

Bernoulli Shifts

Bernoulli trials in probability refer to successive coinflips or rolls of the die. They are modeled as a shift-space over an alphabet $\mathcal{A} = \{1, \dots, N\}$:

$$\Sigma = \mathcal{A}^{\mathbb{Z}} \text{ (two-sided) or } \Sigma = \mathcal{A}^{\mathbb{N}} \text{ (one-sided)}$$

with a probability vector

$$p = (p_1, \dots, p_N) \quad p_i \in [0, 1], \quad \sum_{i=1}^N p_i = 1.$$

The corresponding **Bernoulli measure** μ_p gives **cylinder sets**

$$Z_{[k+1, k+n]}(a) = \{x \in \Sigma : x_{k+1} \dots x_{k+n} = a_1 \dots a_n\}$$

the mass

$$\mu_p(Z_{[k+1, k+n]}(a)) = \prod_{j=1}^n p_{a_j}.$$

Bernoulli Shifts

The Kolmogorov Extension Theorem allows one to extend μ_p to every set in the Borel σ -algebra \mathcal{B} of Σ . (Note that the cylinder sets form a basis of the topology.)

The measure space $(\Sigma, \mathcal{B}, \mu_p)$ can be made into a measure preserving dynamical system by taking the **left-shift**

$$\sigma(\dots x_{-2}x_{-1} \cdot x_0x_1x_2\dots) = \dots x_{-2}x_{-1}x_0 \cdot x_1x_2\dots$$

Exercise: Show that the set of periodic points of $(\Sigma, \mathcal{B}, \mu_p; \sigma)$ has measure zero if and only if $p_i < 1$ for all i .

Exercise: Show that the set of points in $(\Sigma, \mathcal{B}, \mu_p; \sigma)$ with a dense orbit has measure one if and only if $p_i > 0$ for all i .

Isomorphic Systems

Definition: Two measure preserving dynamical systems (X, \mathcal{B}, T, μ) and (Y, \mathcal{C}, S, ν) are called **isomorphic** if there are $X' \in \mathcal{B}$, $Y' \in \mathcal{C}$ and $\phi : Y' \rightarrow X'$ such that

- ▶ $\mu(X') = 1, \nu(Y') = 1$;
- ▶ $\phi : Y' \rightarrow X'$ is a bi-measurable bijection;
- ▶ ϕ is measure preserving: $\nu(\phi^{-1}(B)) = \mu(B)$ for all $B \in \mathcal{B}$.
- ▶ $\phi \circ S = T \circ \phi$.

That is, the below diagram commutes, and $\phi : Y \rightarrow X$ is one-to-one almost everywhere.

$$\begin{array}{ccc} (Y, \mathcal{C}, \nu) & \xrightarrow{S} & (Y, \mathcal{C}, \nu) \\ \phi \downarrow & & \downarrow \phi \\ (X, \mathcal{B}, \mu) & \xrightarrow{T} & (X, \mathcal{B}, \mu) \end{array}$$

Isomorphic Systems

Example: The doubling map

$$T : \mathbb{S}^1 \rightarrow \mathbb{S}^1, \quad T(x) = 2x \bmod 1$$

with Lebesgue measure is isomorphic to the one-sided $(\frac{1}{2}, \frac{1}{2})$ -Bernoulli shift $(\Sigma, \mathcal{B}, \sigma, \mu)$. The (inverse of the) isomorphism is the coding map $\phi^{-1} : X' \rightarrow \Sigma'$:

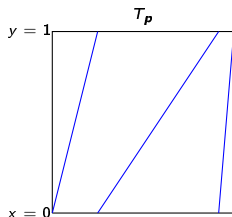
$$\phi^{-1}(x)_k = \begin{cases} 1 & \text{if } T^k(x) \in [0, \frac{1}{2}), \\ 2 & \text{if } T^k(x) \in (\frac{1}{2}, 1]. \end{cases}$$

Here $X' = [0, 1] \setminus \{\text{dyadic rationals in } (0, 1)\}$ because these dyadic rationals map to $\frac{1}{2}$ under some iterate of T , and at $\frac{1}{2}$ the coding map is not well-defined. Note that

$$\Sigma' = \{0, 1\}^{\mathbb{N}} \setminus \{v10^\infty, v01^\infty : v \text{ is a finite word in } \{0, 1\}\}.$$

Isomorphic Systems

Example: Let (p_1, \dots, p_N) be some probability vector with all $p_i > 0$. Then the one-sided (p_1, \dots, p_N) -Bernoulli shift is isomorphic to $([0, 1], \mathcal{B}, T, \text{Leb})$ where $T : [0, 1] \rightarrow [0, 1]$ has N linear branches of slope $1/p_i$.



The map T_p for
 $p = (\frac{1}{4}, \frac{2}{3}, \frac{1}{12})$

The one-sided (p_1, \dots, p_N) -Bernoulli shift is **also** isomorphic to

$([0, 1], \mathcal{B}, S, \nu)$ where $S(x) = Nx \bmod 1$.

But here ν is another measure that gives $[\frac{i-1}{N}, \frac{i}{N}]$ the mass p_i , and $[\frac{i-1}{N} + \frac{j-1}{N^2}, \frac{i-1}{N} + \frac{j}{N^2}]$ the mass $p_i p_j$, etc.

Ornstein Theorem

The **entropy** of the Bernoulli system is defined as:

$$h(p) := - \sum_{i=1}^N p_i \log p_i \quad \text{Convention: } 0 \log 0 = 0.$$

It is preserved under isomorphism, so

Two isomorphic Bernoulli system have the same entropy.

Theorem of Ornstein: Two 2-sided Bernoulli systems are isomorphic **if and only if** they have the same entropy.

Exercise: The 2-sided $(\frac{1}{4}, \frac{1}{4}, \frac{1}{4}, \frac{1}{4})$ shift is isomorphic to the 2-sided $(\frac{1}{8}, \frac{1}{8}, \frac{1}{8}, \frac{1}{8}, \frac{1}{2})$ -shift, but to no other 2-sided shift on ≤ 4 symbols.

Natural Extensions

Invertible systems cannot be isomorphic to non-invertible systems.
So

$$(\{0,1\}^{\mathbb{Z}}, \mathcal{B} : \sigma) \text{ and } (\{0,1\}^{\mathbb{N}}, \mathcal{B} : \sigma)$$

are not isomorphic.

But there is a construction to make a non-invertible system invertible, namely by passing to the **natural extension**.

Natural Extensions

Definition: Let (X, \mathcal{B}, μ, T) be a measure preserving dynamical system. A system (Y, \mathcal{C}, S, ν) is a **natural extension** of (X, \mathcal{B}, μ, T) if there are $X' \in \mathcal{B}$, $Y' \in \mathcal{C}$ and $\phi : Y' \rightarrow X'$ such that

- ▶ $\mu(X') = 1$, $\nu(Y') = 1$;
- ▶ $S : Y' \rightarrow Y'$ is **invertible**;
- ▶ $\phi : Y' \rightarrow X'$ is a measurable surjection;
- ▶ ϕ is measure preserving: $\nu(\phi^{-1}(B)) = \mu(B)$ for all $B \in \mathcal{B}$;
- ▶ $\phi \circ S = T \circ \phi$.

Any two natural extensions can be shown to be isomorphic, so it makes sense to speak of **the** natural extension.

Natural Extensions

Sometimes natural extensions have explicit formulas. For example the **baker map** $B : [0, 1]^2 \rightarrow [0, 1]^2$,

$$B(x, y) = \begin{cases} (2x, \frac{y}{2}) & \text{if } x < \frac{1}{2}; \\ (2x - 1, \frac{y+1}{2}) & \text{if } x \geq \frac{1}{2}. \end{cases}$$

preserving two-dimensional Lebesgue measure is the natural extension of the doubling map via the factor map $\phi(x, y) = x$.

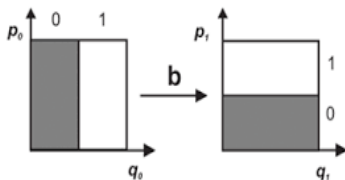


FIG. 1: The baker map propagates a pair of vertical rectangles onto a pair of horizontal rectangles. Inside each separate region the evolution is linear.

Natural Extensions

There is also a general construction: Set

$$Y = \{(x_i)_{i \geq 0} : T(x_{i+1}) = x_i \in X \text{ for all } i \geq 0\}$$

with $S((x_0, x_1, \dots)) = (T(x_0), x_0, x_1, \dots)$. Then S is invertible (with the left shift $\sigma = S^{-1}$) and

$$\nu(A_0, A_1, A_2, \dots) = \inf_i \mu(A_i) \quad \text{for } (A_0, A_1, A_2, \dots) \subset S,$$

is S -invariant.

The factor map $\phi(x_0, x_1, x_2, \dots) := x_0$ satisfies $T \circ \phi = \phi \circ S$.

Also ϕ is measure preserving because, for each $A \in \mathcal{B}$,

$$\phi^{-1}(A) = (A, T^{-1}(A), T^{-2}(A), T^{-3}(A), \dots)$$

and clearly $\nu(A, T^{-1}(A), T^{-2}(A), T^{-3}(A), \dots) = \mu(A)$ because $\mu(T^{-i}(A)) = \mu(A)$ for every i by T -invariance of μ .

The Bernoulli Property

Definition: Let (X, \mathcal{B}, μ, T) be a measure preserving dynamical system.

1. If T is invertible, then the system is called **Bernoulli** if it is isomorphic to a 2-sided Bernoulli shift.
2. If T is non-invertible, then the system is called **one-sided Bernoulli** if it is isomorphic to a 1-sided Bernoulli shift.
3. If T is non-invertible, then the system is called **Bernoulli** if its natural extension is isomorphic to a 2-sided Bernoulli shift.

The Bernoulli Property

The third Bernoulli property is quite general, even though the isomorphism ϕ may be very difficult to find explicitly. Thus, proving that a system is **not** Bernoulli can be hard.

Expanding circle maps $T : \mathbb{S}^1 \rightarrow \mathbb{S}^1$ that satisfy the conditions of the Folklore Theorem are also Bernoulli, i.e., have a Bernoulli natural extension (proven by Ledrappier).

Being **one-sided Bernoulli**, on the other hand, is quite special. If T is piecewise C^2 but not piecewise linear, then it has to be C^2 -conjugate to a piecewise linear expanding map to be one-sided Bernoulli.

Bernoulli (of any of the above forms) **implies mixing**. So non-mixing systems (e.g. circle rotations) are not Bernoulli.