''$  if we use the  $L^2$ -pairing  $X \mapsto (Y \mapsto \int XYdx)$  to embed functions into the space of distributions. We now compute the adjoint of  $\mathsf{ad}(X)$ :

$$ig\langle \operatorname{ad}(X)^*\check{G}_{\mathsf{ld}}(Y),Z ig
angle = \check{G}_{\mathsf{ld}}(Y,\operatorname{ad}(X)Z) = G_{\mathsf{ld}}(Y,-[X,Z])$$
  
=  $\int_{\mathbb{R}} Y'(x) (X'(x)Z(x) - X(x)Z'(x))' dx$   
=  $\int_{\mathbb{R}} Z(x) (X''(x)Y'(x) - (X(x)Y'(x))'') dx$ .

Therefore the adjoint as an element of  $\mathcal{A}_1^*$  is given by

$$\operatorname{\mathsf{ad}}(X)^*\check{G}_{\operatorname{\mathsf{Id}}}(Y)=X''Y'-(XY')''$$
 .

For X = Y we can rewrite this as

$$\mathsf{ad}(X)^* \check{G}_{\mathsf{Id}}(X) = \frac{1}{2} ((X'^2)' - (X^2)''') = \frac{1}{2} \Big( \int_{-\infty}^x X'(y)^2 \, dy - (X^2)' \Big)''$$
  
=  $\frac{1}{2} \check{G}_{\mathsf{Id}} \Big( - \int_{-\infty}^x X'(y)^2 \, dy + (X^2)' \Big) .$ 

If  $X \in \mathcal{A}_1(\mathbb{R})$  then the function  $-\frac{1}{2} \int_{-\infty}^{x} X'(y)^2 dy + \frac{1}{2}(X^2)'$  is again an element of  $\mathcal{A}_1(\mathbb{R})$ . This follows immediately from the definition of  $\mathcal{A}_1(\mathbb{R})$ . Therefore the geodesic equation exists on Diff $_{\mathcal{A}}_1(\mathbb{R})$  and is as given.

However if  $X \in \mathcal{A}(\mathbb{R})$ , a neccessary condition for  $\int_{-\infty}^{x} (X'(y))^2 dy \in \mathcal{A}(\mathbb{R})$  would be  $\int_{-\infty}^{\infty} X'(y)^2 dy = 0$ , which would imply X' = 0. Thus the geodesic equation does not exist on  $\mathcal{A}(\mathbb{R})$ .

The positivity of geodesic distance will follow from the explicit formula for geodesic distance below. QED.

#### Theorem

We define the R-map by:

$$R: \left\{ egin{array}{l} \mathsf{Diff}_{\mathcal{A}_1}(\mathbb{R}) o \mathcal{A}ig(\mathbb{R},\mathbb{R}_{>-2}ig) \subset \mathcal{A}(\mathbb{R},\mathbb{R}) \ arphi \mapsto 2 \left((arphi')^{1/2}-1
ight) \,. \end{array} 
ight.$$

The R-map is invertible with inverse

$$R^{-1}: \begin{cases} \mathcal{A}(\mathbb{R}, \mathbb{R}_{>-2}) \to \mathsf{Diff}_{\mathcal{A}_1}(\mathbb{R}) \\ \gamma \mapsto x + \frac{1}{4} \int_{-\infty}^x \gamma^2 + 4\gamma \ dx \ . \end{cases}$$

The pull-back of the flat  $L^2$ -metric via R is the  $\dot{H}^1$ -metric on  $\text{Diff}_{\mathcal{A}}(\mathbb{R})$ , *i.e.*,

$$R^*\langle \cdot, \cdot \rangle_{L^2} = \langle \cdot, \cdot \rangle_{\dot{H}^1}$$
.

Thus the space  $(\text{Diff}_{\mathcal{A}_1}(\mathbb{R}), \dot{H}^1)$  is a flat space in the sense of Riemannian geometry.

Here  $\langle \cdot, \cdot \rangle_{L^2}$  denotes the  $L^2$ -inner product on  $\mathcal{A}(\mathbb{R})$  with constant volume dx.

### Proof

To compute the pullback of the  $L^2$ -metric via the *R*-map we first need to calculate its tangent mapping. For this let  $h = X \circ \varphi \in T_{\varphi} \text{Diff}_{\mathcal{A}_1}(\mathbb{R})$  and let  $t \mapsto \psi(t)$  be a smooth curve in  $\text{Diff}_{\mathcal{A}_1}(\mathbb{R})$  with  $\psi(0) = \text{Id}$  and  $\partial_t|_0\psi(t) = X$ . We have:

$$\begin{split} T_{\varphi}R.h &= \partial_t|_0 R(\psi(t) \circ \varphi) = \partial_t|_0 2\Big(((\psi(t) \circ \varphi)_x)^{1/2} - 1\Big) \\ &= \partial_t|_0 2((\psi(t)_x \circ \varphi) \varphi_x)^{1/2} \\ &= 2(\varphi_x)^{1/2} \partial_t|_0 ((\psi(t)_x)^{1/2} \circ \varphi) = (\varphi_x)^{1/2} (\frac{\psi_{tx}(0)}{(\psi(0)_x)^{-1/2}} \circ \varphi) \\ &= (\varphi_x)^{1/2} (X' \circ \varphi) = (\varphi')^{1/2} (X' \circ \varphi) \,. \end{split}$$

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Using this formula we have for  $h = X_1 \circ \varphi, k = X_2 \circ \varphi$ :

$$R^*\langle h,k\rangle_{L^2} = \langle T_{\varphi}R.h, T_{\varphi}R.k\rangle_{L^2} = \int_{\mathbb{R}} X_1'(x)X_2'(x)\,dx = \langle h,k\rangle_{\dot{H}^1} \,\,QED$$

# Corollary

Given  $\varphi_0, \varphi_1 \in \text{Diff}_{\mathcal{A}_1}(\mathbb{R})$  the geodesic  $\varphi(t, x)$  connecting them is given by

$$\varphi(t,x) = R^{-1} \Big( (1-t)R(\varphi_0) + tR(\varphi_1) \Big)(x)$$

and their geodesic distance is

$$d(\varphi_0,\varphi_1)^2 = 4 \int_{\mathbb{R}} \left( (\varphi_1')^{1/2} - (\varphi_0')^{1/2} \right)^2 \, dx \; .$$

The methods of Boris and my lecture give local well-posedness of the geodesic equation in  $\operatorname{Diff}_{H_1^{\infty}}(\mathbb{R})$  only. But this construction shows much more: For  $\mathcal{S}_1$ ,  $C_1^{\infty}$ , and even for many kinds of Denjoy-Carleman ultradifferentiable model spaces (not explained here). This shows that Sobolev space methods for treating nonlinear PDEs is not the only method.

**Corollary:** The metric space  $(\text{Diff}_{\mathcal{A}_1}(\mathbb{R}), \dot{H}^1)$  is path-connected and geodesically convex but not geodesically complete. In particular, for every  $\varphi_0 \in \text{Diff}_{\mathcal{A}_1}(\mathbb{R})$  and  $h \in T_{\varphi_0} \text{Diff}_{\mathcal{A}_1}(\mathbb{R})$ ,  $h \neq 0$ there exists a time  $T \in \mathbb{R}$  such that  $\varphi(t, \cdot)$  is a geodesic for |t| < |T| starting at  $\varphi_0$  with  $\varphi_t(0) = h$ , but  $\varphi_x(T, x) = 0$  for some  $x \in \mathbb{R}$ .

**Theorem:** The square root representation on the diffeomorphism group  $\text{Diff}_{\mathcal{A}}(\mathbb{R})$  is a bijective mapping, given by:

$$R: \begin{cases} \mathsf{Diff}_{\mathcal{A}}(\mathbb{R}) \to \big(\mathsf{Im}(R), \|\cdot\|_{L^2}\big) \subset \big(\mathcal{A}\big(\mathbb{R}, \mathbb{R}_{>-2}\big), \|\cdot\|_{L^2}\big) \\ \varphi \mapsto 2\left((\varphi')^{1/2} - 1\right). \end{cases}$$

The pull-back of the restriction of the flat  $L^2$ -metric to Im(R) via R is again the homogeneous Sobolev metric of order one. The image of the R-map is the splitting submanifold of  $\mathcal{A}(\mathbb{R}, \mathbb{R}_{>-2})$  given by:

$$\operatorname{Im}(R) = \left\{ \gamma \in \mathcal{A}(\mathbb{R}, \mathbb{R}_{>-2}) : F(\gamma) := \int_{\mathbb{R}} \gamma(\gamma + 4) \, dx = 0 \right\} \, .$$

On the space  $\text{Diff}_{\mathcal{A}}(\mathbb{R})$  the geodesic equation does not exist. Still: **Corollary:** The geodesic distance  $d^{\mathcal{A}}$  on  $\text{Diff}_{\mathcal{A}}(\mathbb{R})$  coincides with the restriction of  $d^{\mathcal{A}_1}$  to  $\text{Diff}_{\mathcal{A}}(\mathbb{R})$ , i.e., for  $\varphi_0, \varphi_1 \in \text{Diff}_{\mathcal{A}}(\mathbb{R})$  we have

$$d^{\mathcal{A}}(\varphi_0,\varphi_1)=d^{\mathcal{A}_1}(\varphi_0,\varphi_1)$$
 .