

That which is not forbidden is  
compulsory. – T.H. White

That which is not forbidden is  
compulsory.

Unless it is extremely unlikely.

**Free surface flows: Topological: Wave crater collapse**

Axisymmetric:

$\delta \sim (t_o - t)^{2/3}$  length scales

$v \sim (t_o - t)^{-1/3}$  velocity scales

$Pr(v) \sim |v|^{-4}$  velocity probability distribution

Similarity solutions appear to exist

2-D like axisymmetric!

**Quantum fluid flows: Topological: Vortex reconnection**

$\delta \sim (t_o - t)^{1/2}$  length scales

$v \sim (t_o - t)^{-1/2}$  velocity scales

$Pr(v) \sim |v|^{-3}$  velocity probability distribution in turbulence

Similarity solutions exist

Fixed points with reconnection geometry exists

Safety!

Danger!

**Navier-Stokes flows: ?: turbulence**

$Pr(v_i) \sim \exp(-v_i^2/\sigma_v^2)$

$Pr(\partial_j v_i) \sim \exp(-|v_i|/\sigma_s)$

$Pr(\epsilon)$  dissipation interesting!

$Pr(\Omega)$  enstrophy interesting!

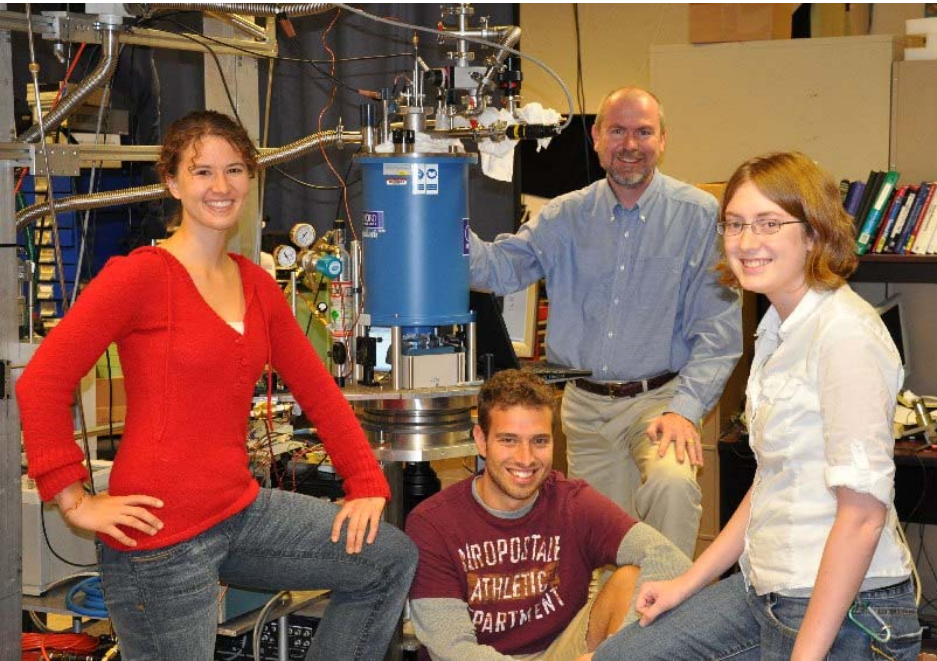
**Euler flows: ?: ?**

Fixed points associated with reconnection?

# Singular events in fluid flow due to topological change

Dan Lathrop

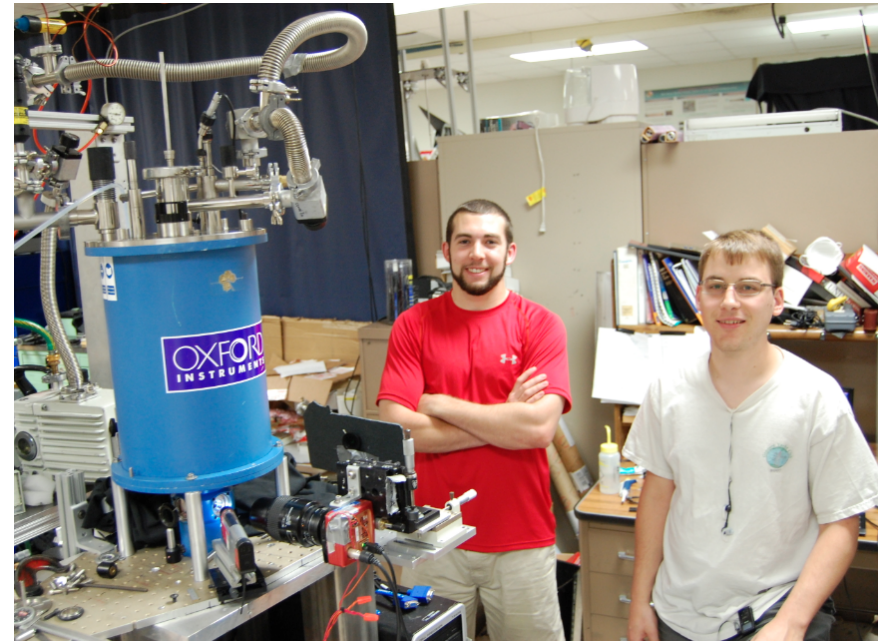
**University of Maryland**  
**National Science Foundation**



**Kristy (Gaff)  
Johnson**

**Matt Paoletti**

**Kaitlyn Tuley**



**Chris Boughter and David Meichle**



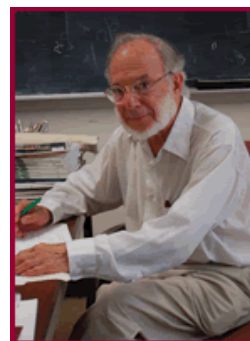
**Enrico Fonda**



**Greg Bewley**



**K.R. Sreenivasan**

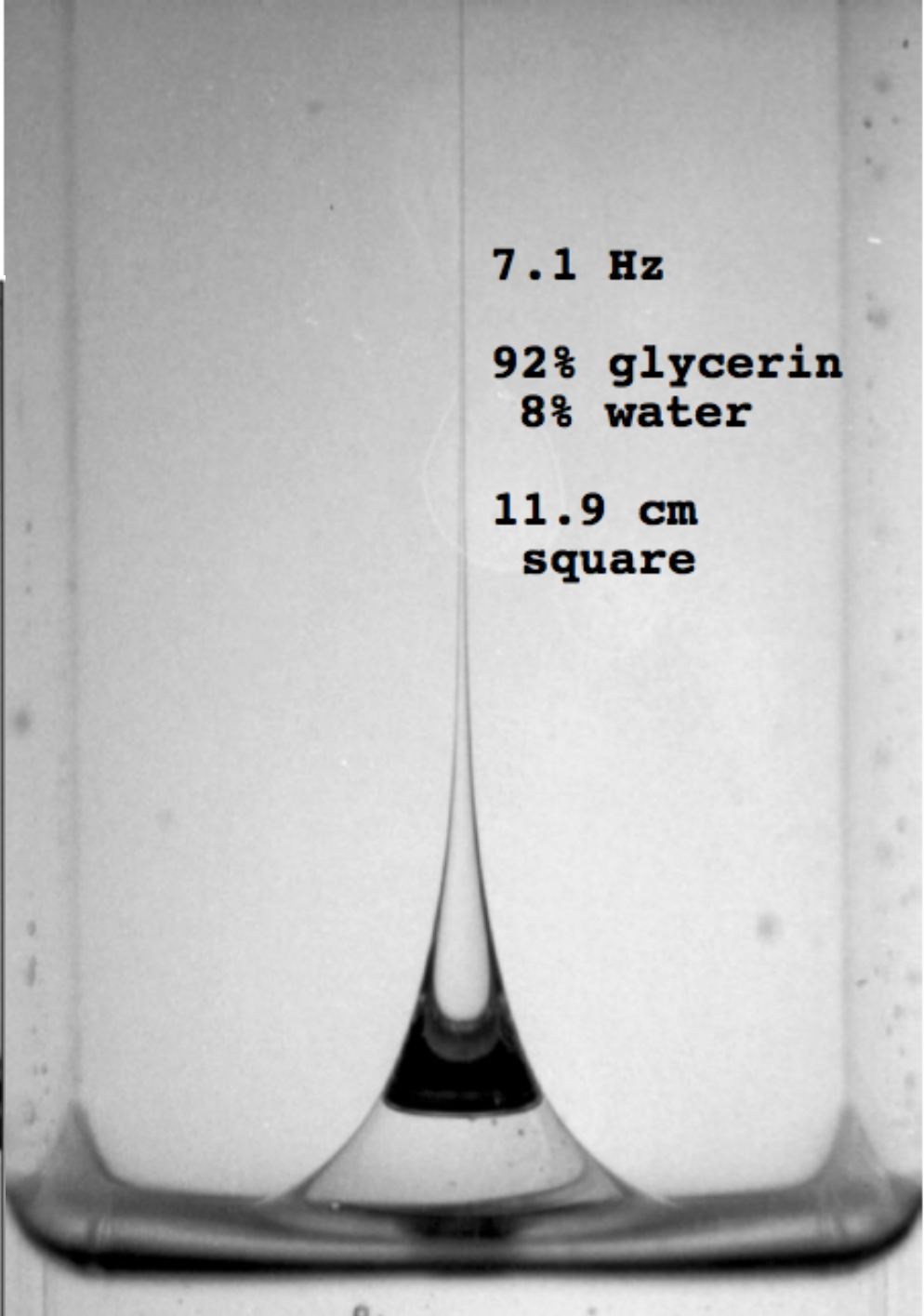
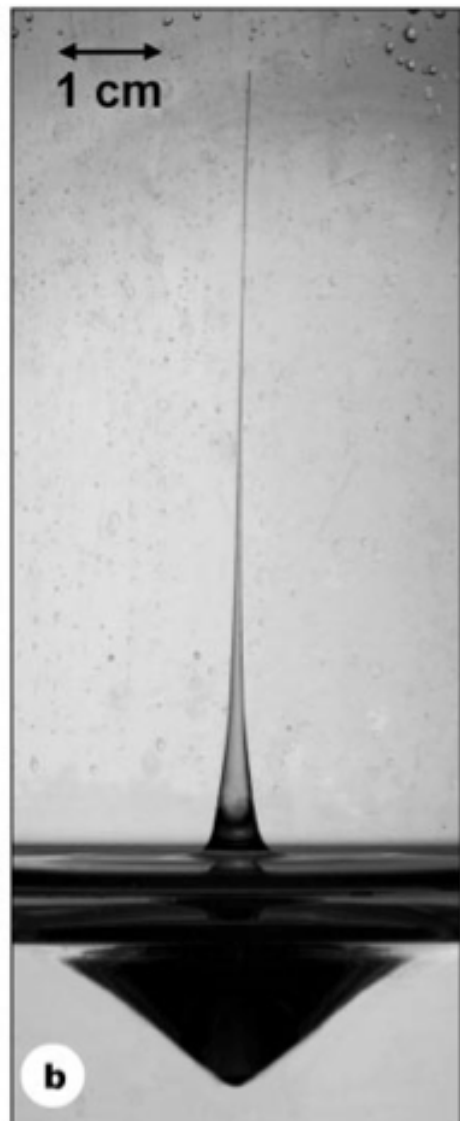
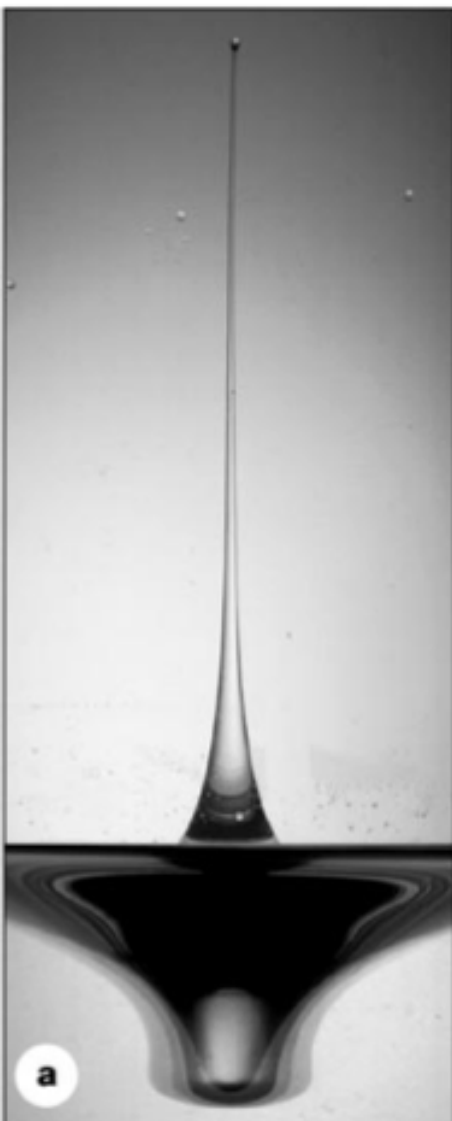


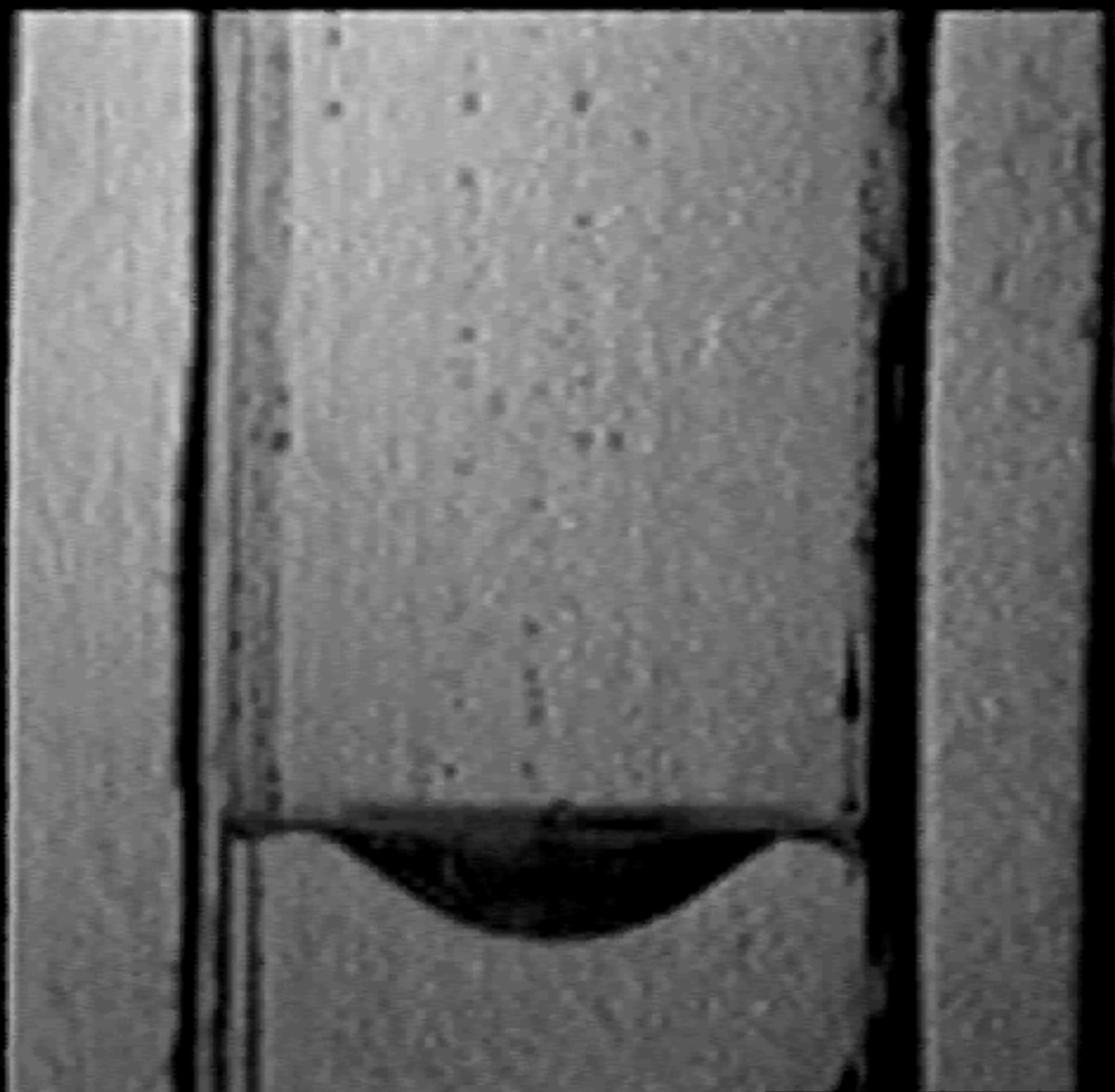
**Michael Fisher**

## **Outline:**

- 1) Free surface singularities**
- 2) Quantum vortices**
- 3) Navier-Stokes turbulence**
- 4) Euler flow**







STATUS

MODE

PLAYBACK

FRAME #

142

TIME OF FRAME

508 MS

EVENT NUMBER

3187

SETUP

F/SEC RECORD

250

SHUTTER SPEED

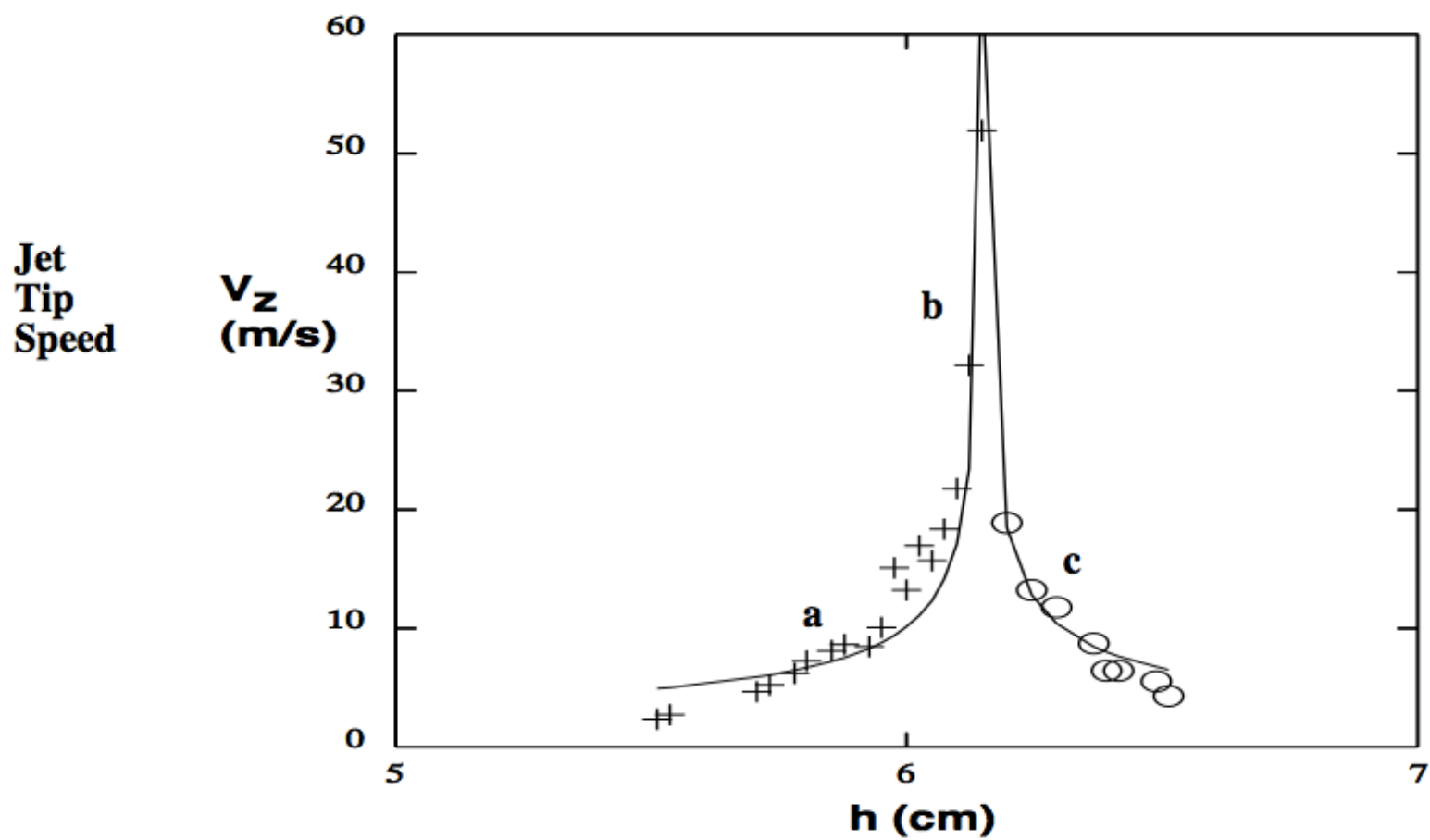
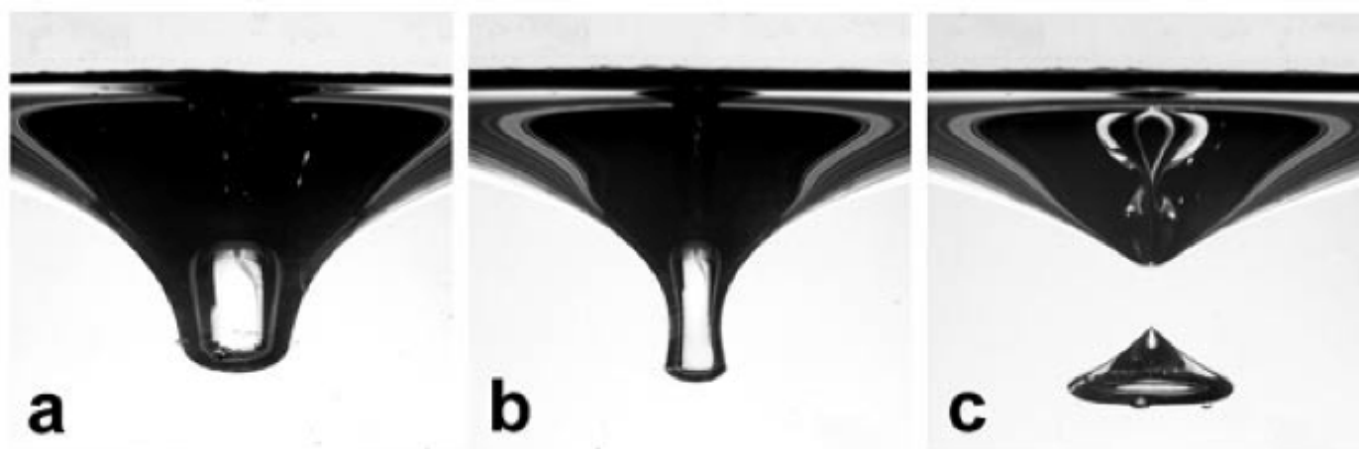
1X

TRIGGER POINT

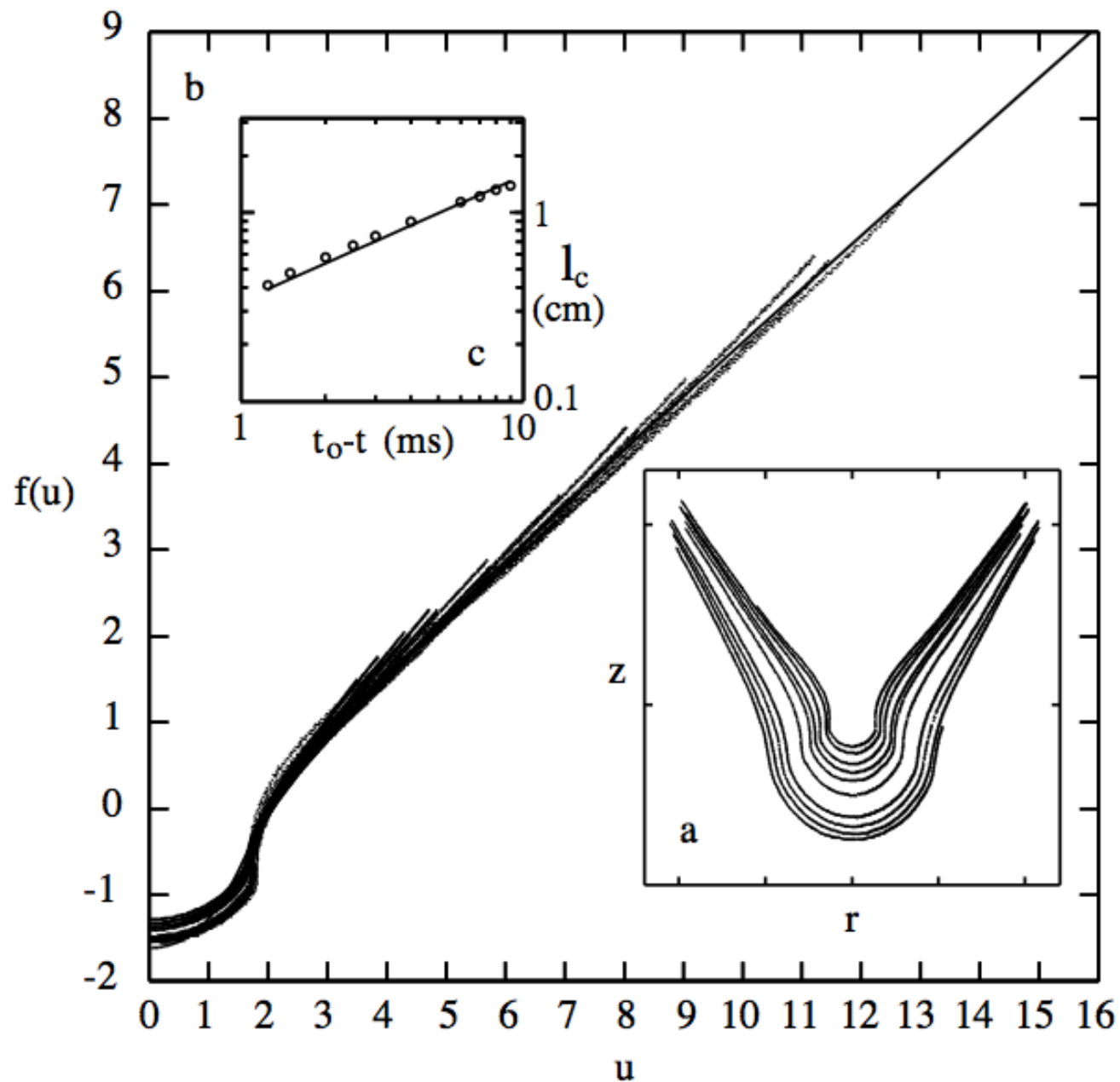
0Z ▲

F/SEC PLAY

30



# Test of Similarity Solution



$$f(u) = z/\tau^{2/3}$$

$$u = r/\tau^{2/3}$$

### Similarity Solution $t < 0$

velocity  $\vec{u} = \vec{\nabla}\phi$

height  $h(r, t)$

$$\nabla^2\phi = 0 \quad (1)$$

$$\frac{\partial h}{\partial t} + (\vec{u} \cdot \vec{\nabla})h = u_z \quad (2)$$

$$\frac{\partial\phi}{\partial t} + \frac{1}{2}(\vec{\nabla}\phi)^2 + \frac{\sigma}{\rho}\left(\frac{1}{R_1} + \frac{1}{R_2}\right) = 0 \quad (3)$$

Similarity Ansatz around  $\tau = (t_o - t) = 0$

$$h(r, t) = \tau^\alpha f[r/\tau^\alpha] = f[u]$$

$$\phi(r, z, t) = \tau^\gamma g[r/\tau^\alpha, z/\tau^\alpha] = \tau^\gamma g[u, v]$$



# Similarity Equations

$$\alpha=2/3 \quad \gamma=1/3$$

Time drops out of problem

$$\begin{aligned} u &= r / (-t)^{2/3} \\ f(u) = v &= z / (-t)^{2/3} \\ g(u,v) &= (-t)^{1/3} \phi(r,z,t) \end{aligned}$$

$$(1) \quad \nabla^2 g = 0$$

$$(2) \quad \frac{\partial f}{\partial u} = \frac{g_v - 2f/3}{g_u - 2u/3}$$

$$(3) \quad -\frac{g}{3} + \frac{2u}{3}g_u + \frac{2f}{3}g_v + \frac{1}{2}(\nabla g)^2 + \kappa = 0$$

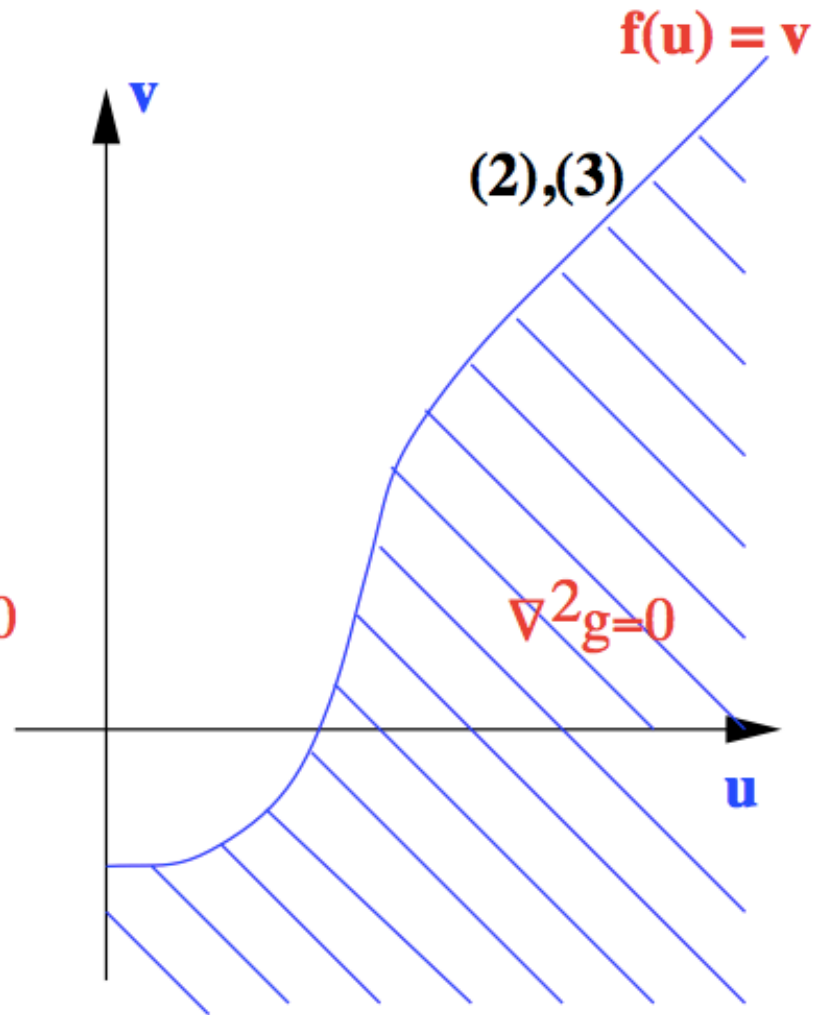
Asymptotics

large radius  $\rightarrow$  cone

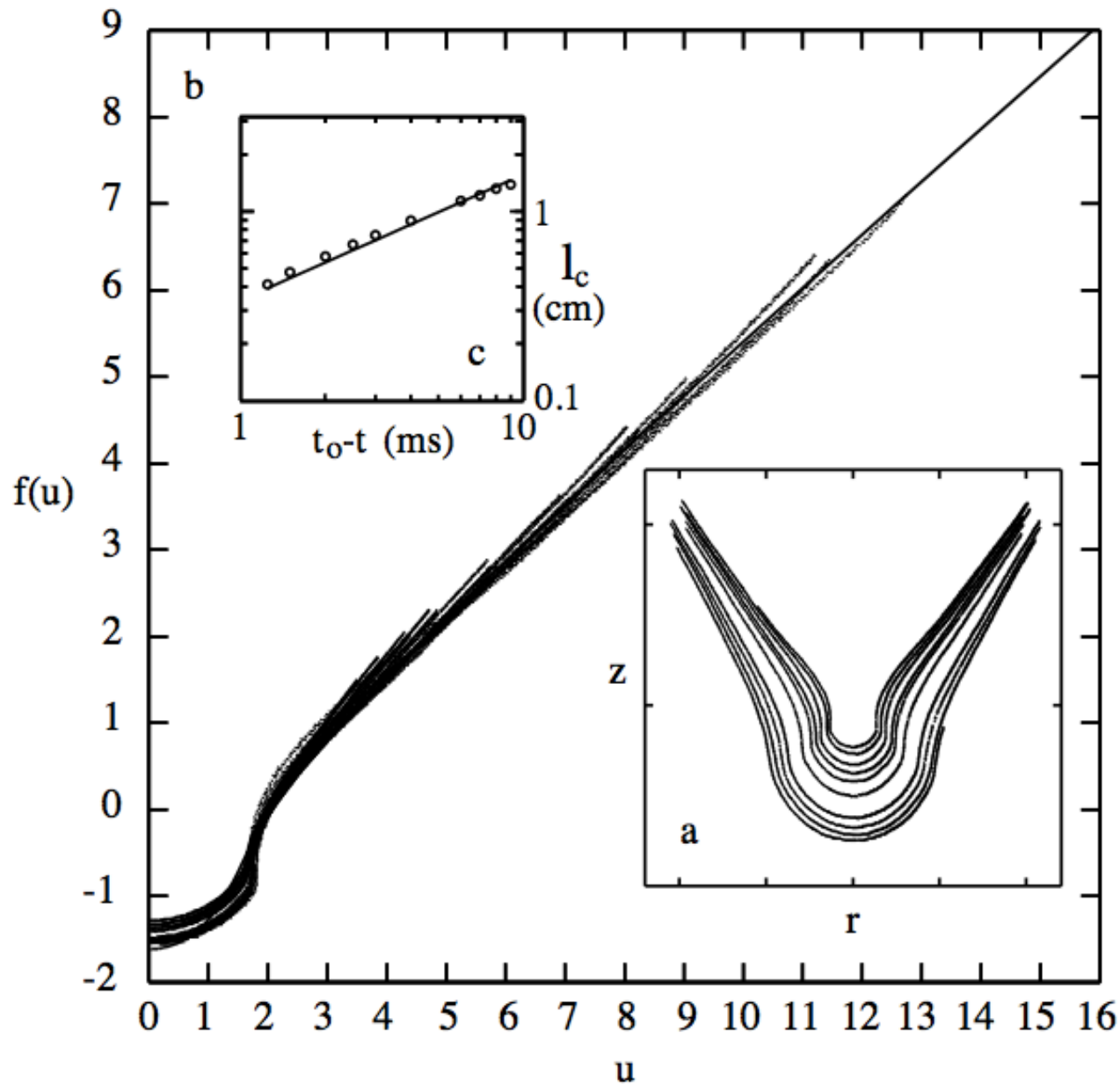
$f \sim u$  angle free

$g \sim u^{1/2}$

small radius  $\rightarrow$  regular



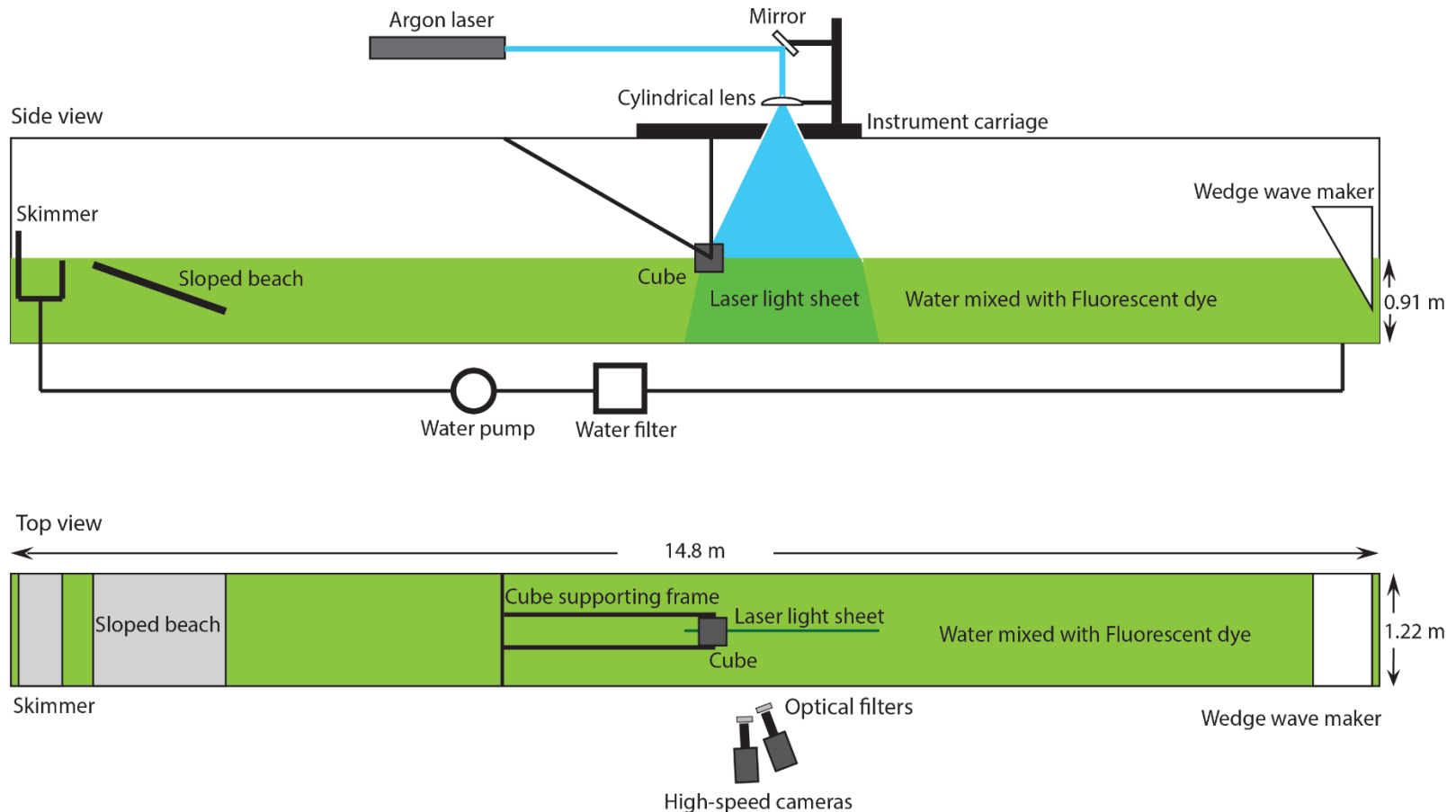
# Test of Similarity Solution



$$f(u) = z/\tau^{2/3}$$

$$u = r/\tau^{2/3}$$

- 2D wave collapse: An Wang, James H. Duncan and Daniel P. Lathrop**

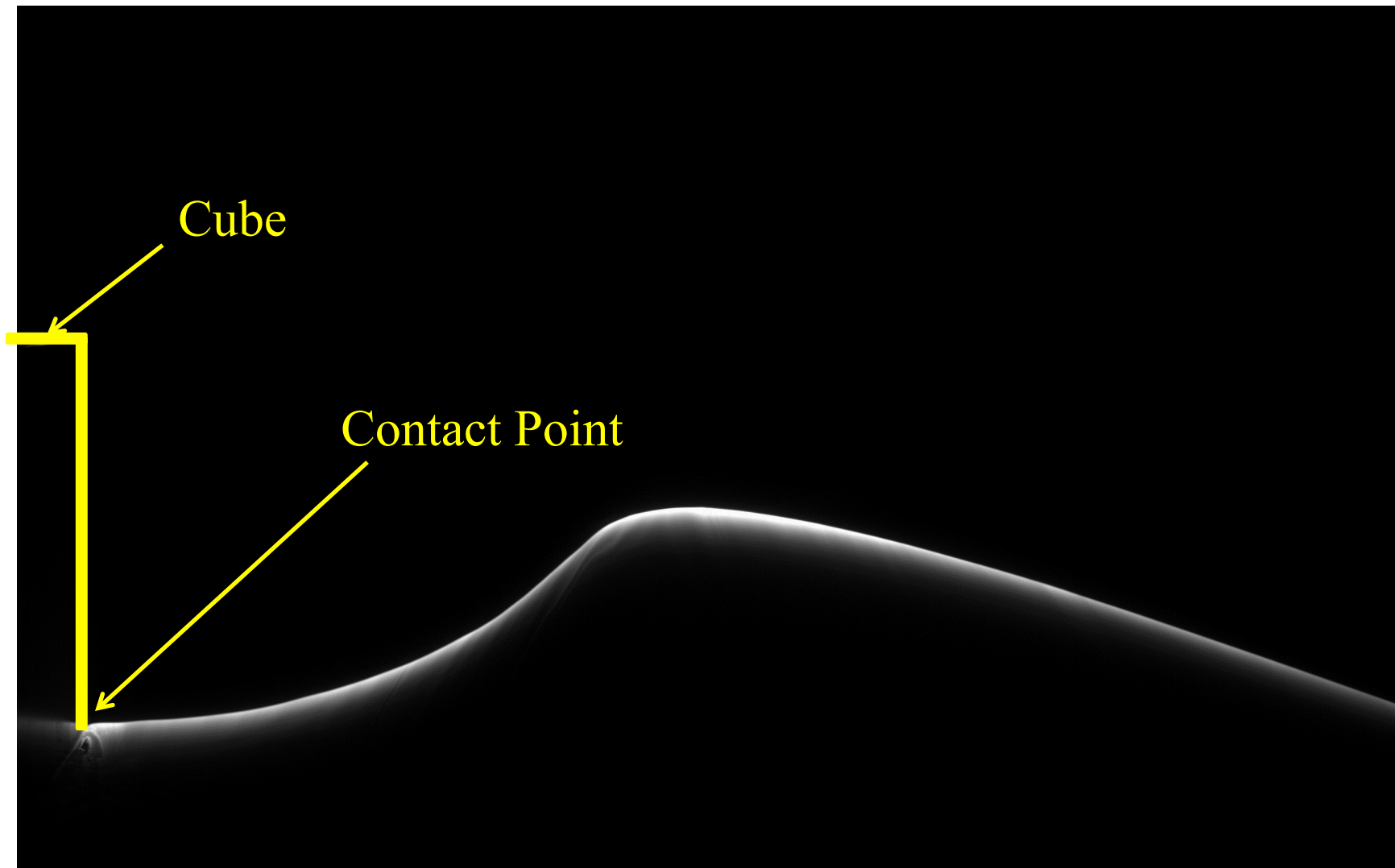


- The cube ( $L = 30.5$  cm) is rigidly mounted at center plane of the wave tank

# High-speed Movie of Wave Impact

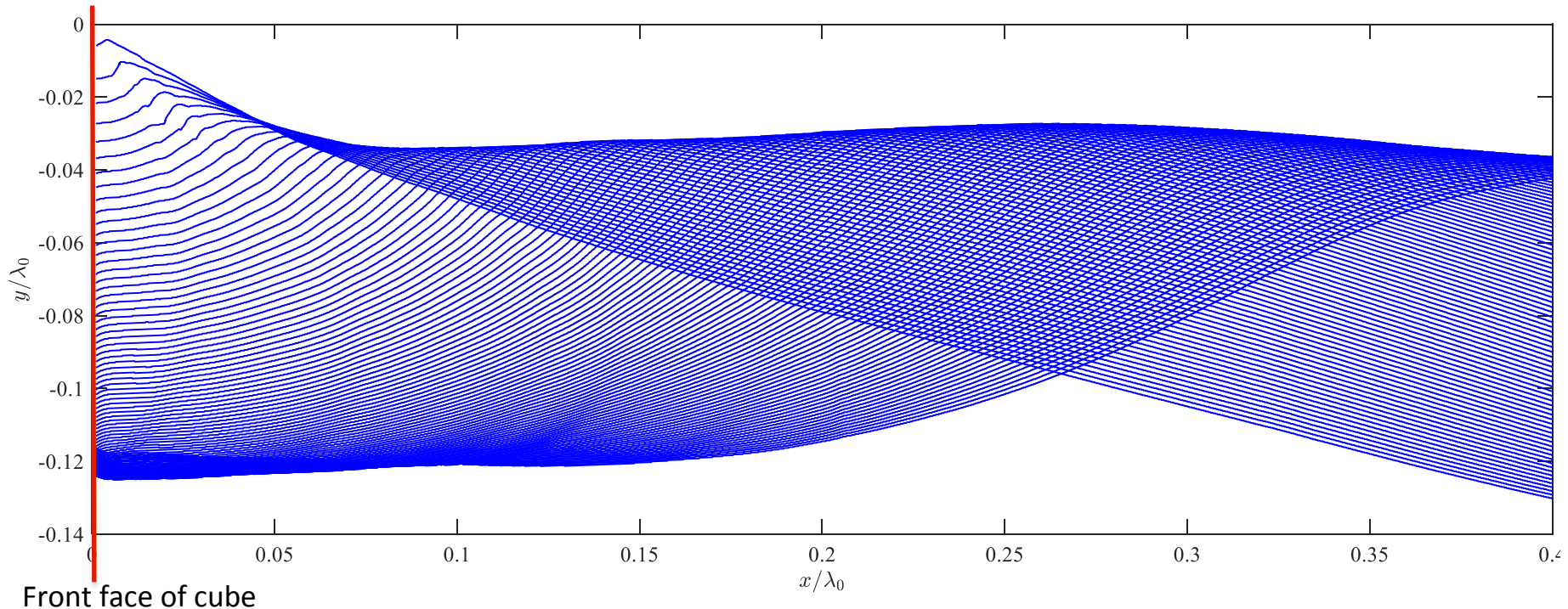
(1500 pps, played at 15 fps, field of view 53 cm)

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# Water Surface Profiles

- The profiles are equally spaced in time, measured at a frame rate of 1500 pps and plotted every 5 frames.
- The last profile represent the **moment of impact** (the water surface between the contact point and the crest collapse to a point and the surface becomes a straight line.)





## Collapse in 2-D Wave-wall interaction

$$(1) \quad \nabla^2 \phi = 0$$

$$(2) \quad \frac{\partial h}{\partial t} + \frac{\partial \phi}{\partial x} \frac{\partial h}{\partial x} = \frac{\partial \phi}{\partial z}$$

$$(3) \quad \frac{\partial \phi}{\partial t} + \frac{1}{2} (\nabla \phi)^2 = 0$$

$$h(x,t) = |t|^\alpha f(x|t|^\alpha)$$

$$\phi(x,z,t) = |t|^\alpha g(x|t|^\alpha, z|t|^\alpha)$$

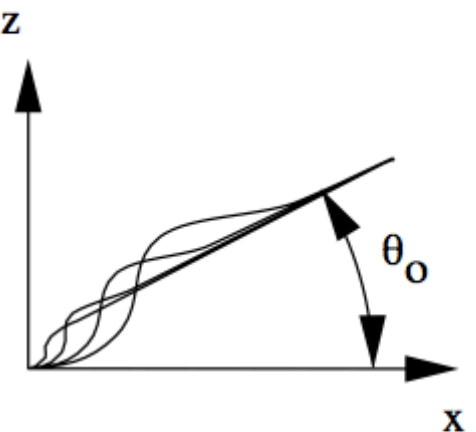
family of exponents allowed  $0 < \alpha < \infty \quad \gamma = 2\alpha - 1$

$$g = a_0 r^{\gamma/\alpha} \cos(\gamma \theta / \alpha) + a_{-1} r^{-1} \cos(\theta) + a_{-2} r^{-2} \cos(2\theta)$$

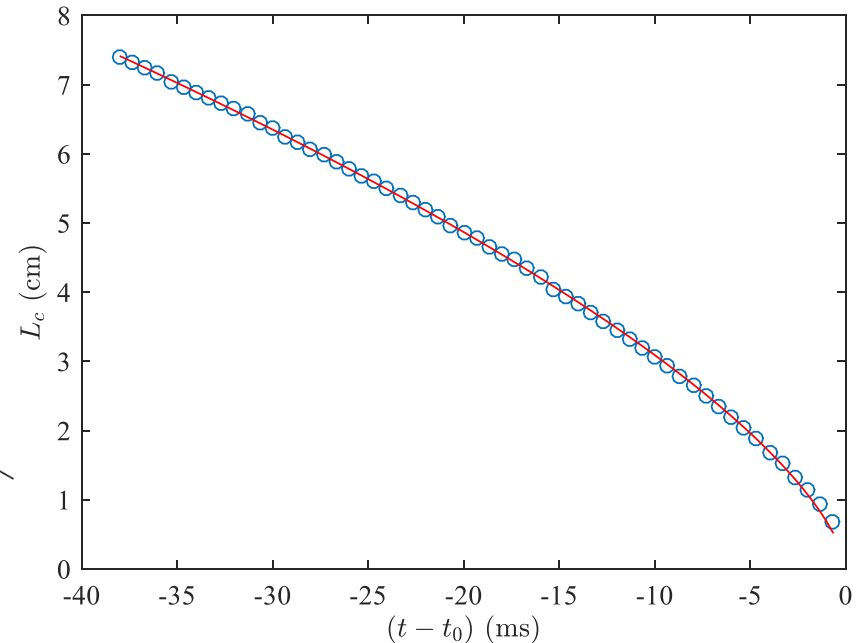
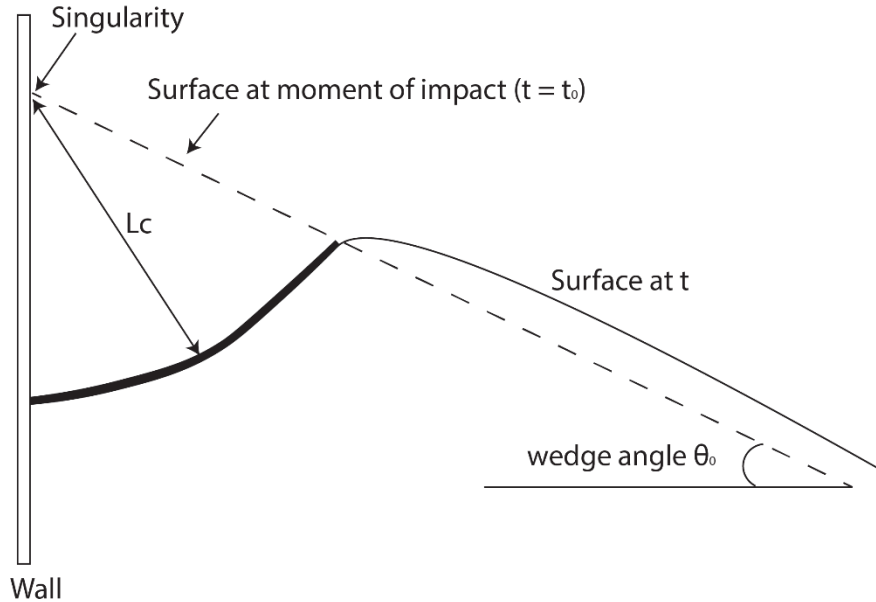
f wedge-like at large x, wedge angle  $\theta_0$

Bounded kinetic energy requires

for  $\alpha > 1/2 \quad a_0 = 0 \Rightarrow$  specifies  $\theta_0$



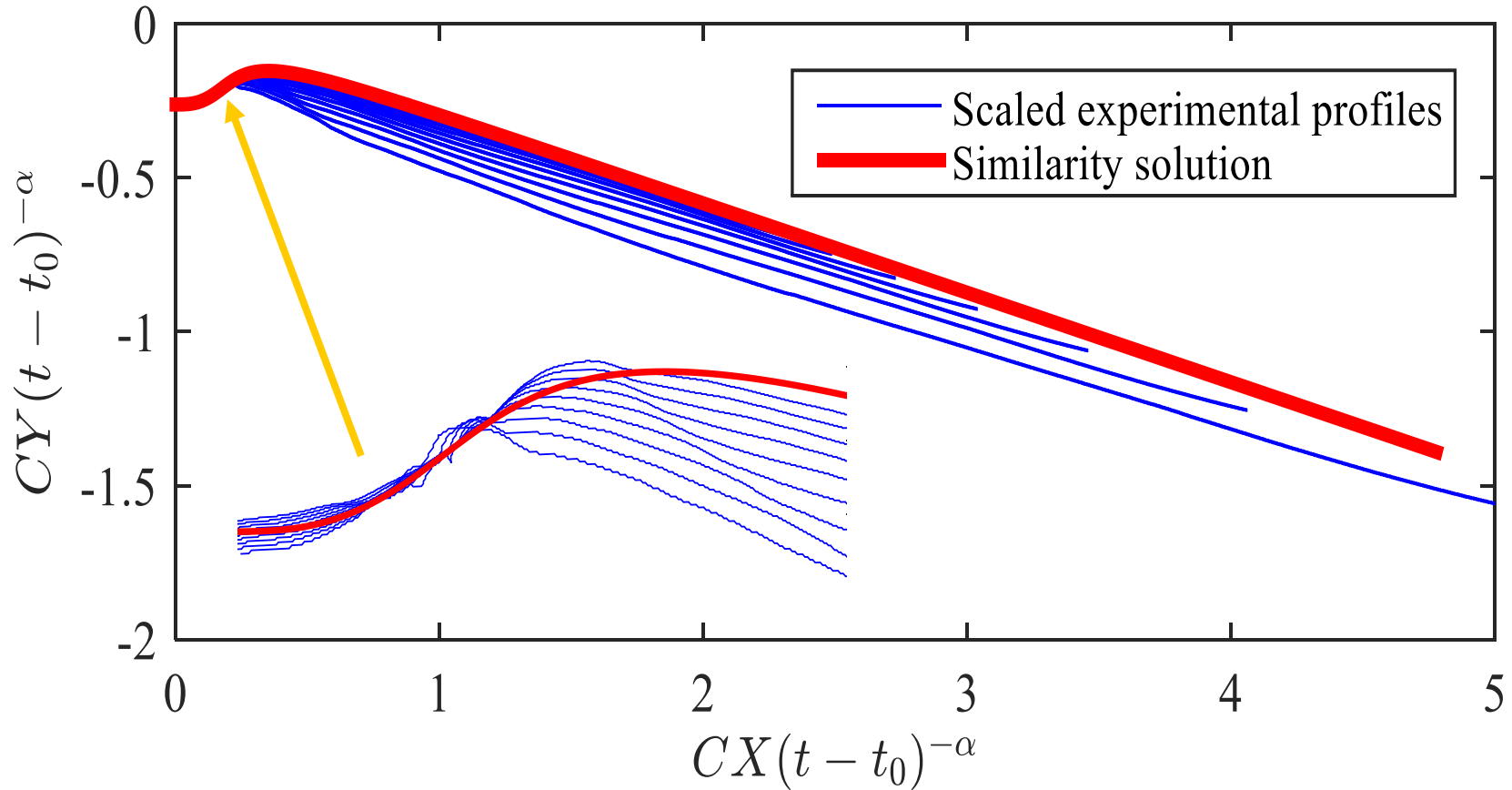
# Collapse to Zero of the Length Scale



- Length scale is defined as the average distance from the origin (singularity point) to points on the curved portion of the surface (between contact point and crest).
- The decay of length scale follows the power law with exponent  $\alpha = 0.655$
- The wedge angle remains nearly a constant about  $15^\circ$

# Comparison

From  $t = 38$  ms to  $t = 0$  ms (57 frames)



# Quantum Fluids

A **state of matter** with long range quantum order

Type of synchronization

partial phase sync of the individual atomic wave functions

E.g.

BEC atomic systems

$^4\text{He}$

$^3\text{He}$

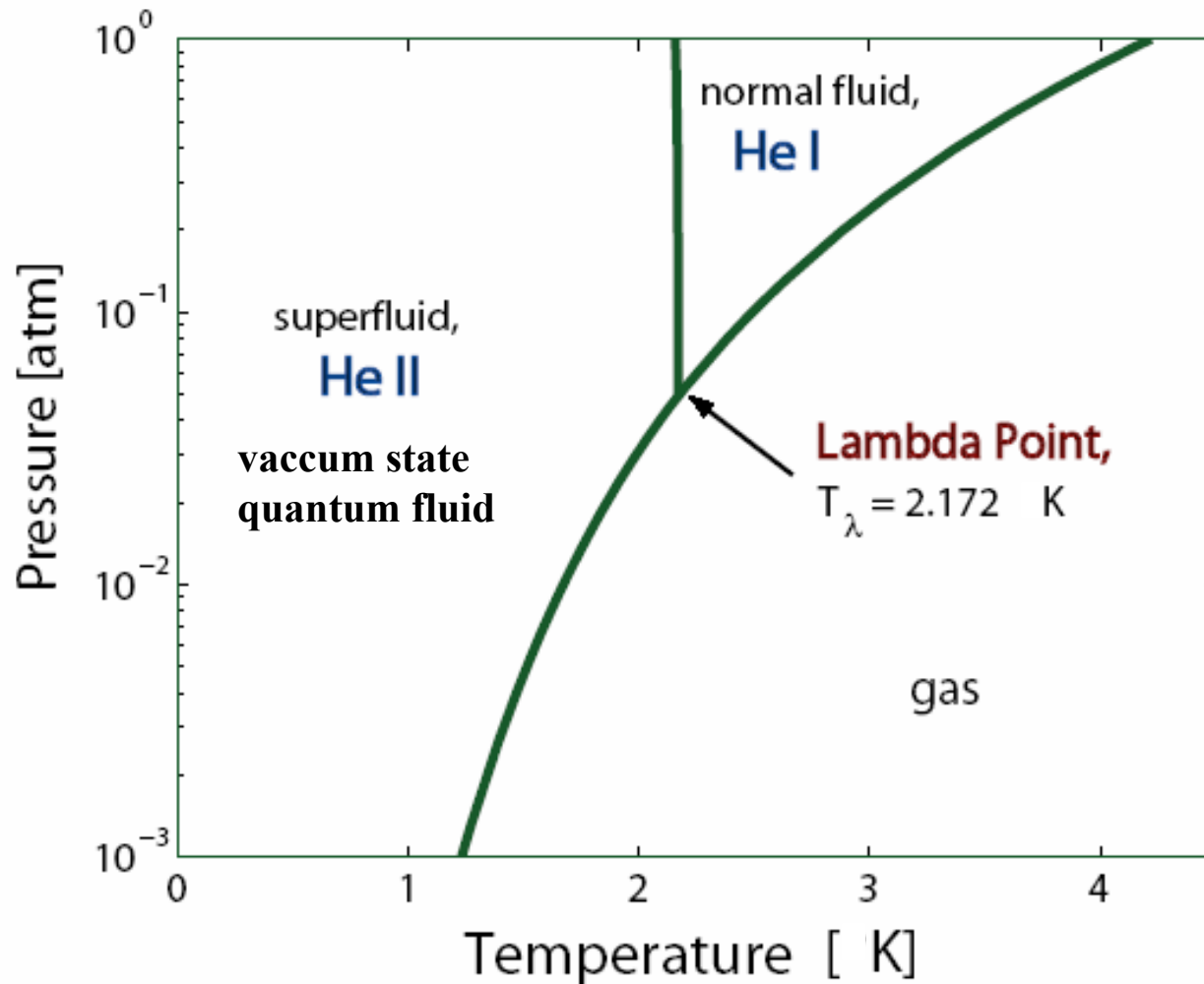
Cooper pair electrons in superconductors

Physical vacuum

**Quantum Turbulence**  $\rightarrow$  turbulence in a quantum fluid

**Why does it matter?**

# Background: Superfluid Helium





# Two-Fluid Model

- Order parameter for superfluid helium is a complex field,

$$\Psi(\mathbf{x}) = Ae^{i\phi}$$

A is amplitude,

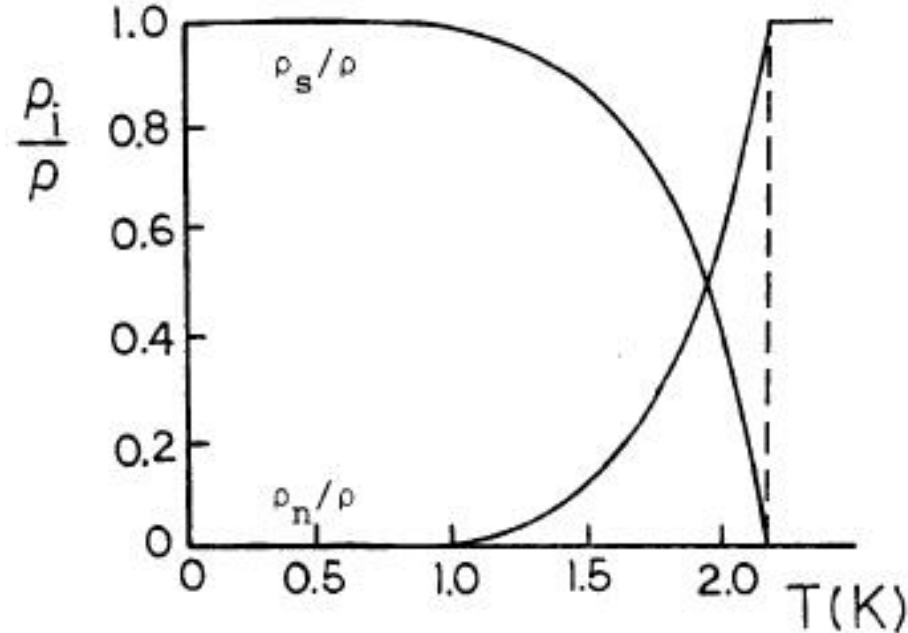
and  $\phi$  is the phase

- Superfluid velocity given by

$$v_s = \kappa \nabla \phi \quad \kappa = \frac{h}{m}$$

$h$  = Planck's constant

$m$  = mass of helium atom

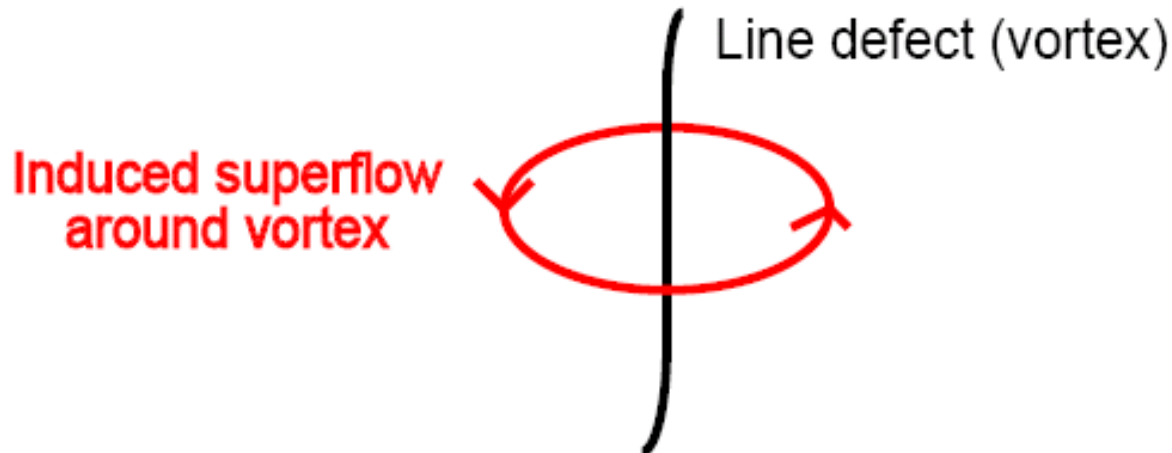


# Quantized Vortices

- Lowest energy state:  $n=1$ , so  $\phi$  wraps  $2\pi$  around a defect
- Induces a superflow around the line:

$$\mathbf{v}_{\Phi} = \frac{\mathbf{K}}{r}$$

$s$  is distance from defect



# What is quantum turbulence?

An evolving set of quantized vortices:

Aperiodic

Large range of length scales  
and curvatures

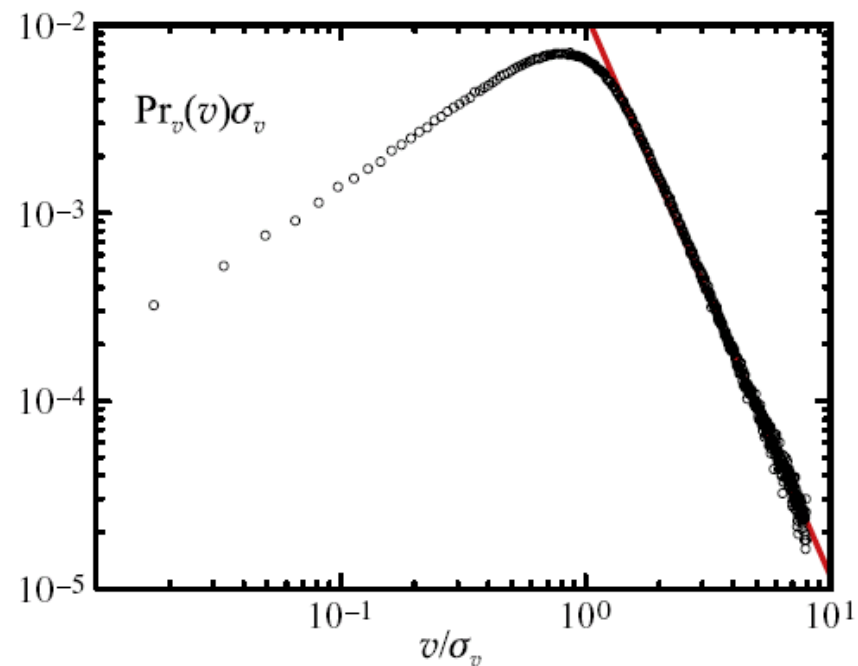
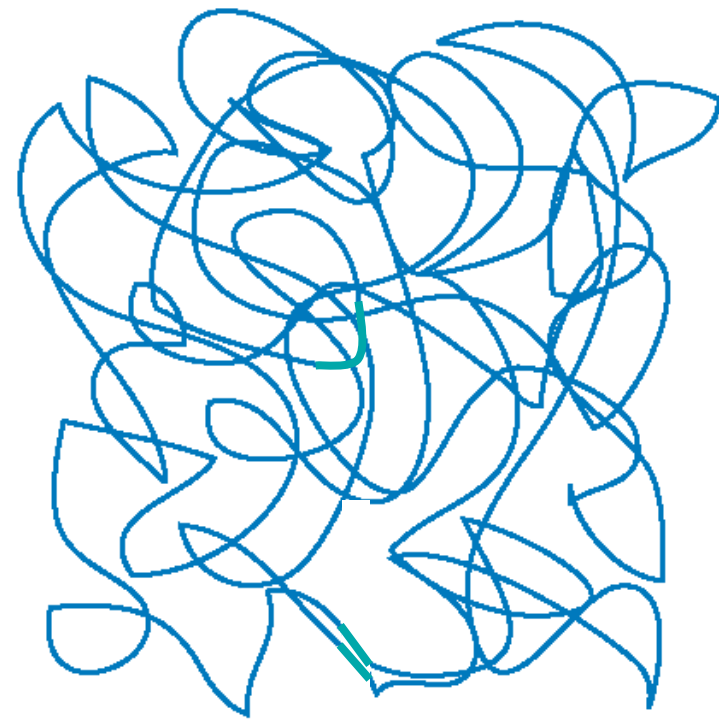
Rings

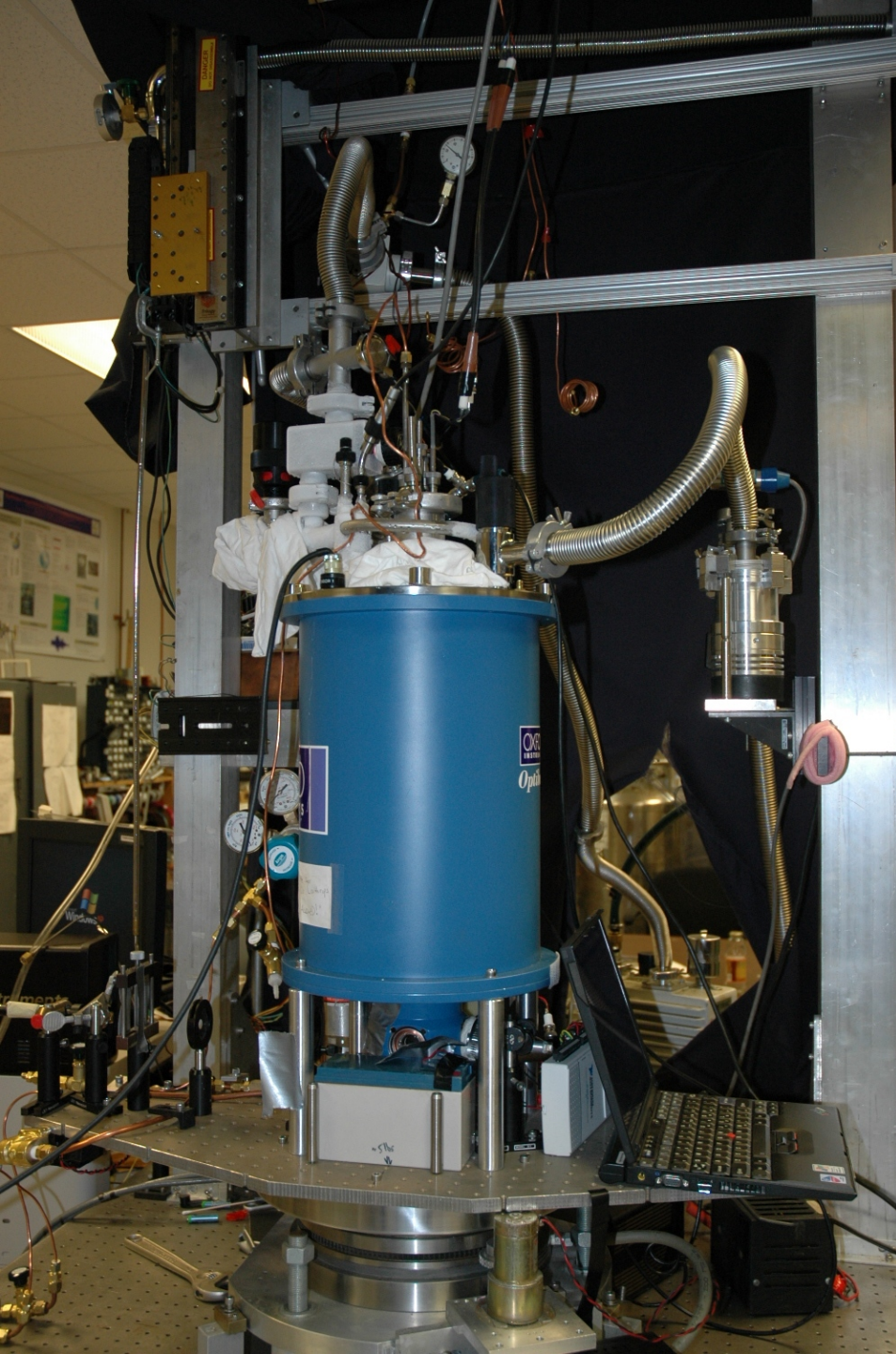
Vortices ending at walls

Knots

Quantum turbulence  
is dominated by:

Reconnection  
Ring collapse





# Superfluid helium Visualization

## Apparatus

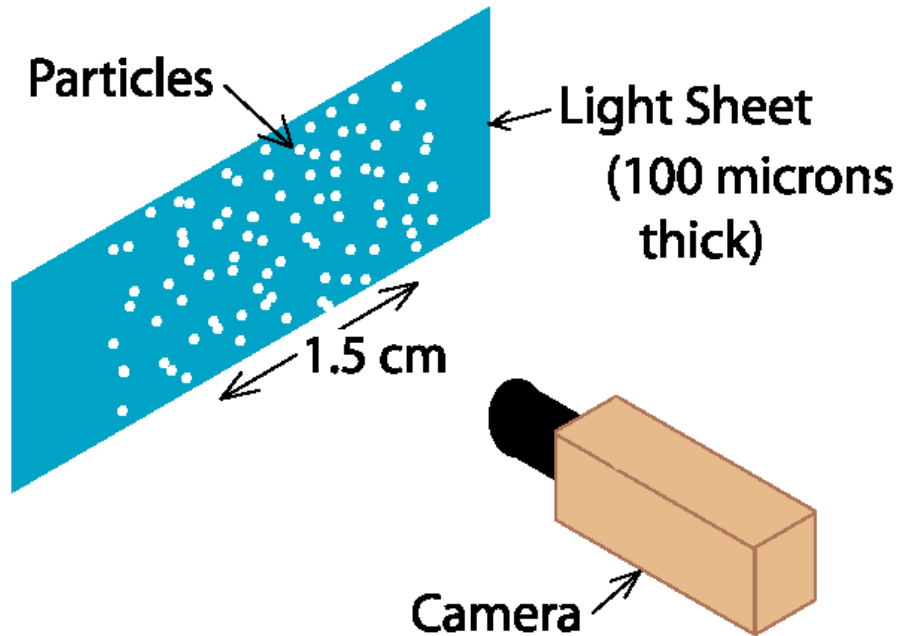
Optical cryostat

Particle injector

Laser sheet

Low light camera

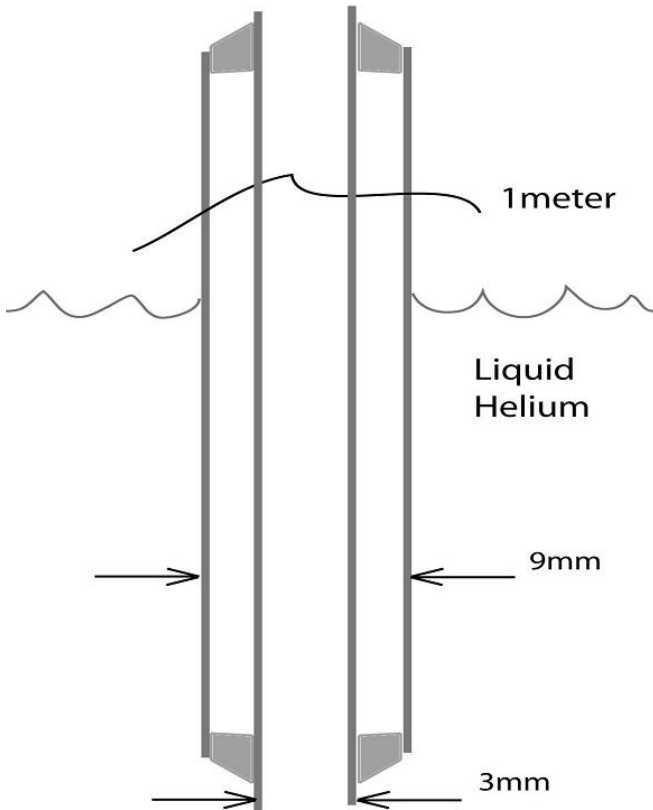
$T > 1.6$  K



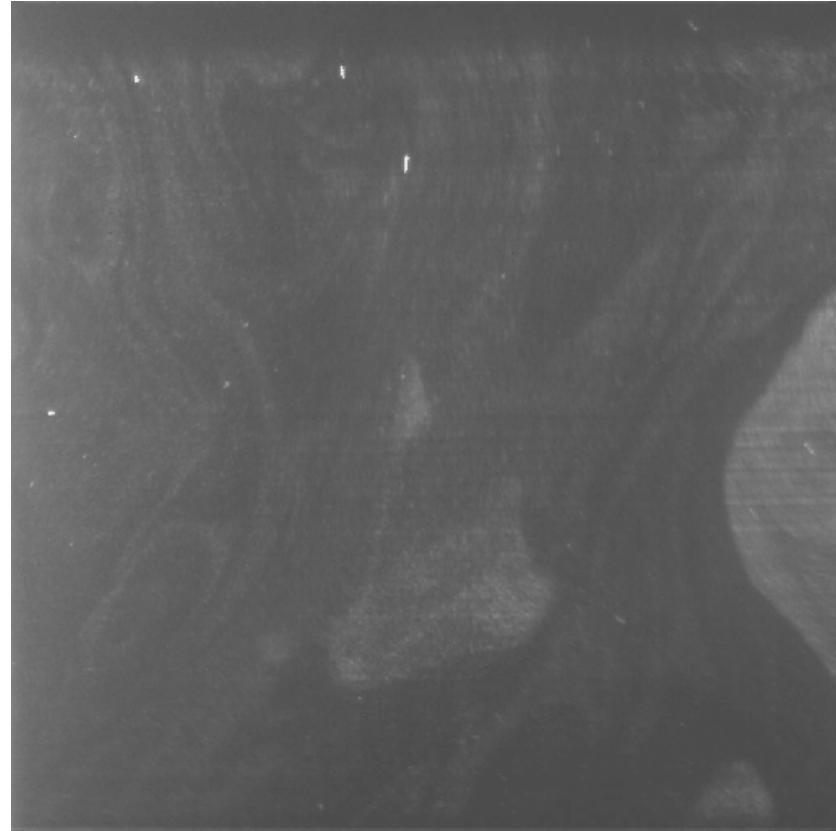
# Particle Production

$1 \text{ H}_2 : \chi \text{ He } \chi \gg 1$

↓ mixture pressure applied here



$T > T_\lambda$



↑  
1/2 image  
8 mm  
↓

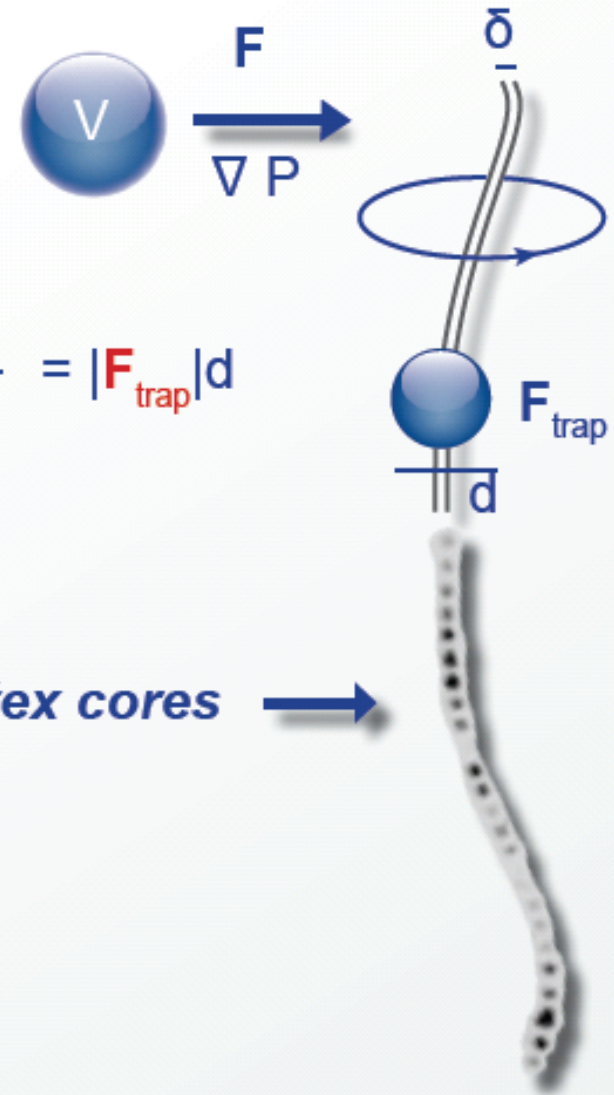
Bewley et al., *Experiments in Fluids* 2008



# Visualization of vortices - particle trapping

$$P = -\frac{\rho_s \kappa^2}{8\pi^2 r^2}$$

$$\mathbf{F} = \oint_{\partial\Omega} P \hat{n} dA$$



Decrease of energy

$$\Delta\varepsilon = \frac{\rho_s \kappa^2}{4\pi} d \ln \frac{d}{2\delta} = |\mathbf{F}_{\text{trap}}| d$$

$$\frac{\mathbf{F}_{\text{trap}}}{V} \propto \frac{\ln d}{d^3}$$

**Particles get trapped on vortex cores** →

Ions in liquid He     $\ominus$  16 Å     $\oplus$  6 Å

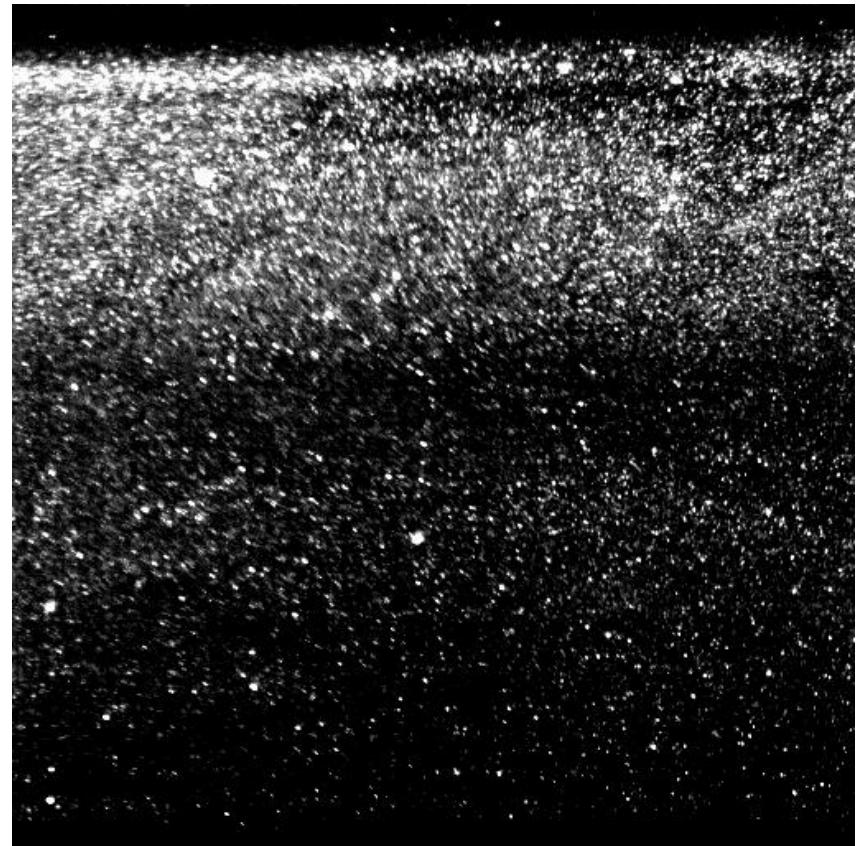
*Parks and Donnelly (1966), Williams and Packard (1974)*

Solid hydrogen particles > 1 μm

*Bewley, Lathrop, Sreenivasan, Nature 441 588 (2006)*

# Visualizing Superfluid Vortices in He II

- Below  $T_\lambda$  hydrogen particles collect onto filaments
- Previous work has shown these filaments are particles trapped on the superfluid vortices (Bewley, *et al.*, *Nature* 2006)

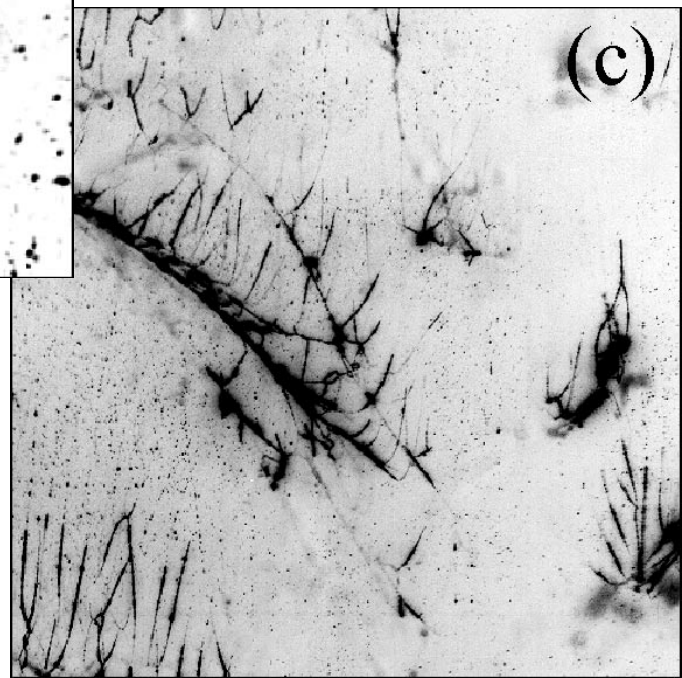
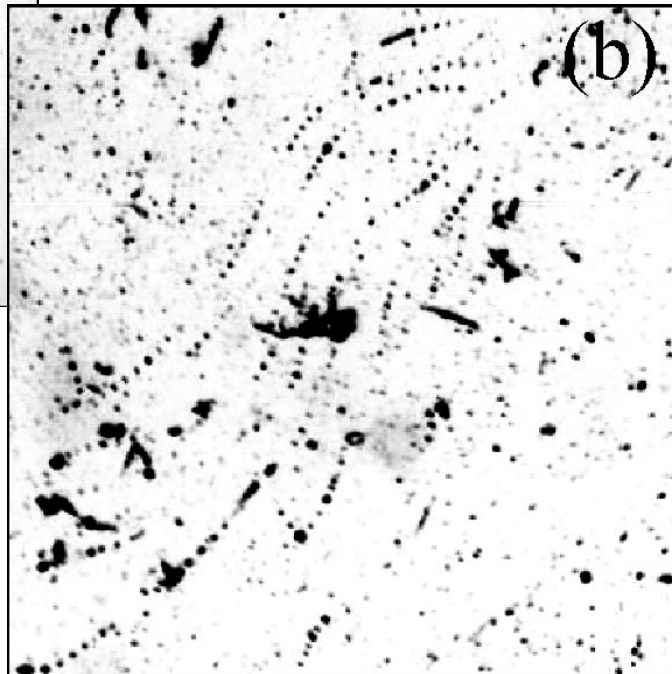
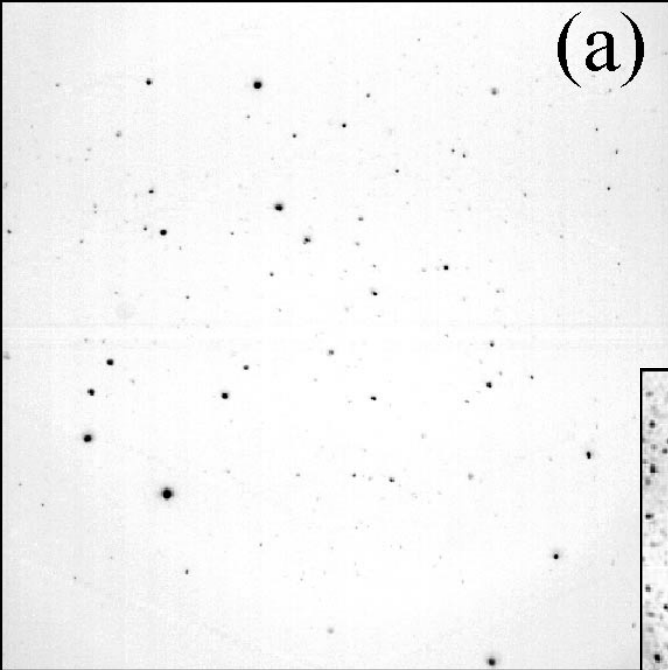


8  
mm

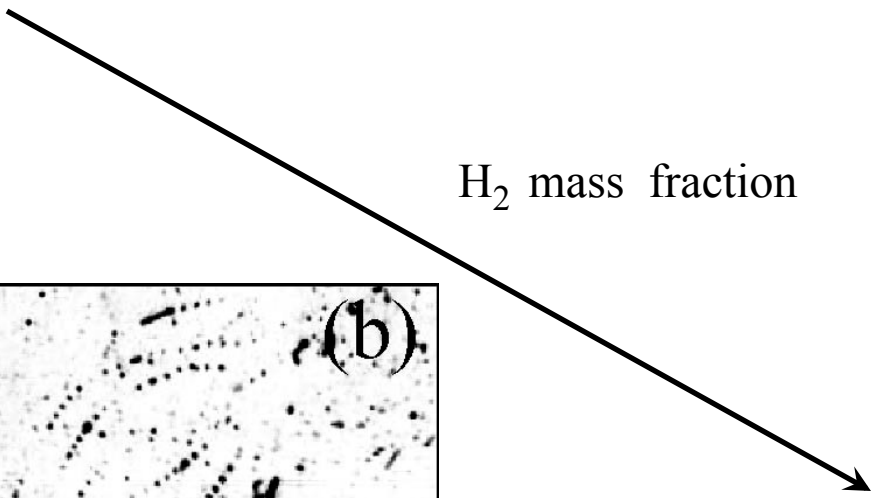
Movie in real time  
Begins 180 s after transition  
 $T_\lambda - T \sim 50$  mK

Sounds through transition caught with  
MEMS microphone





H<sub>2</sub> mass fraction



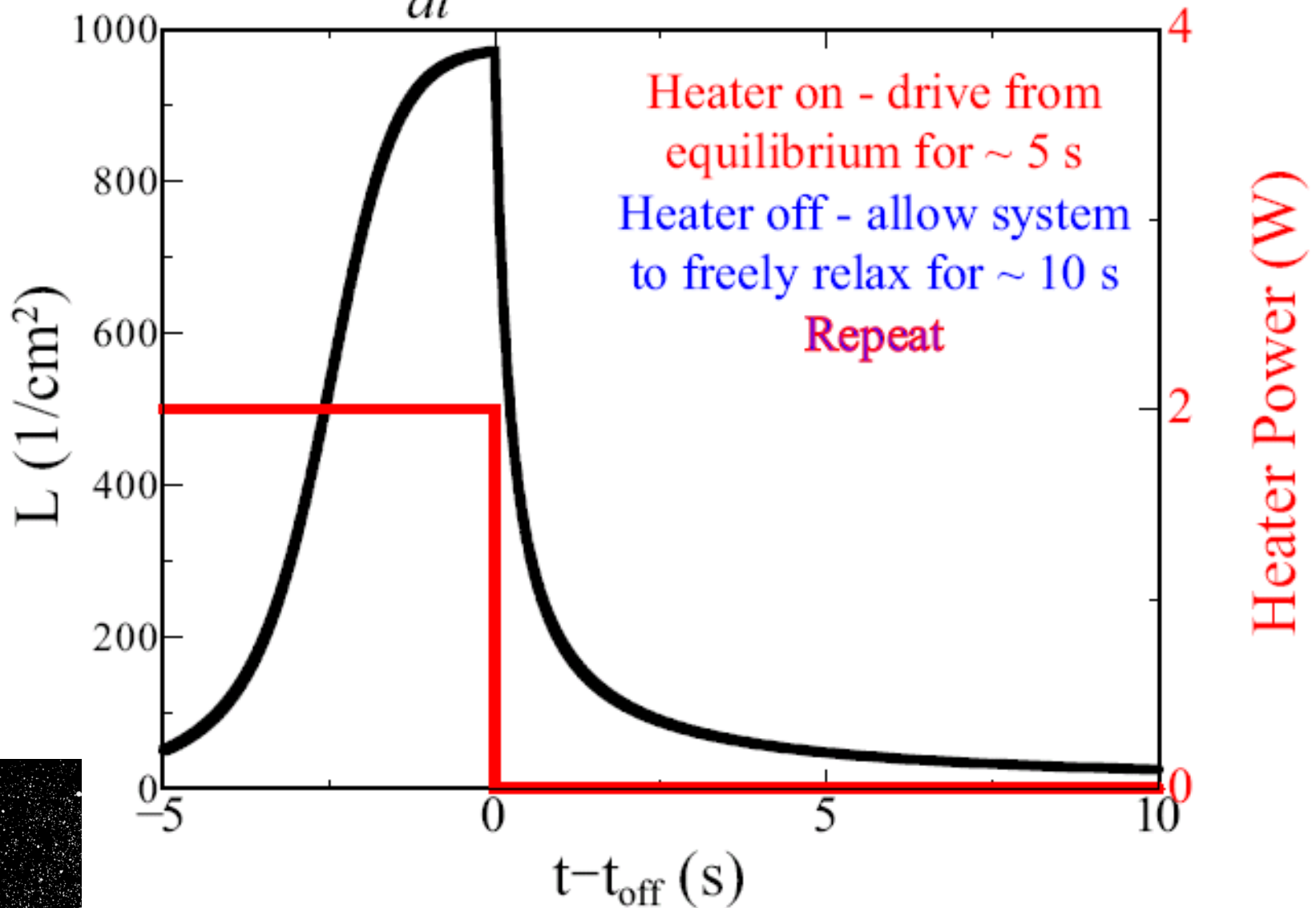
Smaller is better

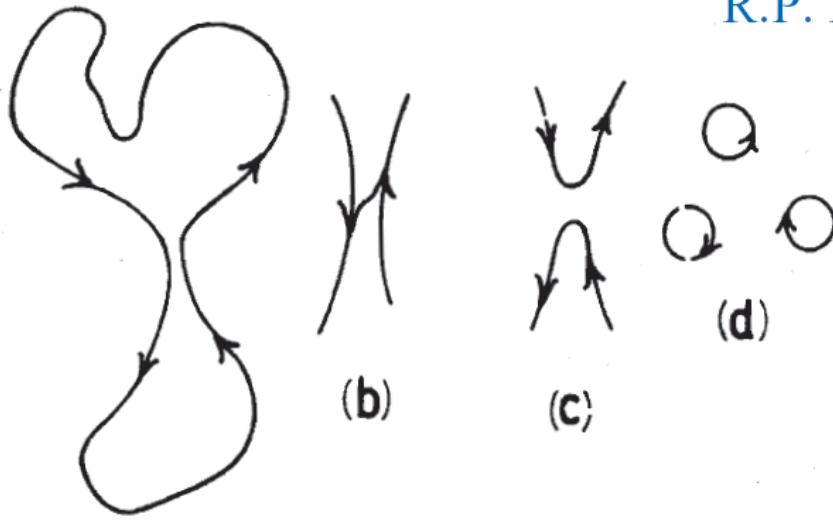
Fewer is better

Particles are not passive!

# Pulsed Counterflow Experiments

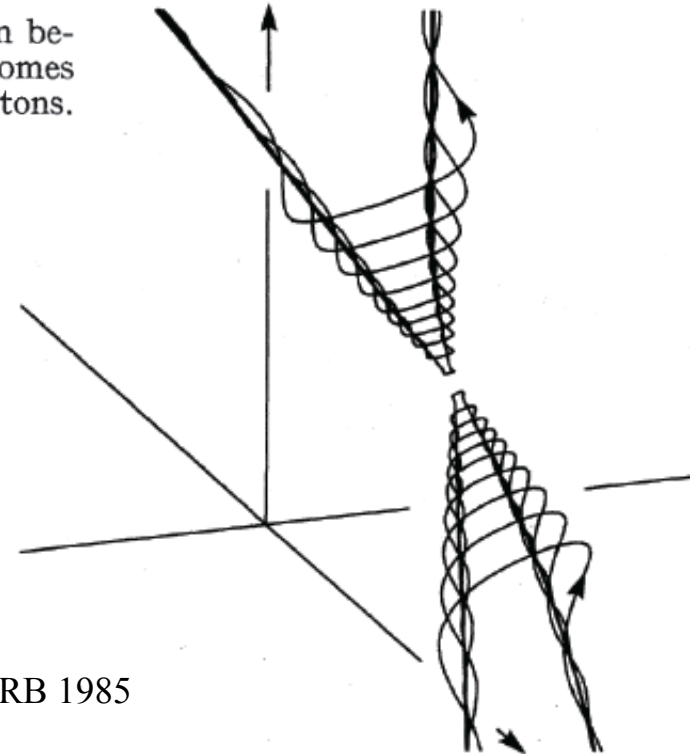
$$\frac{dL}{dt} = \alpha |\mathbf{v}_{ns}| L^{3/2} - \beta \kappa L^2$$





(a) Prog. Low Temp. Phys. 1, 17 (1955)

Fig. 10. A vortex ring (a) can break up into smaller rings if the transition between states (b) and (c) is allowed when the separation of vortex lines becomes of atomic dimensions. The eventual small rings (d) may be identical to rotons.



# Vortex reconnection

Theoretical work

Schwarz, PRB 1985 (LV)

de Waele and Aarts, PRL 1994 (LV)

Koplik and Levine, PRL 1993 (NLSE)

Tsubota and Maekawa, JPSJ 1992 (LV)

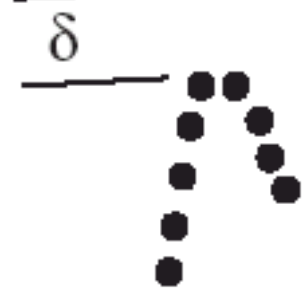
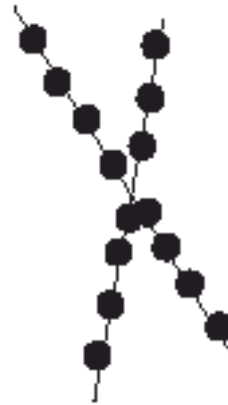
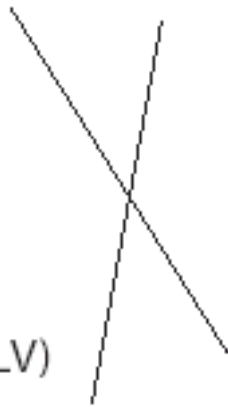
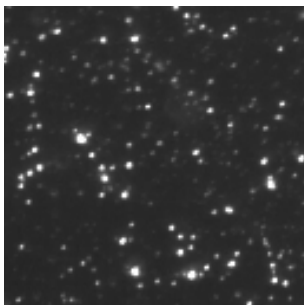
Nazarenko and West 2003 (NLSE)

Much more recent work!

Kerr PRL 2011 (NLSE)

$$\delta \sim \kappa^{1/2}(t_0 - t)^{1/2}$$

$$\delta \sim \kappa^{1/2}(t - t_0)^{1/2}$$



Pre-reconnection:  $\delta(t) = A[\kappa(t_0-t)]^{1/2}[1+c(t_0-t)]$

Post-reconnection:  $\delta(t) = A[\kappa(t-t_0)]^{1/2}[1+c(t-t_0)]$

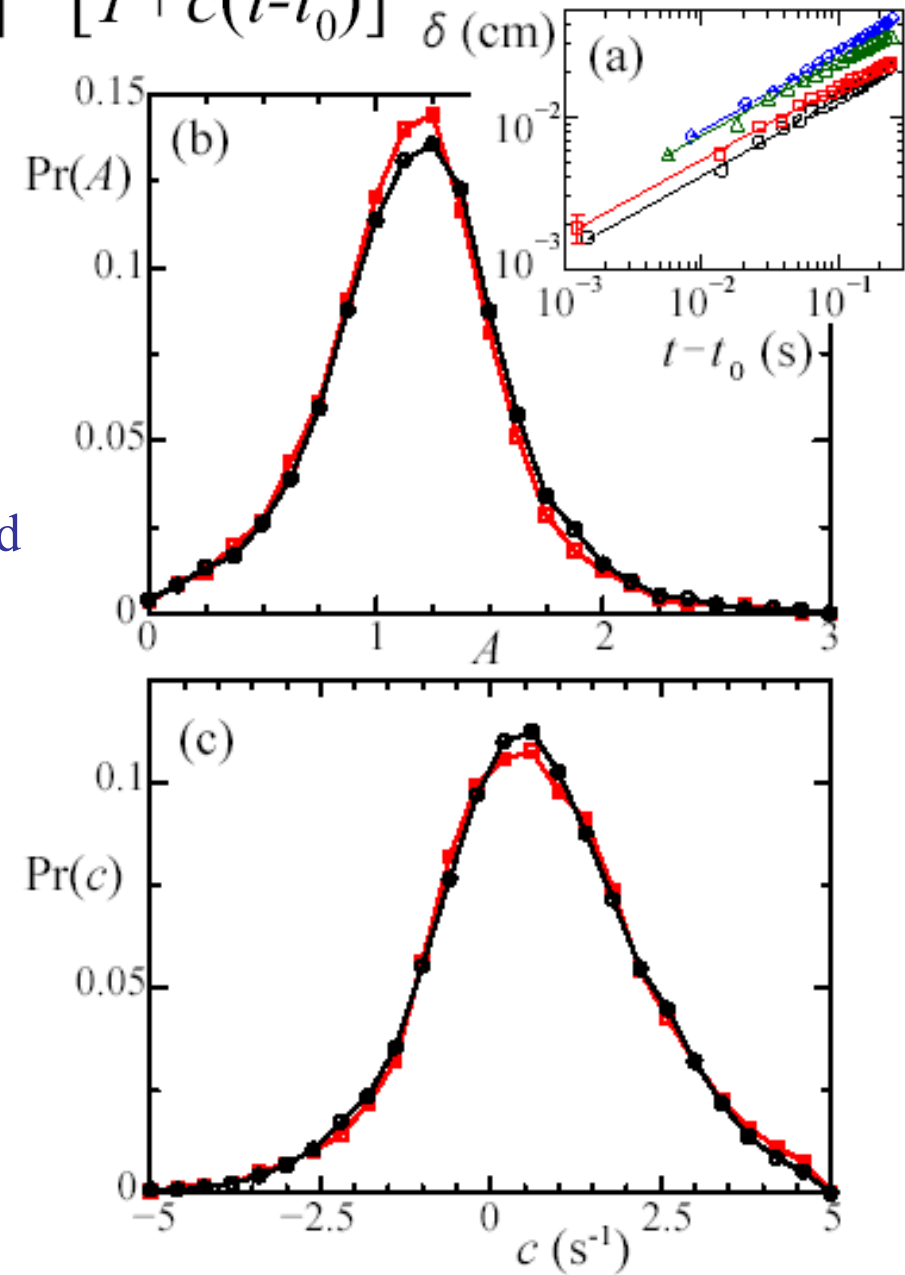
Only small pre- and post- differences

$c$  may represent the affect of local strains

reconnection and ring collapse represented

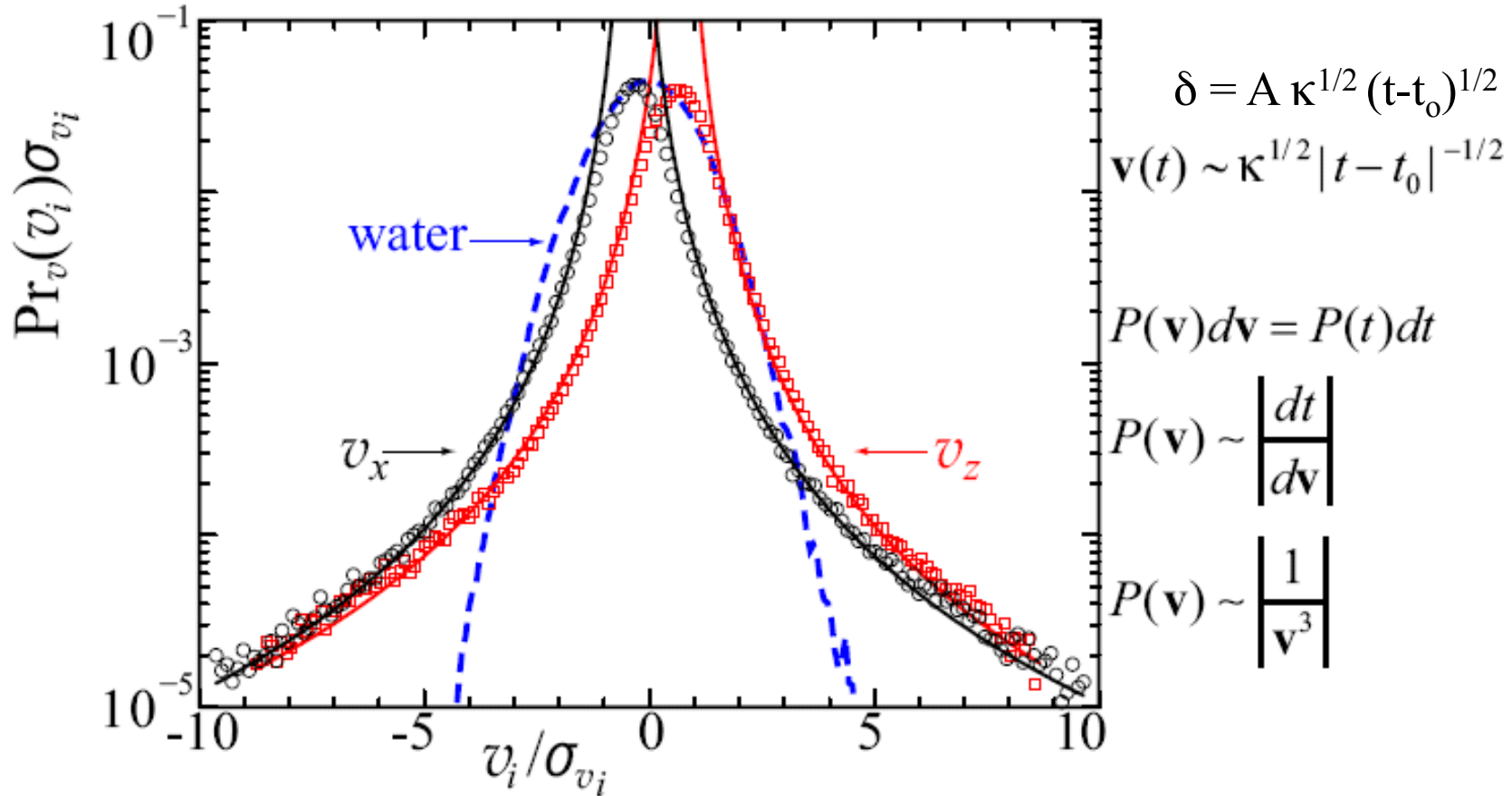
**NEARLY TIME REVERSIBLE!**

M.S. Paoletti, M.E. Fisher, and D.P. Lathrop,  
 "Reconnection dynamics for quantized vortices,"  
 Physica D (2010)





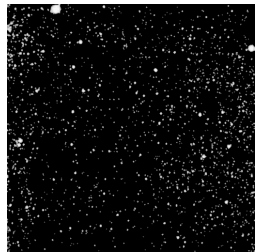
# Velocity Statistics

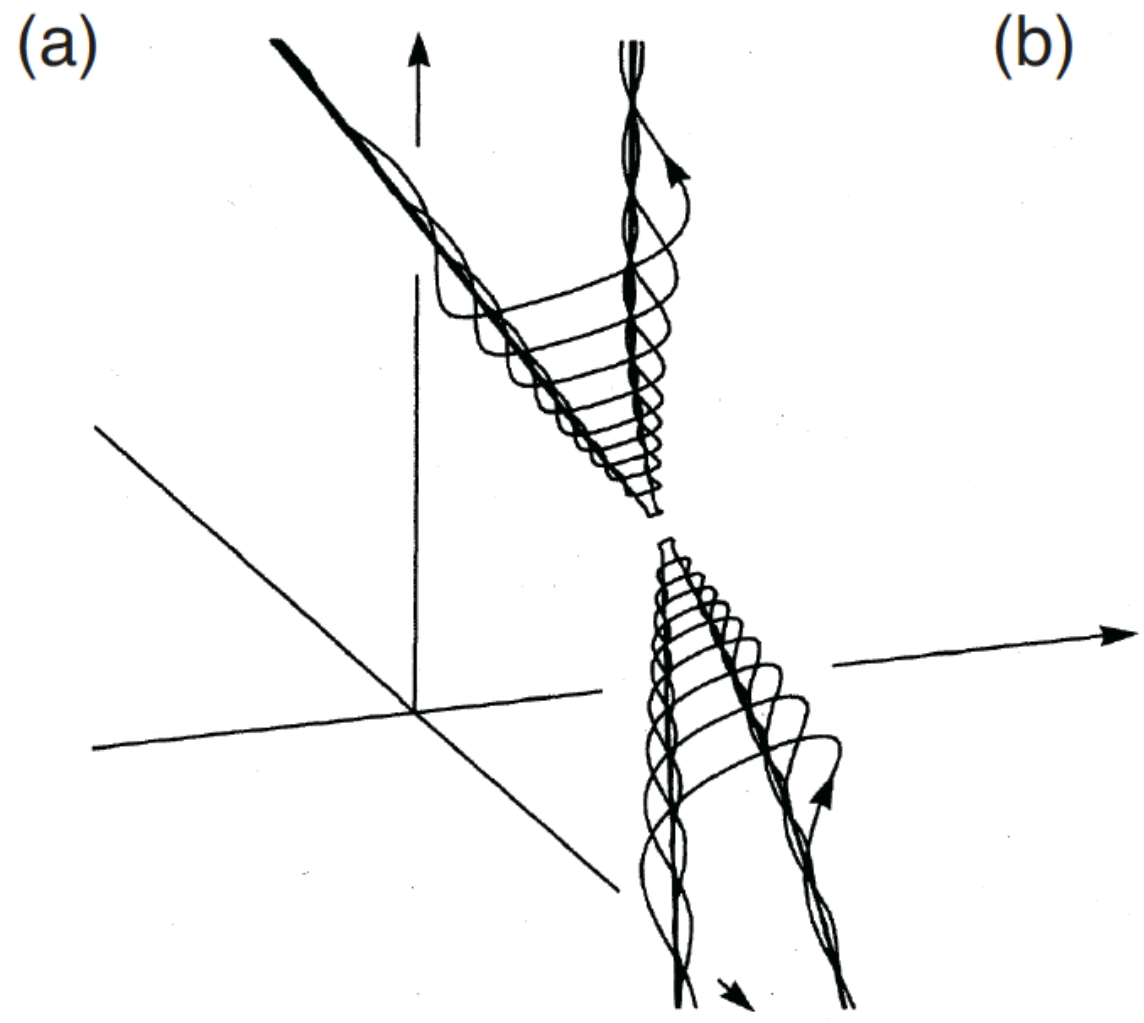
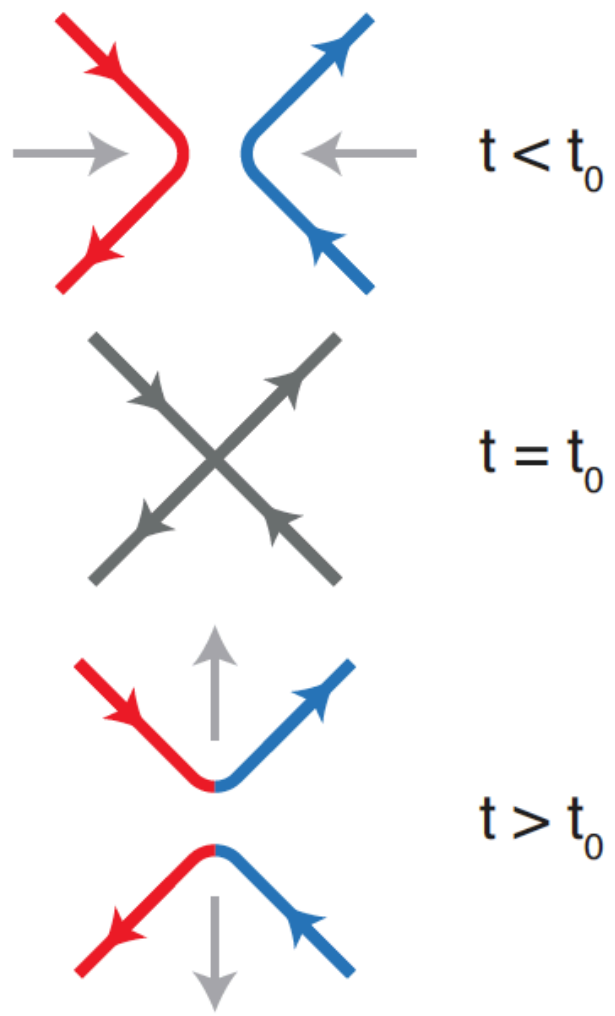


Reconnection produces predictable  
power-law velocity tails quite distinct  
from classical turbulence

M.S. Paoletti, M.E. Fisher, K.R. Sreenivasan, and D.P. Lathrop, "Velocity statistics distinguish quantum from classical turbulence," Phys. Rev. Lett. (2008)

Bagaley and Barenghi, PRE (2011).





# Vortex Filament Models

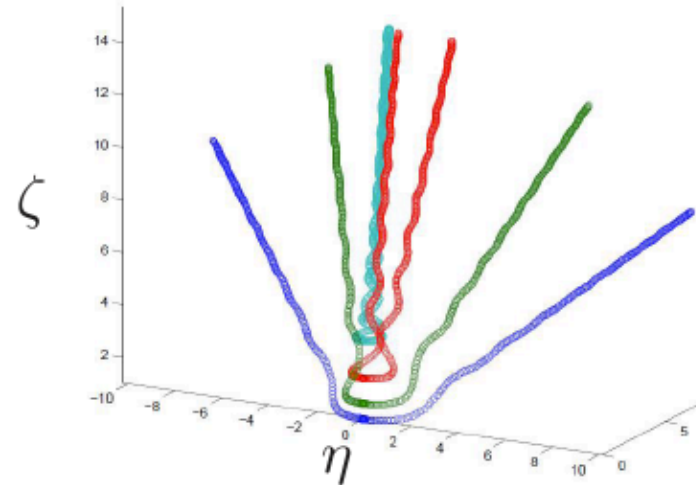
- Local Induction Approximation (LIA)

$$\frac{\partial \vec{s}(\sigma, t)}{\partial t} = \beta \frac{\partial \vec{s}(\sigma, t)}{\partial \sigma} \times \frac{\partial^2 \vec{s}(\sigma, t)}{\partial \sigma^2} + \alpha(T) \frac{\partial^2 \vec{s}(\sigma, t)}{\partial \sigma^2}$$

- LIA has one-parameter family of self-similar solutions in dimensionless similarity coordinates
- Adopt dimensionless similarity coordinates

$$\eta = (x - x_0) / \sqrt{\kappa(t - t_0)}$$
$$\zeta = (z - z_0) / \sqrt{\kappa(t - t_0)}$$

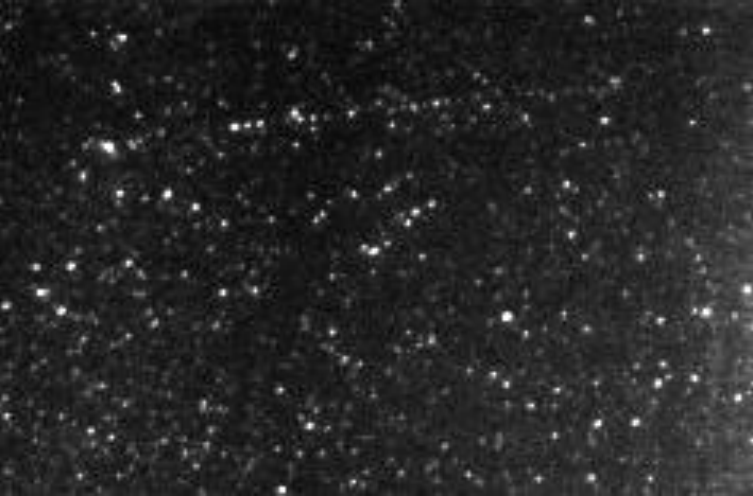
LIA Curves vs Vortex Angle



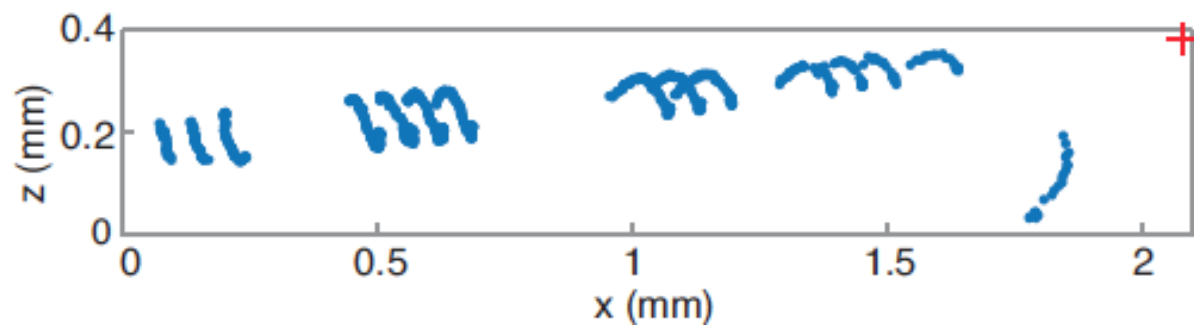
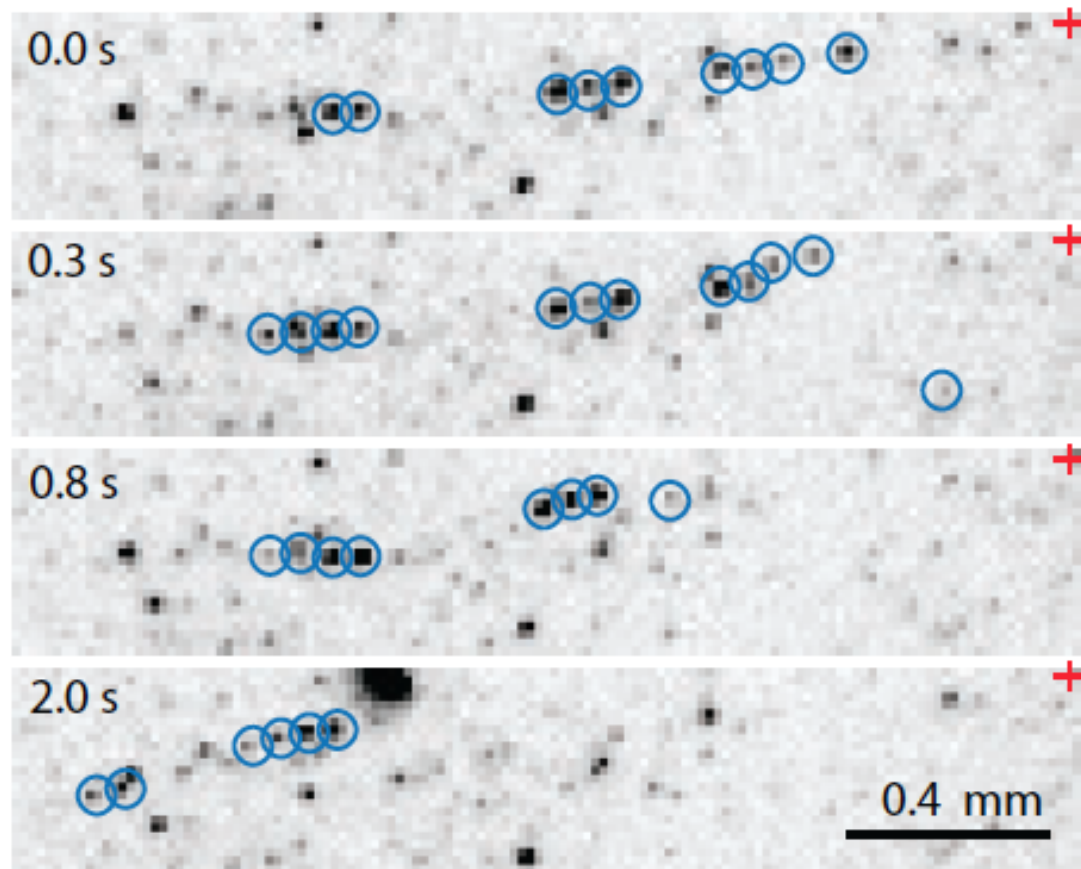
*Eur. J. Mech. B - Fluids* 19 (2000) 361-378

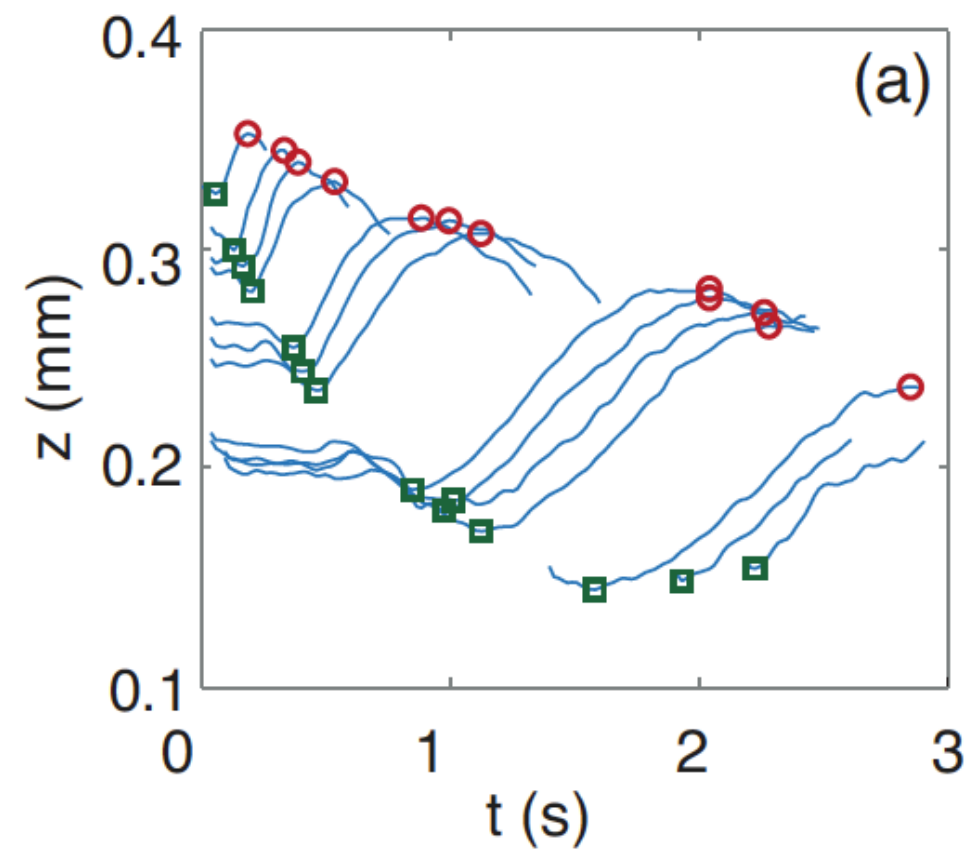
© 2000 Éditions scientifiques et médicales Elsevier SAS. All rights reserved  
S0997-7546(00)00123-0/FLA

**Evolution of quantum vortices following reconnection**

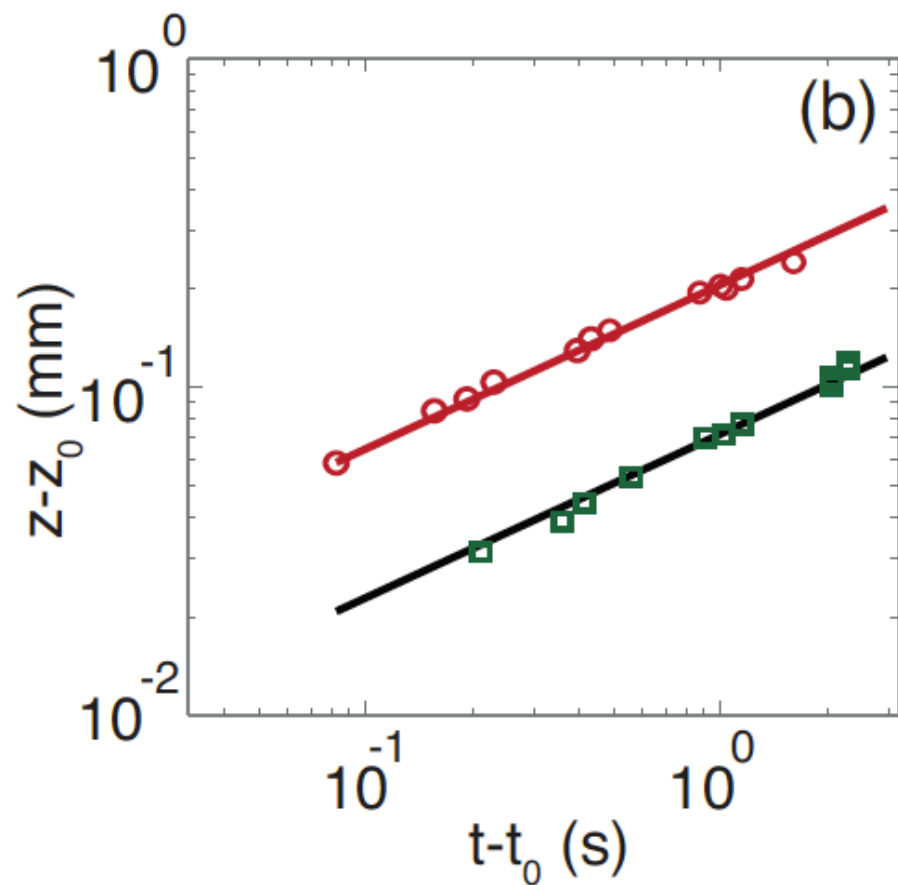


# Kelvin waves!

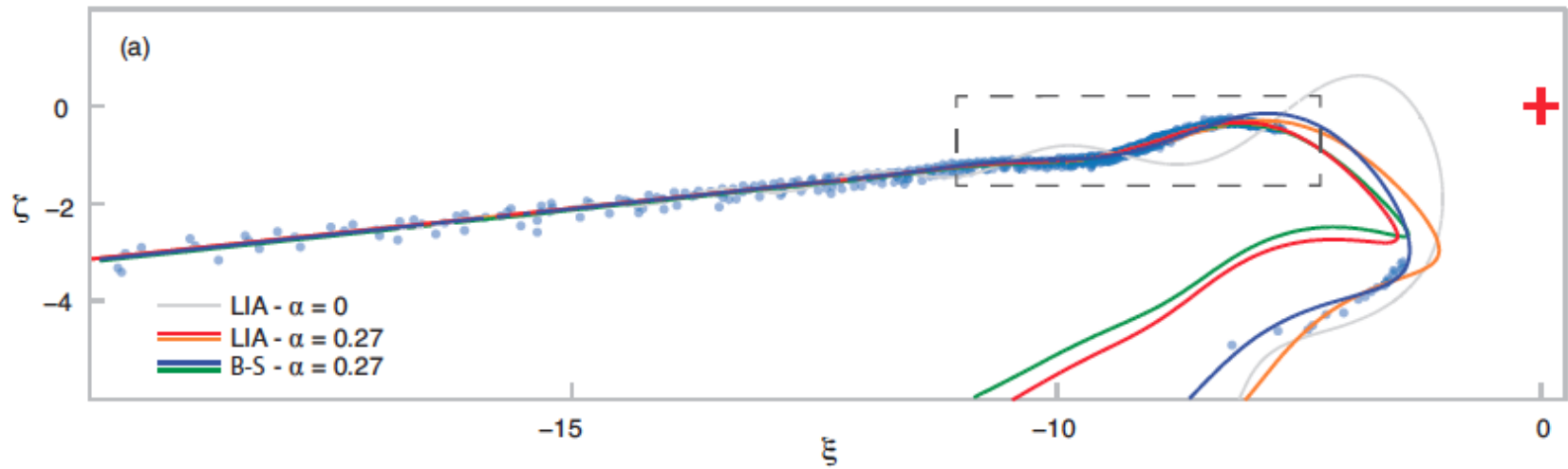
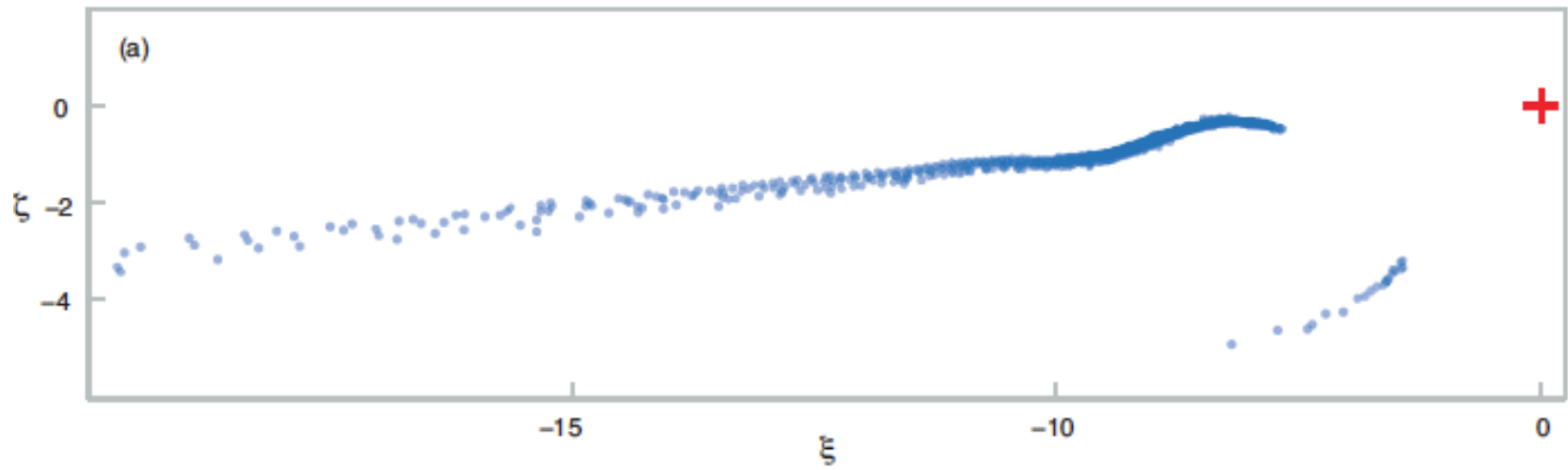




$$\xi = (x - x_0) / \sqrt{\kappa(t - t_0)}$$

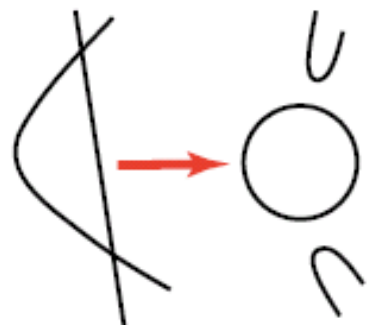
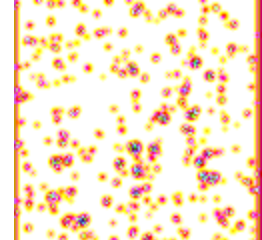


$$\zeta = (z - z_0) / \sqrt{\kappa(t - t_0)}$$

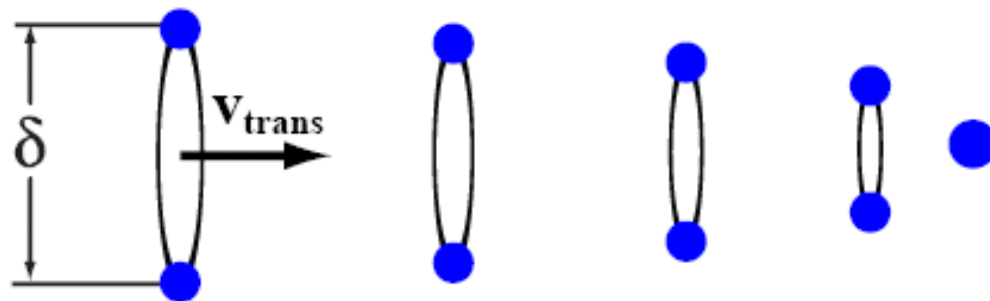


Collapse in similarity coordinates and a comparison with theoretical models

# Quantized Vortex Rings



Reconnection  
can produce  
vortex rings



Pair of particles on collapsing  
rings look like reconnection  
backwards in time with additional  
transverse velocity





# Fixed points of the nonlinear Schrodinger equation

$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \Psi + g |\Psi|^2 \Psi - \mu \Psi$$

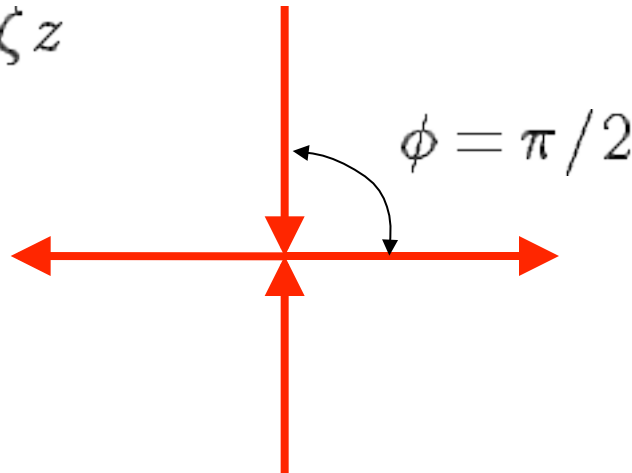
$$\Psi_0 = \sqrt{(\mu/g)}$$

$$\Psi_1 = f(r) e^{i\theta}$$



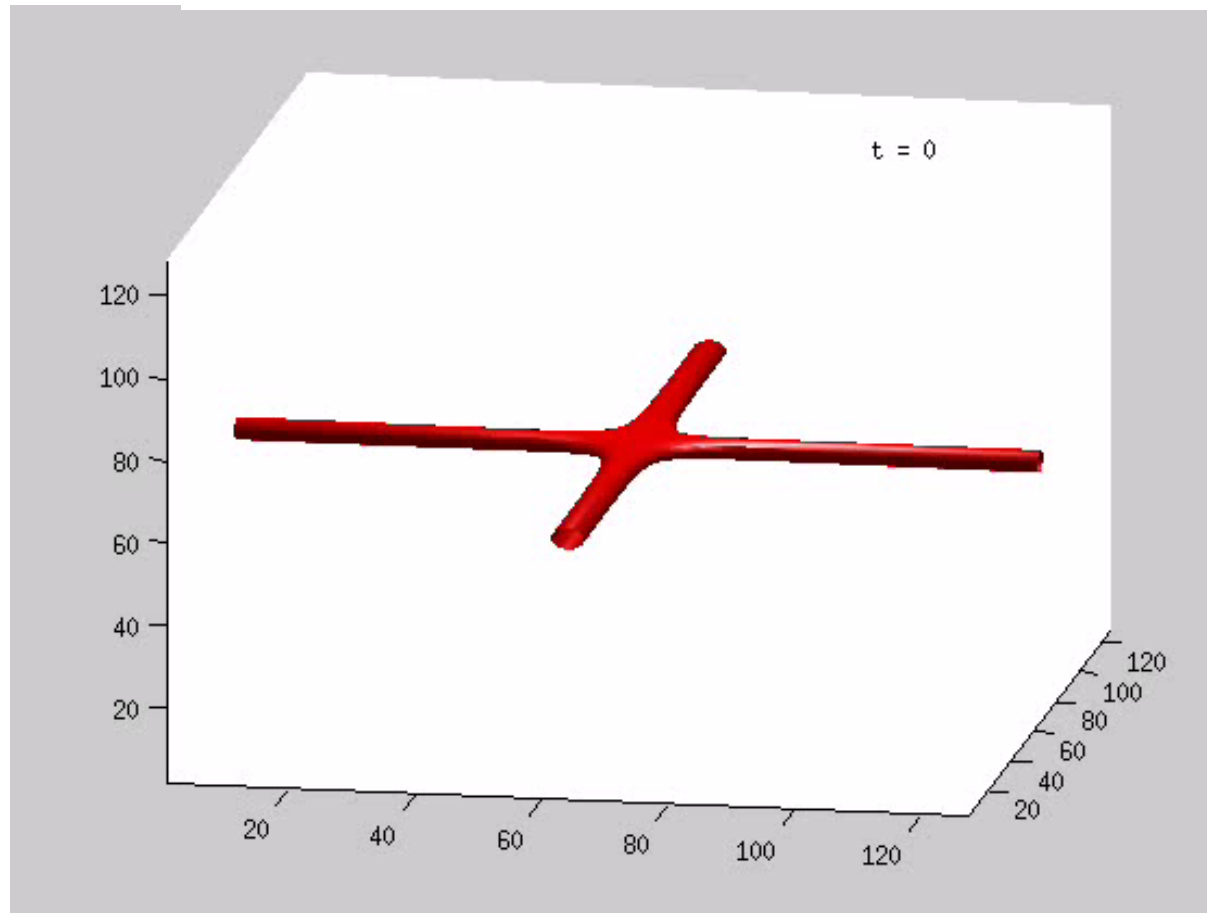
New fixed point. Near origin:

$$\Psi_2 = xy + i\zeta z$$



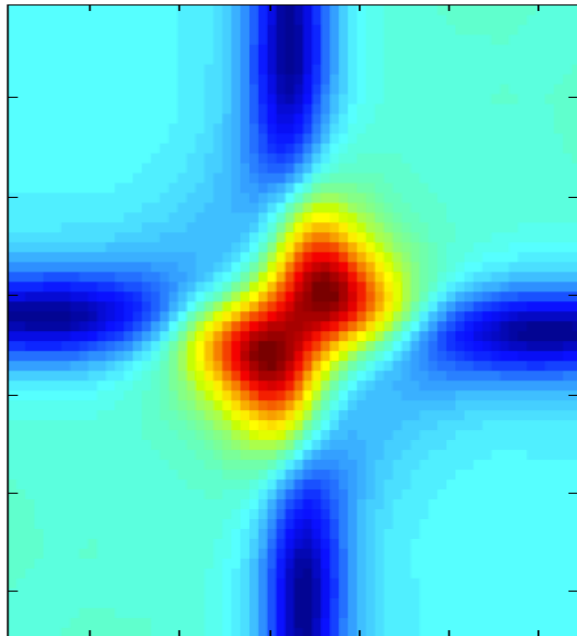
# New fixed point is a saddle

David P. Meichle, Cecilia Rorai, Michael E. Fisher, and D. P. Lathrop  
Phys. Rev. B **86**, 014509 – Published 10 July 2012

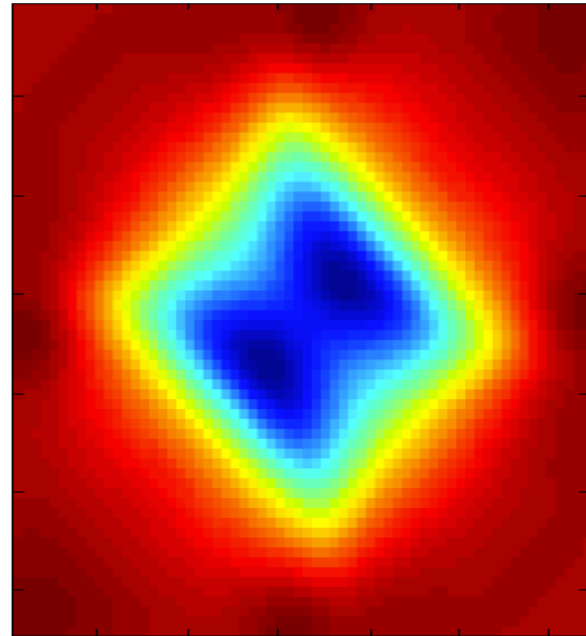


Eigenvector for saddle directions away from  $\Psi_2$

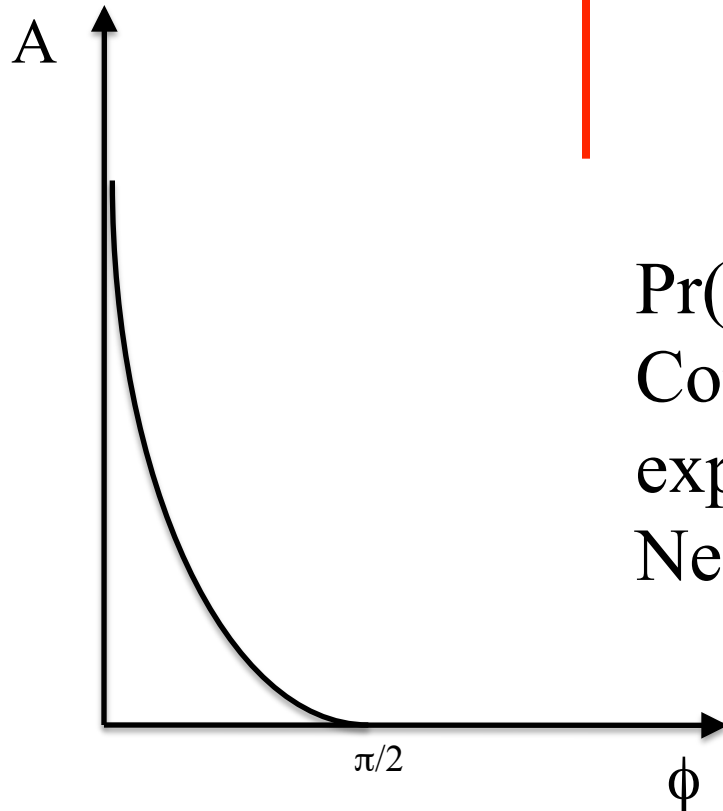
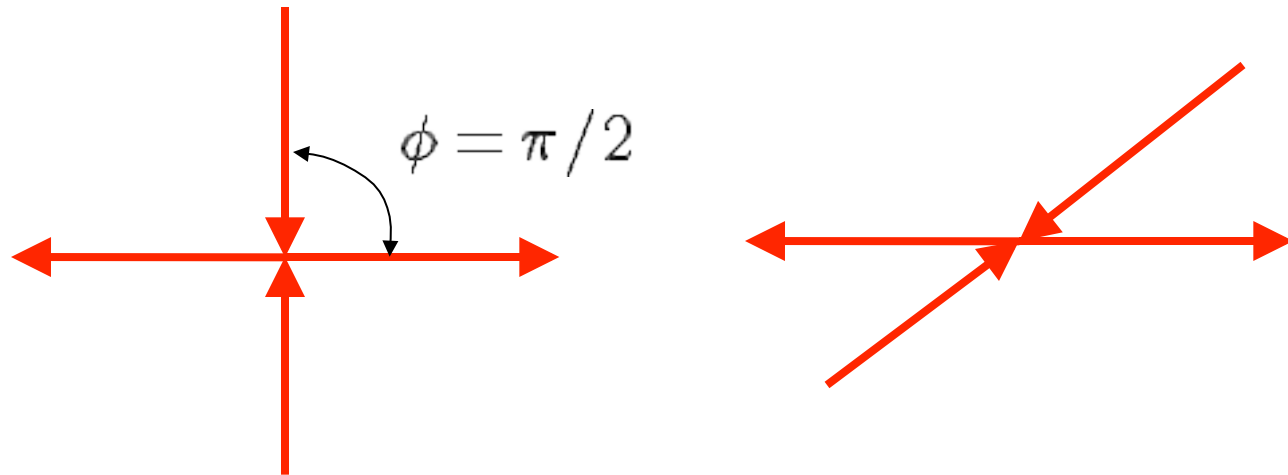
Real part



Imaginary part



# Prefactor depends on reconnection geometry

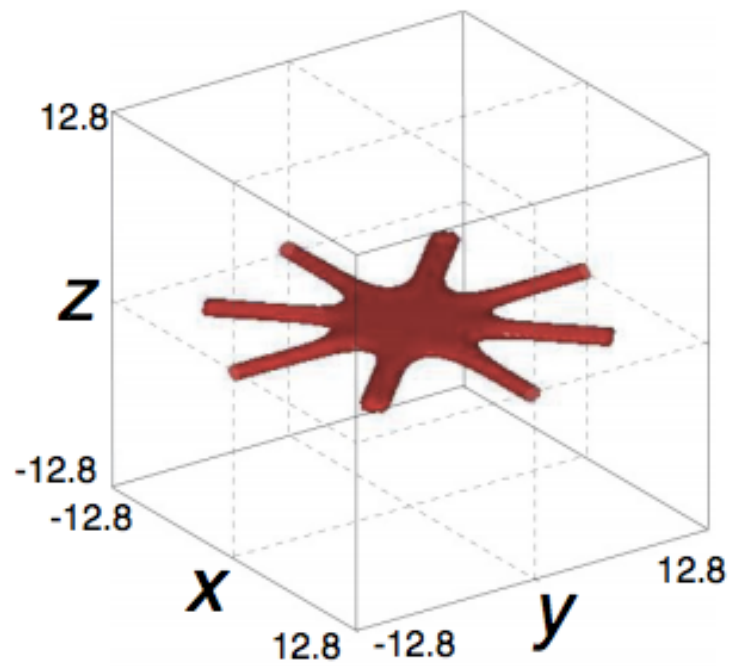


$$\Pr(A) = \Pr(\phi) |d\phi/dA|$$

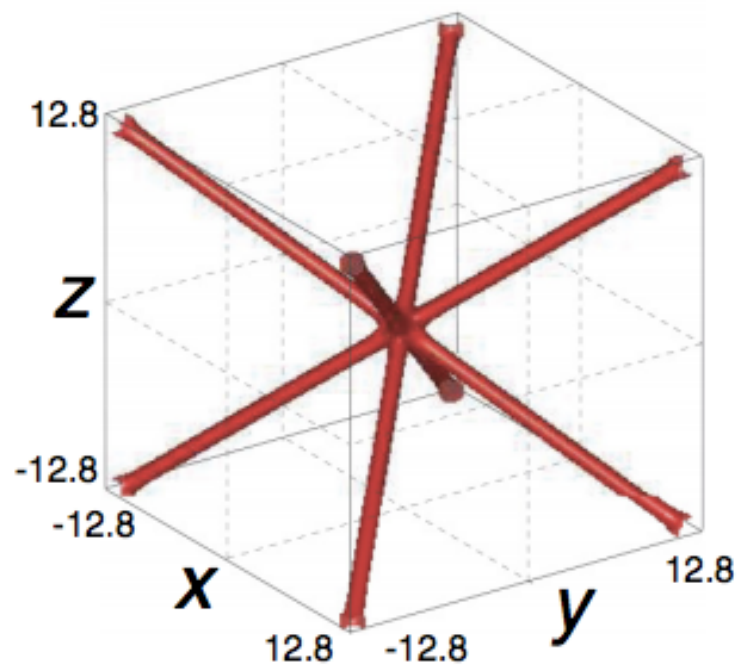
Connection possible between  
experiment and theory

Need  $A(\phi)$ !

(a)



(b)



Dissertations: [complex.umd.edu](http://complex.umd.edu)

Youtube channel: n3umh

Bewley, Lathrop, and Sreenivasan Nature 2006

Bewley, Sreenivasan, and Lathrop, Exp. in Fluids 2008

Paoletti, Fiorito, Sreenivasan, and Lathrop, J. Phys. Soc. Japan 2008

Bewley, Paoletti, Sreenivasan, and Lathrop, Proc. Nat. Acad. Sci. 2008

Paoletti, Fisher, Sreenivasan, and Lathrop, PRL 2008

Paoletti, Fisher, and Lathrop, Physica D 2008

Paoletti and Lathrop, Ann. Rev. of Cond. Matter Phys. 2011

Meichle, Rorai, Fisher, and Lathrop, PRB 2012

Fonda, Meichle, Ouellette, Hormoz, and Lathrop PNAS 2014

Meichle and Lathrop Rev. Sci. Inst. 2014

Next steps: 3-D tracking

# We love fluid turbulence

Velocity field rough  $\vec{v}(\vec{x},t)$

Large range of length scales

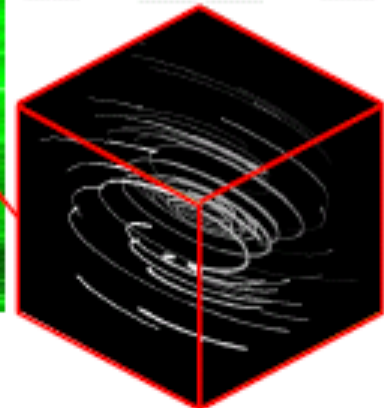
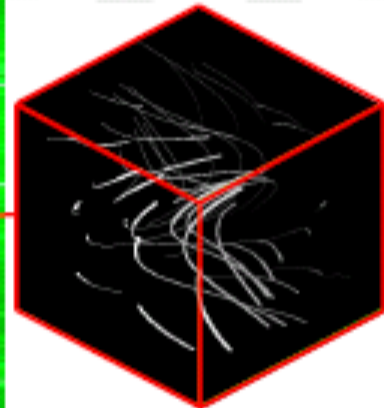
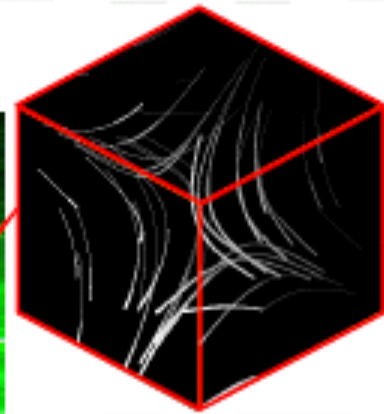
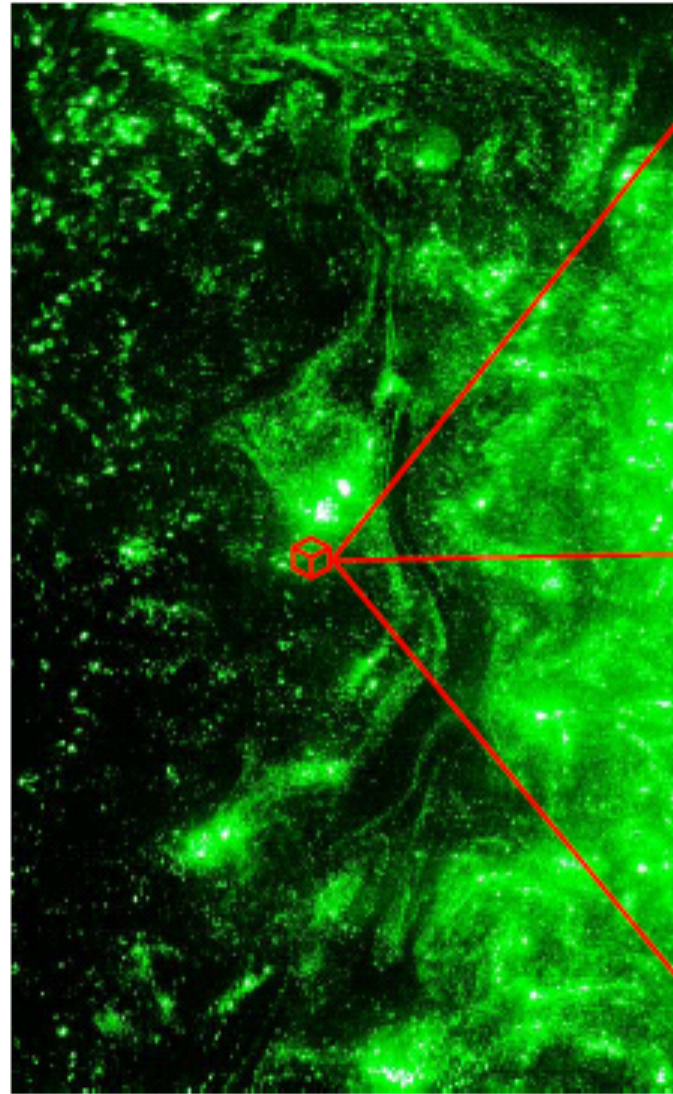
Large range of time scales

$\Pr(v_i)$  Gaussian

$E(k) \sim k^{-5/3}$  Inertial range

Known – agreed on equations  
of motion

Navier-Stokes Eq.





**Navier-Stokes flows: ? : turbulence**

$$Pr(v_i) \sim \exp(-v_i^2/\sigma_v^2)$$

$$Pr(\partial_j v_i) \sim \exp(-|v_i|/\sigma_s)$$

$Pr(\epsilon)$  dissipation interesting!

$Pr(\Omega)$  enstrophy interesting!

**Euler flows: ? : ?**

Fixed points associated with reconnection?

# New optical measurement of gradients and velocity at Kolmogorov scale

$$M = \partial u_i / \partial x_j = S + A$$

**Dissipation (strain)**

$$\varepsilon = \nu \|\mathbf{S}\|^2 / 2$$

$$\varepsilon(\vec{x}, t)$$

**Enstrophy (rotation)**

$$\Omega = \|\mathbf{A}\|^2 / 2$$

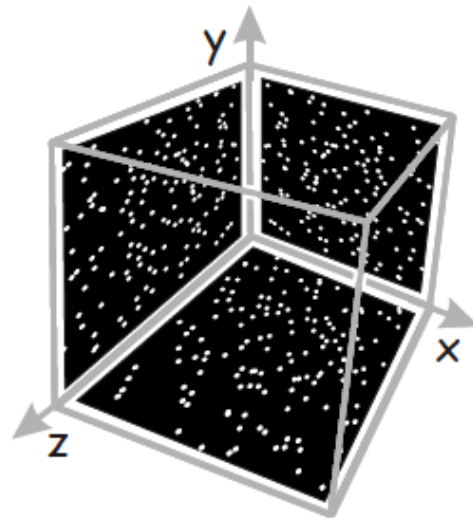
$$\Omega(\vec{x}, t)$$

Two interacting fields

What are the causal connections?

Oscillating grid turbulence at  $R_\lambda = 54$

# Measuring Gradients



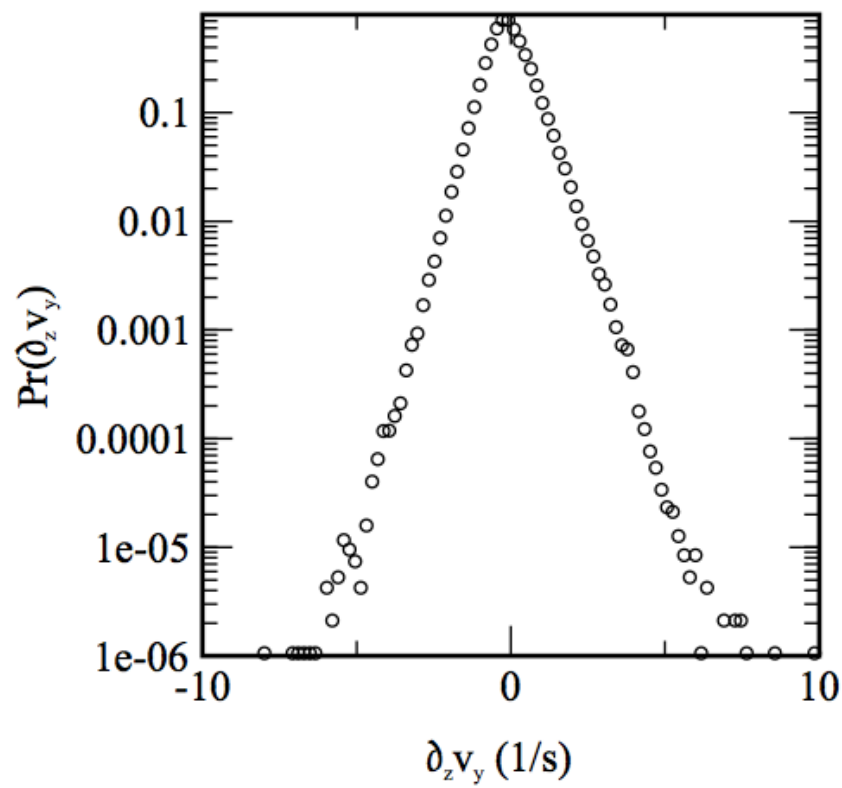
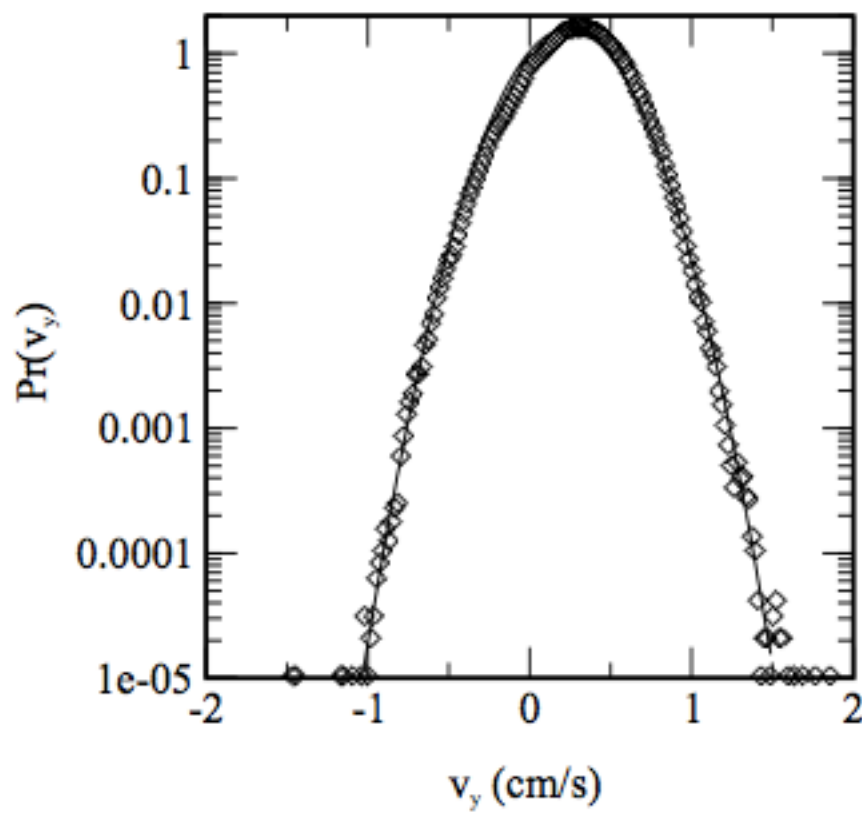
3 camera  
views



fit data to model

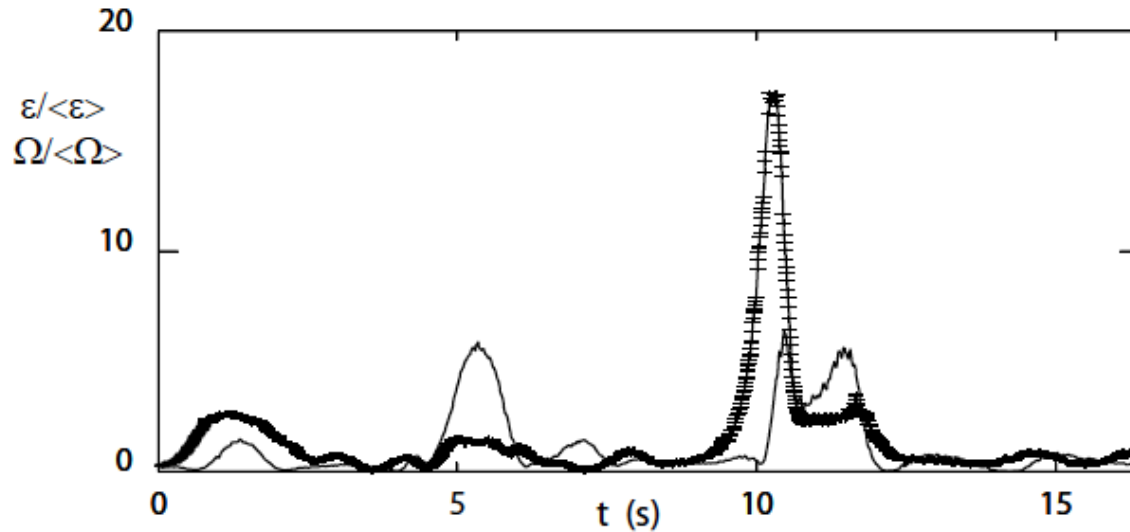
$$\vec{v} = \vec{v}_o + M \cdot \vec{x}$$

$$M = \begin{bmatrix} \partial_x V_x & \partial_y V_x & \partial_z V_x \\ \partial_x V_y & \partial_y V_y & \partial_z V_y \\ \partial_x V_z & \partial_y V_z & \partial_z V_z \end{bmatrix}$$



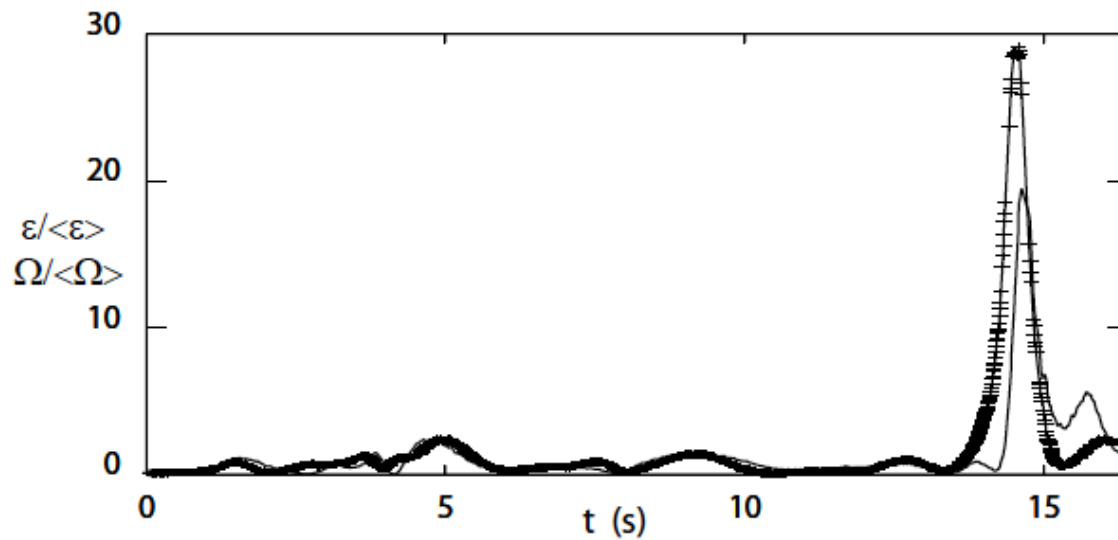
## Dissipation and Enstrophy Time Series – Intense Events

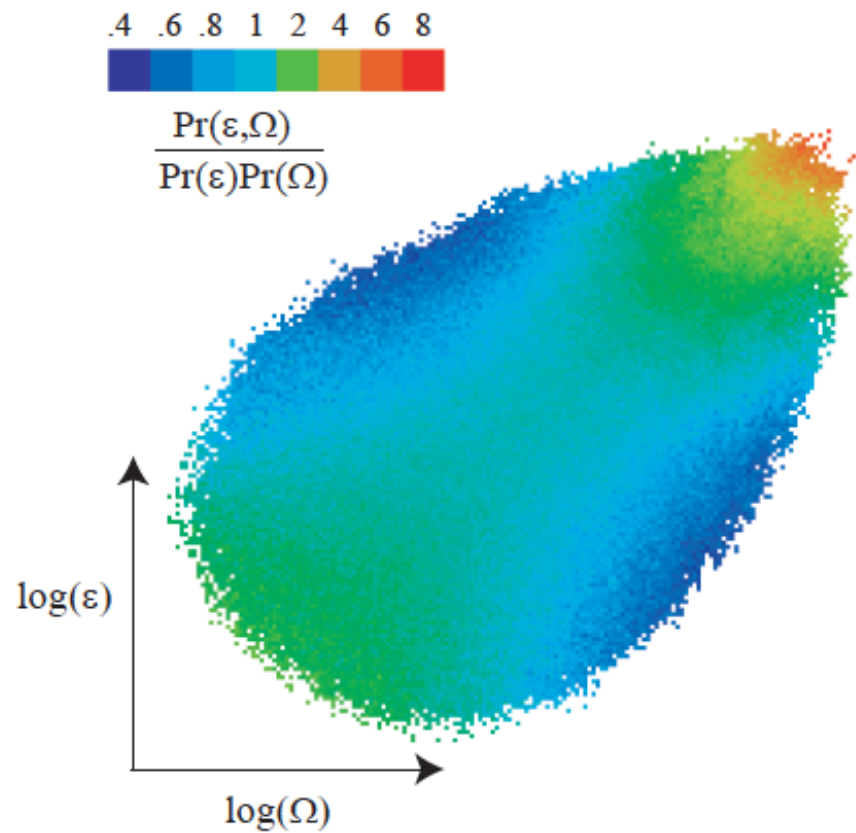
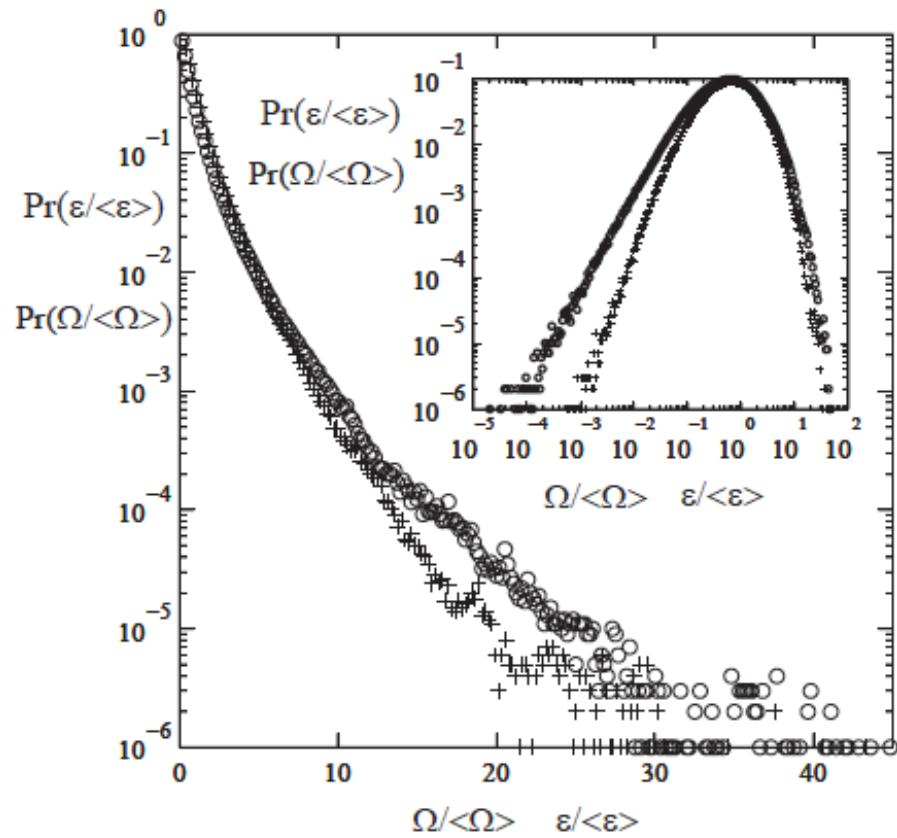
Zeff, Lanterman, McAllister, Roy, Kostelich, and Lathrop, Nature 421, 146 (2003)

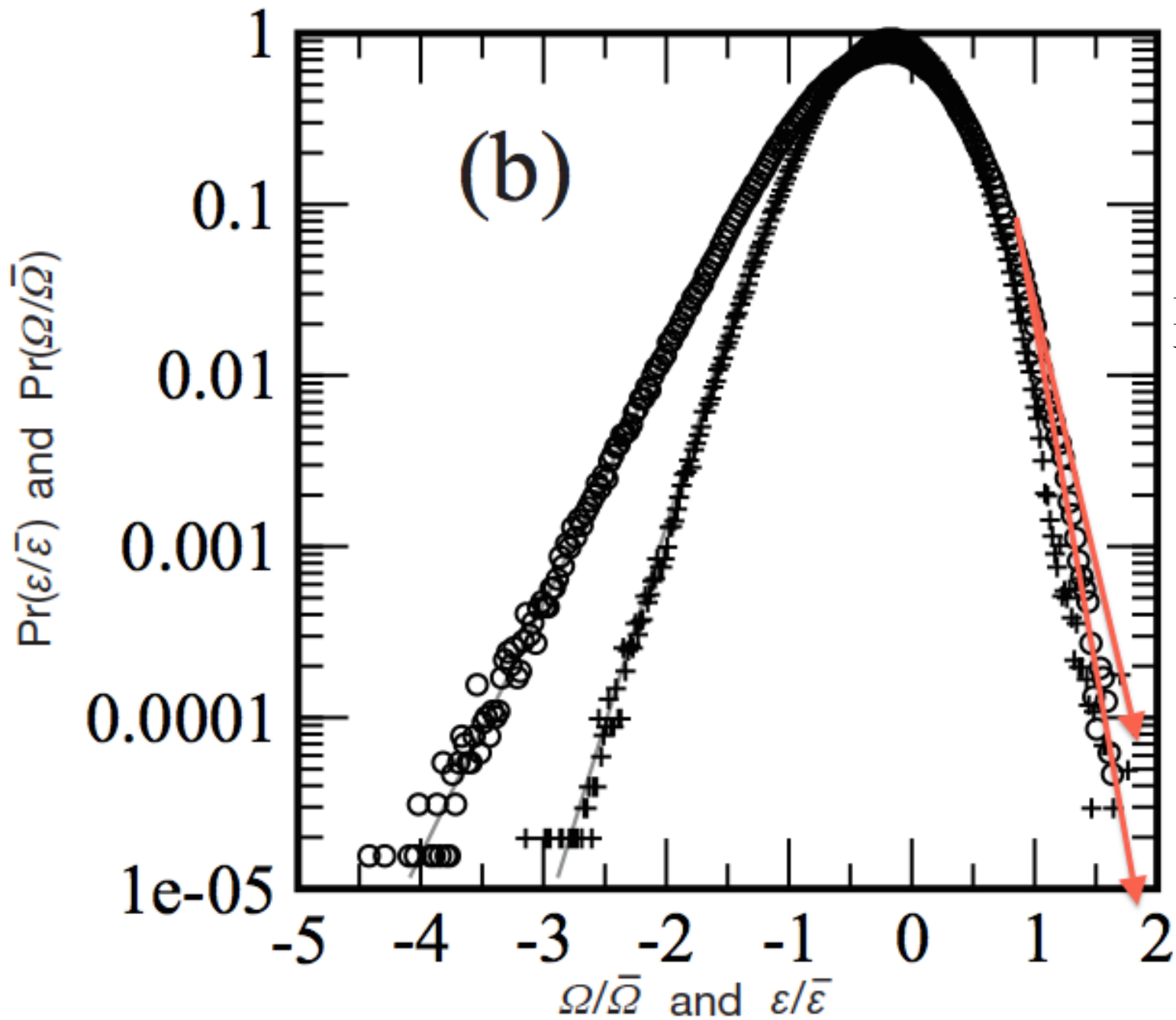


$$\varepsilon = \nu / 2 \parallel M_{ij} + M_{ji} \parallel_2$$

$$\Omega = 2 \parallel M_{ij} - M_{ji} \parallel_2 = \omega^2 / 2$$







Does the  
Variance exist?

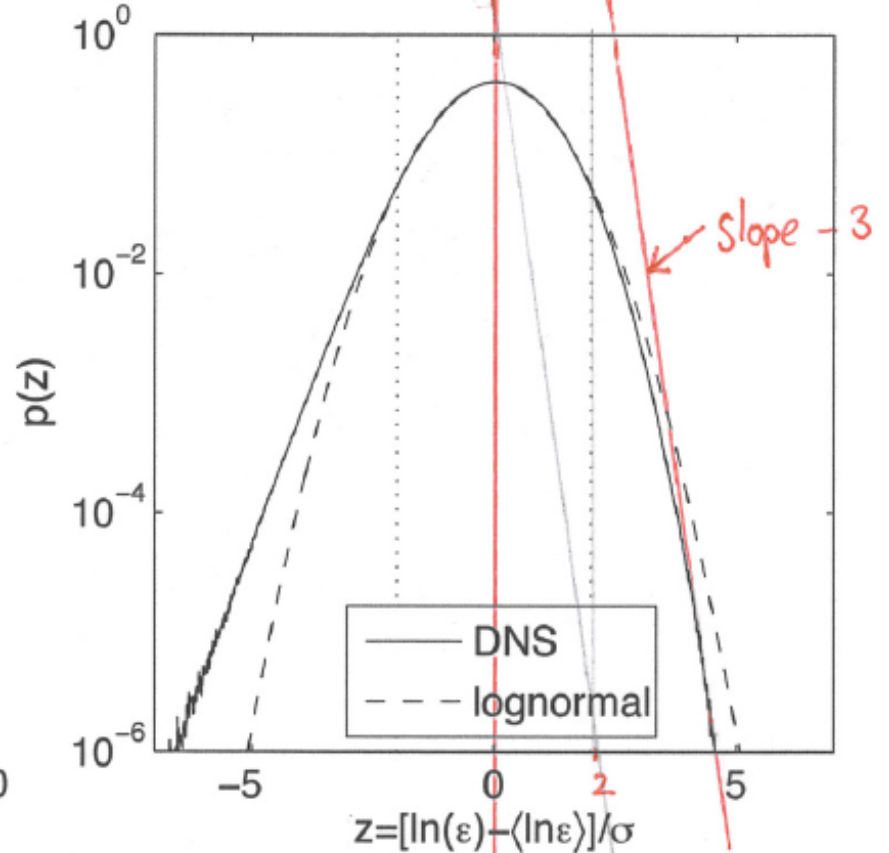
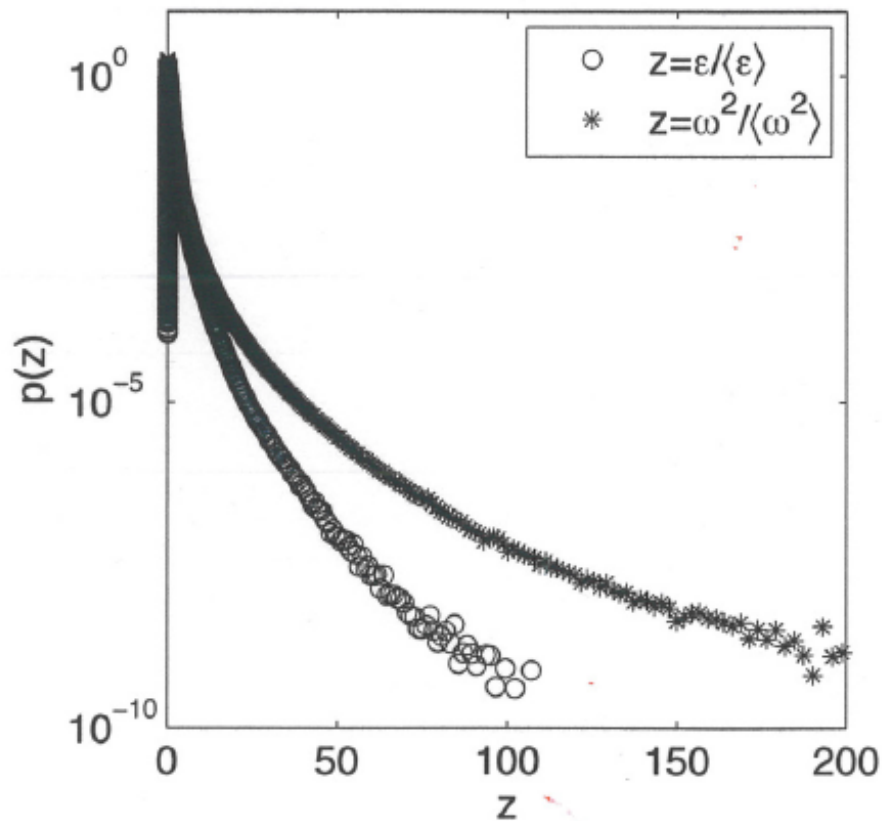


Hi Dan,

this is what I got for

$\varepsilon$ .

Cheers, Jörg



**Euler flows: ? : ?**

Fixed points associated with reconnection?

[arXiv.org](#) > [cond-mat](#) > [arXiv:cond-mat/0311487](#)

[Condensed Matter](#) > [Soft Condensed Matter](#)

## **Turbulent intermittency and Euler similarity solutions**

[Daniel P. Lathrop](#)

*(Submitted on 20 Nov 2003)*

## Euler fixed points:

Stationary

Uniform velocity

Pure strain

Axisymmetric vortices  $v_z(r)$  profiles

Reconnection fixed point?

$$\vec{\omega} = \nabla x \vec{u}$$

$$\vec{u} x \vec{\omega} = \nabla \Pi$$

Nontrivial  $\Pi$  topological constraints

Trivial  $\Pi$  yields  $\vec{u} x \vec{\omega} = 0$

$\nabla x \vec{u} = \lambda \vec{u}$  Beltrami flows, i.e. force free fields

$\vec{u} = \vec{S} + \vec{T}$  Chandrasekhar 1954.  $S_2^2$ ,  $T_2^2$  solution similar to reconnection fixed point in GP, but oscillatory in the radial direction and helical...

**Free surface flows: Topological: Wave crater collapse**

Axisymmetric:

$\delta \sim (t_o - t)^{2/3}$  length scales

$v \sim (t_o - t)^{-1/3}$  velocity scales

$Pr(v) \sim |v|^{-4}$  velocity probability distribution

Similarity solutions appear to exist

2-D like axisymmetric!

**Quantum fluid flows: Topological: Vortex reconnection**

$\delta \sim (t_o - t)^{1/2}$  length scales

$v \sim (t_o - t)^{-1/2}$  velocity scales

$Pr(v) \sim |v|^{-3}$  velocity probability distribution in turbulence

Similarity solutions exist

Fixed points with reconnection geometry exists

Safety!

Danger!

**Navier-Stokes flows: ?: turbulence**

$Pr(v_i) \sim \exp(-v_i^2/\sigma_v^2)$

$Pr(\partial_j v_i) \sim \exp(-|v_i|/\sigma_s)$

$Pr(\epsilon)$  dissipation interesting!

$Pr(\Omega)$  enstrophy interesting!

**Euler flows: ?: ?**

Fixed points associated with reconnection?