Hilbert-Schmidt integral operators vs. systems of elliptic Calogero-Moser type

Simon Ruijsenaars

School of Mathematics University of Leeds

ESI Vienna EIS Workshop, 23 March 2017

4 **A** N A **B** N A **B** N









1. Introduction

• Assume *F* is a bounded subset of \mathbb{R}^N and K(x, y) is a function that is continuous on $\overline{F} \times \overline{F}$ (with \overline{F} the closure of *F*). Thus we have

$$\int_{F^2} |K(x,y)|^2 dx dy < \infty.$$

As a consequence the integral operator

$$(\mathcal{I}f)(x) := \int_F K(x,y)f(y)dy,$$

is a Hilbert-Schmidt operator on $\mathcal{H} := L^2(F, dx)$.

 Recall this entails that there exist two sets of pairwise orthogonal unit vectors f₀, f₁,... and g₀, g₁,... such that

$$\mathcal{K}(x,y) = \sum_{m=0}^{M} s_m f_m(x) \overline{g_m(y)}, \ s_0 \ge s_1 \ge s_2 \ge \cdots > 0, \ M \le \infty,$$

with the so-called singular values s_m satisfying $\sum s_m^2 < \infty$.

• The self-adjoint operators $\mathcal{I}^*\mathcal{I}$ and \mathcal{II}^* are then given by

$$\mathcal{I}^*\mathcal{I} = \sum_{m=0}^M s_m^2 g_m \otimes \overline{g_m}, \ \mathcal{II}^* = \sum_{m=0}^M s_m^2 f_m \otimes \overline{f_m},$$

so they are trace class and non-negative.

- Let us call an HS operator complete when it has trivial null space and dense range. Equivalently, the vectors f_0, f_1, \ldots and g_0, g_1, \ldots are ONBs (orthonormal bases) for \mathcal{H} .
- It seems there is no useful general way to recognize a complete HS operator when you meet one.
- For large classes of special HS operators, however, completeness can be shown. The proofs only involve elementary Fourier analysis (S. R., 2013). These results apply in particular to the elliptic kernel functions of my survey talk, provided the parameters are suitably restricted.

 A long-standing goal is to reinterpret the 2N commuting A_{N-1} AΔOs A_{k,δ}(x) from my survey talk as commuting self-adjoint operators on the Hilbert space

$$\mathcal{H}_{A}:=L^{2}(F_{A},dx),$$

$$F_{A} := \{-\pi/2r < x_{N} < \cdots < x_{1} \le \pi/2r\}.$$

(They are at least formally self-adjoint on \mathcal{H}_A , by contrast to $A_{k,\delta}(x) = C(x)\mathcal{A}_{k,\delta}(x)C(x)^{-1}$.)

 Likewise, the 2 commuting BC_N A∆Os A_δ(x) ought to be promoted to commuting self-adjoint operators on the Hilbert space

$$\mathcal{H}_B := L^2(F_B, dx),$$

$$F_B := \{ 0 < x_N < \cdots < x_1 \le \pi/2r \}.$$

• To this end, we need 'only' show existence of an ONB of joint eigenfunctions with real eigenvalues.

 Under suitable restrictions on the parameters, the kernel functions from my survey talk give rise to complete HS integral operators *I*_ξ, ξ ∈ C, and *I* on *H*_A and *H*_B, resp. Thus the operators

$$\mathcal{T}_{\xi} := \mathcal{I}_{\xi} \mathcal{I}_{\xi}^*, \quad \mathcal{T} := \mathcal{I} \mathcal{I}^*,$$

are positive trace class operators.

- Crux: it can be expected that the *T*_ξ- and *T*-eigenvectors extend to meromorphic eigenfunctions of the AΔOs *A*_{k,δ} and *A*_δ with real eigenvalues.
- Reason: the A Δ Os are formally self-adjoint and formally satisfy

$$[\mathcal{A}_{k,\delta},\mathcal{T}_{\xi}]=0, \quad [\mathcal{A}_{\delta},\mathcal{T}]=0,$$

due to the kernel identities. Thus the eigenvector ONB of the trace class operators 'should' yield an ONB of joint eigenfunctions of the commuting $A\Delta Os$.

- This approach is easily understood and formally convincing, but a lot of analysis is needed to make it work. This involves in particular complex analysis to prove the meromorphy of the *T*-eigenfunctions, and functional analysis to control dense domains for the AΔOs. (No general Hilbert space theory for AΔOs exists to date.)
- It can be expected that a similar approach applies to the nonrelativistic (PDO) case. A difficulty in this setting is that the eigenfunctions are (generically) not meromorphic. However, for the rank-one cases (i. e. Lamé and Heun), one can invoke Sturm-Liouville and Frobenius theory to push it through. This gives rise to a novel S₄ spectral invariance of Heun Hamiltonians (S. R., 2009). In this seminar we only supply some further information about the AΔO case.

2. The A_{N-1} case

 For the N = 2 (relativistic Lamé) case and special couplings, the 'expected' results for the AΔOs A_± were shown to hold true (without using kernel functions) some 15 years ago (S. R., 2003). Specifically, letting

$$b = (N_+ + 1)a_+ - N_-a_- \in (0, a_+ + a_-),$$

with

$$N_+,N_-\in\mathbb{N}:=\{0,1,2,\ldots\},\ a_+/a_-\notin\mathbb{Q},$$

the Hilbert space $\mathcal{H} = L^2((-\pi/2r, \pi/2r], dx)$ has an ONB that consists of restrictions of meromorphic joint eigenfunctions with real eigenvalues to $(-\pi/2r, \pi/2r]$. A crucial ingredient in the proof is a coupled system of Bethe Ansatz equations.

Specializing my recent results for the BC₁ case, this joint eigenfunction ONB admits an interpolation to any b ∈ (0, a₊ + a₋). These BC₁ results hinge on suitable use of the BC₁ kernel function and are sketched below.

 For N > 2 there is work in progress; requiring once more b ∈ (0, a₊ + a_−), there is circumstantial evidence for the conjecture that the joint eigenfunction ONB can be labelled by

$$n \in \mathbb{Z}^N_{\geq} \equiv \{n \in \mathbb{Z}^N \mid n_1 \geq \cdots \geq n_N\},\$$

in such a way that when the minimum of the gaps $n_j - n_{j+1}$, j = 1, ..., N - 1, tends to ∞ one has asymptotics proportional to

$$\sum_{\sigma \in S_N} \frac{C(x_{\sigma})}{C(x)} \exp(2irn \cdot x_{\sigma}).$$

(If so, the commuting dual dynamics yield a factorized S-matrix.)

For the cases b = a₊ and b = a_− the joint eigenfunction ONB amounts to 'free fermions' (~ Schur polynomials), the A∆O-eigenvalues are obvious, and the eigenvalues for a modified HS family are explicitly known too (S. R., 2009).

3. The BC₁ case

• Here the 'initial' kernel identity reads

$$\mathcal{A}_{\delta}(\gamma; \mathbf{x}) \mathcal{S}(\sigma(\gamma); \mathbf{x}, \mathbf{y}) = \mathcal{A}_{\delta}(\gamma'; \mathbf{y}) \mathcal{S}(\sigma(\gamma); \mathbf{x}, \mathbf{y}), \quad \delta = +, -,$$

where

$$egin{aligned} &\gamma'\equiv -J\gamma,\ &\sigma(\gamma)\equiv -rac{1}{4}\sum_{\mu=0}^7\gamma_\mu=-rac{1}{4}\langle\zeta,\gamma
angle,\quad \zeta\equiv(1,\ldots,1), \end{aligned}$$

and J can be viewed as the reflection associated with the highest E_8 root $\zeta/2$, i. e.,

$$J \equiv \mathbf{1}_8 - \frac{1}{4}\zeta \otimes \zeta.$$

The BC₁ kernel function is given by

$$\mathcal{S}(t; x, y) \equiv \prod_{\delta_1, \delta_2 = +, -} G(\delta_1 x + \delta_2 y - ia + it),$$

with G(z) the elliptic gamma function.

< ロ > < 同 > < 回 > < 回 >

For the AΔOs A_±(γ; x) the relevant Hilbert space is the weighted L² space

$$\mathcal{H}_{w} \equiv L^{2}([0, \pi/2r], w_{e}(\gamma; x) dx).$$

• It is crucial to switch from this A Δ O pair to the D_8 -invariant A Δ Os

$$\mathcal{A}_{\delta}(\gamma; \mathbf{x}) = c_{e}(\gamma; \mathbf{x})^{-1} \mathcal{A}_{\delta}(\gamma; \mathbf{x}) c_{e}(\gamma; \mathbf{x}),$$

which are formally self-adjoint on

$$\mathcal{H}=L^2([0,\pi/2r],dx),$$

for suitable γ (in particular for $\gamma \in \mathbb{R}^8$).

They satisfy the kernel identity

$$\mathcal{A}_{\delta}(\gamma; \mathbf{x})\mathcal{K}(\gamma; \mathbf{x}, \mathbf{y}) = \mathcal{A}_{\delta}(\gamma'; -\mathbf{y})\mathcal{K}(\gamma; \mathbf{x}, \mathbf{y}),$$

with

$$\mathcal{K}(\gamma; \mathbf{x}, \mathbf{y}) \equiv \frac{\mathcal{S}(\sigma(\gamma); \mathbf{x}, \mathbf{y})}{c_{e}(\gamma; \mathbf{x})c_{e}(\gamma'; -\mathbf{y})}$$

 With further restrictions on γ, the kernel function K(γ; x, y) yields a complete HS integral operator I(γ) on H. Requiring γ ∈ ℝ⁸ from now on, it suffices to further restrict γ by

$$\gamma_{\mu}, \ \gamma'_{\mu} \in (-a, a), \ \ \sigma(\gamma) \in (0, a).$$

- With this restriction, we can show that the resulting eigenvector *H*-ONB *f_n*(*γ*), *n* = 0, 1, 2, ..., for the self-adjoint trace class operator *I*(*γ*)*I*(*γ*)* has the following features:
 - $f_n(\gamma)$ is the restriction to $[0, \pi/2r]$ of a meromorphic function $f_n(\gamma; x)$ with known pole locations depending only on γ ;
 - Setting

$$a_s \equiv \min(a_+, a_-), a_l \equiv \max(a_+, a_-),$$

and assuming a_l is not a multiple of a_s , the functions $f_n(\gamma; x)$ are joint eigenfunctions of $A_{\pm}(\gamma; x)$ with real eigenvalues.

- Consequence: With the above restrictions on a_± and γ understood, the AΔOs give rise to commuting self-adjoint operators Â_±(γ) on H with discrete spectra.
- Further results include:
 - The definition of $\hat{\mathcal{A}}_{\pm}(\gamma)$ implies that the operators are invariant under D_8 -transformations of γ .
 - For γ in the ball $\|\gamma\|_2 < a$ (with the origin deleted), the operators are isospectral under E_8 -transformations. Generically, this yields 135 (= $|W(E_8)/W(D_8)|$) distinct isospectral operators.
 - For generic γ, we also get 64 distinct commuting HS operators.
 The asymptotic behavior as n → ∞ of the eigenfunctions f_n(γ; x) is the same as that of an *H*-ONB of functions P_n(γ; x)/c_P(γ; x), with P_n(γ; x) orthonormal polynomials; this relation also leads to detailed information on eigenvalue asymptotics.

Recent references re HS approach:

- Hilbert-Schmidt-operators vs. integrable systems of elliptic Calogero-Moser type. II. The A_{N-1} case: First steps, Comm. Math. Phys. **286** (2009), 659–680
- Hilbert-Schmidt-operators vs. integrable systems of elliptic Calogero-Moser type. III. The Heun case, SIGMA 5 (2009), 049, 21 pages
- On positive Hilbert-Schmidt operators, Integr. Equ. Oper. Theory, 75 (2013), 393–407
- Hilbert-Schmidt-operators vs. integrable systems of elliptic Calogero-Moser type. IV. The relativistic Heun (van Diejen) case, SIGMA 11 (2015), 004, 78 pages