

ON GENERALIZED STRONG AND NORM RESOLVENT CONVERGENCE

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ABSTRACT. We present a streamlined approach for generalized strong and norm convergence of self-adjoint operators in different Hilbert spaces. In particular, we establish convergence of associated (semi-)groups, (essential) spectra and spectral projections. In addition, we give some applications to Sturm–Liouville operators.

1. INTRODUCTION

The standard notions of convergence for unbounded self-adjoint operators are strong and norm resolvent convergence. In the classical setting, all operators are defined on the same Hilbert space. However, there are many situations of practical interest where two operators are defined on different Hilbert spaces. For example, this occurs when an operator is restricted to a sequence of subspaces. To accommodate such situations, it is natural to work with a corresponding sequence of embedding operators and this is commonly known as generalized strong/norm resolvent convergence and dates at least back to Stolz and Weidmann [10, 11] who applied this to the approximation of singular Sturm–Liouville operators with regular ones. Further extensions were given by Weidmann in his monograph [15, Section 9.3] (see also [13] and [12, Section 6.6] and the extension to non-self-adjoint operators in [3]). Other applications include the case of Sturm–Liouville operators (or, more general, canonical systems, elliptic partial differential operators) with varying weight functions or the approximation of Schrödinger operators on different graphs. In particular, for the last purpose, generalized strong/norm resolvent convergence was further extended by Post in [6] and in his monograph [7, Chapter 4]. See also the comparison of different concepts by Post and Zimmer [9] and the overview [8]. Furthermore, perturbation theory for such situations also plays an important role in the study of indefinite Sturm–Liouville operators [2]. Our main motivation here stems from applications to Sturm–Liouville operators with varying weight functions.

The main purpose of the present paper is to advocate a streamlined approach to both generalized strong and norm resolvent convergence. Our motivation is to reduce the number of requirements for the embedding operators and only work with one natural condition, equations (2.2) and (2.1), for generalized strong and norm resolvent convergence, respectively. In contradistinction to Stolz and Weidmann, we do not require our embedding operators to be projections. While in the case of strong resolvent convergence projections are sufficient for the applications we are aware of, the present slightly more general approach seems natural also in this

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context and does not impose any additional difficulties. In contradistinction to Post, our conditions are stronger (but still sufficient for the applications we have in mind) and we do not introduce a quantified concept of δ -quasi-unitary, since we are only interested in limiting results. This allows for a streamlined approach with a reduced number of requirements providing the (from our point of view) main results, including full proofs, in just seven pages. Overall we follow the presentation of [12, Section 6.6] (which in turn is inspired by [15, Section 9.3]) and we have tried to further simplify proofs (to compensate for the extra complications added by the embedding operators) and we also include some amplifications, which even go beyond classical results at a few points.

In our final section we will provide some applications of these results and show how to improve some recent results which have been obtained using different methods.

Overall we hope that this result will foster further research in this direction as we consider this as an important topic in operator theory which deserves more attention.

2. GENERALIZED STRONG AND NORM RESOLVENT CONVERGENCE

Let $\mathfrak{H}$ and $\mathfrak{H}_n$, $n \in \mathbb{N}$, be separable Hilbert spaces together with bounded operators $J_n \in \mathcal{B}(\mathfrak{H}, \mathfrak{H}_n)$ such that

$$\|J_n\| \leq C, \quad \|J_n^* J_n - \mathbb{I}\| = \|J_n J_n^* - \mathbb{I}\| \rightarrow 0, \quad (2.1)$$

where C is a positive constant and J_n^* is the adjoint operator of J_n . While this is similar to [7, Chapter 4], we have opted for a slightly more streamlined definition, which is simpler and more natural for our intended purpose. In this connection note that (2.1) implies that $J_n^* J_n$ and $J_n J_n^*$ are both invertible for sufficiently large n . Hence, it is not hard to see that both $J_n^{-1} = J_n^* (J_n J_n^*)^{-1}$ and $(J_n^*)^{-1} = (J_n J_n^*)^{-1} J_n$ exist in this case with $\|J_n^* - J_n^{-1}\| \rightarrow 0$, $\|J_n\| = \|J_n^*\| \rightarrow 1$ and $\|J_n^{-1}\| = \|(J_n^*)^{-1}\| \rightarrow 1$. In particular, the condition $\|J_n\| \leq C$ is redundant.

We will also look at the case where

$$\|J_n\| \leq C, \quad \|J_n^* J_n \psi - \psi\| \rightarrow 0, \quad \psi \in \mathfrak{H}. \quad (2.2)$$

In this case, J_n does not have to be eventually invertible and we only have that $\liminf_{n \rightarrow \infty} \|J_n\| \geq 1$. However, in the applications we have in mind, J_n will be surjective and hence J_n^* will have a bounded inverse by Banach's closed range theorem (the converse being also true).

Example 2.1. The first example we have in mind are weighted Lebesgue spaces $\mathfrak{H} := L^2(\mathbb{R}^d, w_\infty(x)dx)$, $\mathfrak{H}_n := L^2(\mathbb{R}^d, w_n(x)dx)$ with $w_n(x) \leq Cw_\infty(x)$ and $\|\frac{w_n}{w_\infty} - 1\|_\infty \rightarrow 0$, where we set $\frac{0}{0} := 1$. In this case $J_n \psi = \psi$ is simply the natural embedding and J_n^* is multiplication with $\frac{w_n}{w_\infty}$. Equation (2.1) will hold if $\frac{w_n}{w_\infty} \rightarrow 1$ uniformly.

The other example is where $\mathfrak{H}_n \subseteq \mathfrak{H}$ is a sequence of subspaces. In this case J_n is the corresponding orthogonal projection and J_n^* is the natural embedding. Note that we have $J_n J_n^* = \mathbb{I}$ and (2.2) will hold if $J_n \rightarrow \mathbb{I}$ strongly.

Let $A \in \mathfrak{C}(\mathfrak{H})$, $A_n \in \mathfrak{C}(\mathfrak{H}_n)$ be (densely defined) closed operators. Then we say that A_n converges to A in generalized norm resolvent sense if (2.1) holds and

$$\lim_{n \rightarrow \infty} \|J_n^* R_{A_n}(z) J_n - R_A(z)\| = 0 \quad (2.3)$$

for one $z \in \Gamma = \mathbb{C} \setminus \Sigma$, $\Sigma := \overline{\sigma(A) \cup \bigcup_n \sigma(A_n)}$. Using (2.1) this is equivalent to

$$\lim_{n \rightarrow \infty} \|R_{A_n}(z)J_n - J_n R_A(z)\| = 0. \quad (2.4)$$

Similarly, we say that A_n converges to A in generalized strong resolvent sense if (2.2) holds and

$$\lim_{n \rightarrow \infty} \|J_n^* R_{A_n}(z)J_n \psi - R_A(z)\psi\| = 0, \quad \psi \in \mathfrak{H}, \quad (2.5)$$

for one $z \in \Gamma$. As before, by (2.2), this is implied by

$$\lim_{n \rightarrow \infty} \|R_{A_n}(z)J_n \psi - J_n R_A(z)\psi\| = 0, \quad \psi \in \mathfrak{H}, \quad (2.6)$$

but the converse might not be true unless J_n is surjective.

Remark 2.2. In the classical case, one often avoids taking the closure in the definition of Σ . Indeed, since strong/norm convergent sequences are bounded, a z for which (2.3) or (2.5) holds will always be in the interior of Γ in this case. In our situation, this is still true for (2.3) since J_n is invertible as pointed out before. It is also true for (2.5) if J_n is an orthogonal projection or if J_n is invertible. However, in general it could fail and it does not seem worth to add extra conditions to include these additional points in Γ .

Theorem 2.3. *Let A_n , A be self-adjoint operators and suppose that A_n converges to A in generalized norm resolvent sense. Then*

$$\lim_{n \rightarrow \infty} J_n^* f(A_n)J_n = \lim_{n \rightarrow \infty} J_n^{-1} f(A_n)J_n = f(A) \quad (2.7)$$

for every bounded continuous function $f: \Sigma \rightarrow \mathbb{C}$ with $\lim_{\lambda \rightarrow -\infty} f(\lambda) = \lim_{\lambda \rightarrow \infty} f(\lambda)$.

If A_n converges to A in generalized strong resolvent sense, then

$$\text{s-lim}_{n \rightarrow \infty} J_n^* f(A_n)J_n = f(A) \quad (2.8)$$

for every bounded continuous function $f: \Sigma \rightarrow \mathbb{C}$.

Proof. The set of functions for which the claim holds forms a $*$ -subalgebra. That it is a subalgebra follows since addition, scalar multiplication and products are continuous with respect to norm/strong convergence. To see that the claim holds for f^* if it holds for f , use that $\|J_n^* f(A_n)J_n - f(A)\| = \|J_n^* f^*(A_n)J_n - f^*(A)\|$ in the case of norm convergence. In the case of strong convergence, note that $J_n^* f(A_n)J_n \psi \rightarrow f(A)\psi$ implies weak convergence $J_n^* f^*(A_n)J_n \psi \rightharpoonup f^*(A)\psi$ and $\|J_n^* f(A_n)J_n \psi\| = \|J_n^* f^*(A_n)J_n \psi\| \rightarrow \|f(A)\psi\| = \|f^*(A)\psi\|$ since $f(A_n)$ and $f(A)$ are normal. Together, both show $J_n^* f^*(A_n)J_n \psi \rightarrow f^*(A)\psi$.

Since this $*$ -subalgebra contains $f(\lambda) = 1$ and $f(\lambda) = \frac{1}{\lambda - z_0}$, $z_0 \in \Gamma$, it is dense by the Stone–Weierstraß theorem. The usual $\frac{\varepsilon}{3}$ argument shows that this $*$ -subalgebra is also closed. $\square$

Two easy consequences are worth noticing:

Corollary 2.4. *Suppose A_n converges to A in generalized strong or norm resolvent sense for one $z_0 \in \Gamma$. Then this holds for all $z \in \Gamma$.*

Corollary 2.5. *Suppose A_n converges to A in generalized strong resolvent sense. Then*

$$J_n^* e^{itA_n} J_n \xrightarrow{s} e^{itA}, \quad t \in \mathbb{R}, \quad (2.9)$$

and if all operators are semibounded by the same bound

$$J_n^* e^{-tA_n} J_n \xrightarrow{s} e^{-tA}, \quad t \geq 0. \quad (2.10)$$

The last claim also holds in norm, provided we have generalized norm resolvent convergence.

Next we turn to criteria to verify resolvent convergence. The first one deals with the case when the difference is relatively bounded.

Lemma 2.6. *Let A_n, A be self-adjoint operators with $J_n \mathfrak{D}(A) \subseteq \mathfrak{D}(A_n)$. Then A_n converges to A in generalized norm resolvent sense if there are non-negative sequences a_n and b_n converging to zero such that*

$$\|(A_n J_n - J_n A)\psi\| \leq a_n \|A\psi\| + b_n \|\psi\|, \quad \psi \in \mathfrak{D}(A). \quad (2.11)$$

Proof. From the second resolvent identity

$$R_{A_n}(z)J_n - J_n R_A(z) = R_{A_n}(z)(J_n A - A_n J_n)R_A(z), \quad z \in \Gamma, \quad (2.12)$$

we infer

$$\begin{aligned} \|(R_{A_n}(i)J_n - J_n R_A(i))\psi\| &\leq \|R_{A_n}(i)\| \left(a_n \|A R_A(i)\psi\| + b_n \|R_A(i)\psi\| \right) \\ &\leq (a_n + b_n) \|\psi\| \end{aligned}$$

and hence $\|R_{A_n}(i)J_n - J_n R_A(i)\| \leq a_n + b_n \rightarrow 0$. $\square$

The second deals with the case when the difference is relatively form bounded.

Let $A \in \mathfrak{C}(\mathfrak{H})$ be a closed linear operator. Then the expression

$$\mathbf{q}_A(\psi) = \langle \psi, A\psi \rangle, \quad \psi \in \mathfrak{D}(A),$$

is called the quadratic form associated to A . The operator A is called nonnegative if for all $\psi \in \mathfrak{D}(A)$ we have $\mathbf{q}_A(\psi) \geq 0$. For a nonnegative operator A , introduce the scalar product

$$\langle \varphi, \psi \rangle_A := \langle \varphi, (A + 1)\psi \rangle, \quad \varphi, \psi \in \mathfrak{D}(A).$$

Let $\mathfrak{H}_A$ be the completion of $\mathfrak{D}(A)$ with respect to the above scalar product. It is easy to verify that $\mathfrak{D}(A) \subseteq \mathfrak{H}_A \subseteq \mathfrak{H}$. Clearly, the quadratic form $\mathbf{q}_A$ can be extended to every $\psi \in \mathfrak{H}_A$ by setting

$$\mathbf{q}_A(\psi) = \langle \psi, \psi \rangle_A - \|\psi\|^2, \quad \psi \in \mathfrak{Q}(A) = \mathfrak{H}_A.$$

The set $\mathfrak{Q}(A)$ is also called the form domain of A . For more details, see [12].

Lemma 2.7. *Suppose $A, A_n \geq \gamma$ are bounded from below with $J_n \mathfrak{Q}(A) \subseteq \mathfrak{Q}(A_n)$. Then A_n converges to A in generalized norm resolvent sense if there are non-negative sequences a_n and b_n converging to zero such that*

$$|\mathbf{q}_A(\psi) - \mathbf{q}_{A_n}(J_n \psi)| \leq a_n \mathbf{q}_{A-\gamma}(\psi) + b_n \|\psi\|^2, \quad \psi \in \mathfrak{Q}(A),$$

where $\mathbf{q}_A$ and $\mathbf{q}_{A_n}$ are the associated quadratic forms.

Proof. We can assume $a_n, b_n < 1$ without loss of generality. Applying [12, Theorem 6.25] with $\mathbf{q}(\psi) = \mathbf{q}_{A_n}(J_n \psi) - \mathbf{q}_A(\psi)$ for any $\psi \in \mathfrak{Q}(A)$ we have

$$\|C_{\mathbf{q}}(\lambda)\| \leq \max \left(a_n, \frac{b_n}{\lambda + \gamma} \right) \rightarrow 0, \quad \lambda > -\gamma,$$

and

$$R_{B_n}(-\lambda) = R_A(-\lambda)^{\frac{1}{2}} (\mathbb{I} + C_{\mathbf{q}}(\lambda))^{-1} R_A(-\lambda)^{\frac{1}{2}},$$

where $C_{\mathbf{q}}(\lambda)$ is the bounded operator corresponding to the quadratic form $\mathbf{q}(R_A(-\lambda)^{\frac{1}{2}} \cdot)$ and $B_n := J_n^* A_n J_n$ is the self-adjoint operator associated with the quadratic form $\mathbf{q}_{A_n}(J_n \cdot)$. Hence we have

$$\begin{aligned} \|R_{B_n}(-\lambda) - R_A(-\lambda)\| &= \|R_A(-\lambda)^{\frac{1}{2}} ((1 + C_{\mathbf{q}}(\lambda))^{-1} - 1) R_A(-\lambda)^{\frac{1}{2}}\| \\ &\leq \|R_A(-\lambda)^{\frac{1}{2}}\|^2 \|(1 + C_{\mathbf{q}}(\lambda))^{-1} - 1\| \rightarrow 0. \end{aligned}$$

Moreover, we have

$$\begin{aligned} \|J_n R_{B_n}(-\lambda) - R_{A_n}(-\lambda) J_n\| &= \|R_{A_n}(-\lambda)(A_n + \lambda)(J_n R_{B_n}(-\lambda) - R_{A_n}(-\lambda) J_n)\| \\ &\leq \|R_{A_n}(-\lambda)\| \|(A_n + \lambda) J_n R_{B_n}(-\lambda) - J_n\| \\ &\leq \|R_{A_n}(-\lambda)\| \|(J_n^*)^{-1}\| \|(B_n + \lambda + \lambda(J_n^* J_n - \mathbb{I})) R_{B_n}(-\lambda) - J_n^* J_n\| \\ &\leq \|R_{A_n}(-\lambda)\| \|(J_n^*)^{-1}\| \|(\lambda R_{B_n}(-\lambda) - 1)\| \|J_n^* J_n - \mathbb{I}\| \rightarrow 0, \end{aligned} \quad \square$$

from which the claim follows.

Next, we establish that strong convergence of the operators implies strong resolvent convergence.

Let A be a symmetric operator with dense domain $\mathfrak{D}(A)$ and $\bar{A}$ is its only self-adjoint extension. Then A is called essentially self-adjoint and $\mathfrak{D}(A)$ is called a core for $\bar{A}$. For more details, see [12].

Lemma 2.8. *The sequence A_n converges to A in generalized strong resolvent sense if there is a core $\mathfrak{D}_0$ of A such that for every $\psi \in \mathfrak{D}_0$ we have $J_n \psi \in \mathfrak{D}(A_n)$ for n sufficiently large and $(A_n J_n - J_n A) \psi \rightarrow 0$.*

Proof. Using the second resolvent formula (2.12), we have

$$(R_{A_n}(\mathbf{i}) J_n - J_n R_A(\mathbf{i})) \phi = -R_{A_n}(\mathbf{i}) (A_n J_n - J_n A) \psi \rightarrow 0$$

for $\phi = (A - \mathbf{i}) \psi$ with $\psi \in \mathfrak{D}_0$ which is dense, since $\mathfrak{D}_0$ is a core. The rest follows from [12, Lemma 1.14]. $\square$

Finally, we mention that generalized weak resolvent convergence, is equivalent to the generalized strong resolvent convergence.

Lemma 2.9. *Let A_n, A be self-adjoint operators. Suppose (2.2) and $\|J_n J_n^* - \mathbb{I}\| \rightarrow 0$. If $J_n^* R_{A_n}(z) J_n$ converges to $R_A(z)$ weakly for some $z \in \Gamma \setminus \mathbb{R}$, then A_n converges to A in generalized strong resolvent sense.*

Proof. By $J_n^* R_{A_n}(z) J_n \rightharpoonup R_A(z)$ we also have $J_n^* R_{A_n}(z)^* J_n \rightharpoonup R_A(z)^*$ and thus by the first resolvent formula

$$\begin{aligned} &|\|J_n^* R_{A_n}(z) J_n \psi\|^2 - \|R_A(z) \psi\|^2| \\ &= |\langle \psi, J_n^* R_{A_n}(z^*) J_n J_n^* R_{A_n}(z) J_n \psi - R_A(z^*) R_A(z) \psi \rangle| \\ &\leq |\langle \psi, J_n^* R_{A_n}(z^*) (J_n J_n^* - \mathbb{I}) R_{A_n}(z) J_n \psi \rangle| \\ &\quad + |\langle \psi, (J_n^* R_{A_n}(z^*) R_{A_n}(z) J_n - R_A(z^*) R_A(z)) \psi \rangle| \\ &\leq \|J_n^*\| \|R_{A_n}(z^*)\| |\langle \psi, (J_n J_n^* - \mathbb{I}) R_{A_n}(z) J_n \psi \rangle| \\ &\quad + \frac{1}{|z - z^*|} |\langle \psi, (J_n^* R_{A_n}(z) J_n) \psi - (J_n^* R_{A_n}(z)^* J_n) \psi \rangle| \\ &\rightarrow 0. \end{aligned}$$

Together with $J_n^* R_{A_n}(z) J_n \rightharpoonup R_A(z)$ we have $J_n^* R_{A_n}(z) J_n \psi \rightarrow R_A(z) \psi$ by virtue of [12, Lemma 1.12 (iv)]. $\square$

Now we turn to convergence of spectra. As usual, we denote by $\lim_{n \rightarrow \infty} S_n$ the set of all λ for which there is a sequence $\lambda_n \in S_n$ converging to λ .

Theorem 2.10. *Let A_n and A be self-adjoint operators. If A_n converges to A in generalized strong resolvent sense, we have $\sigma(A) \subseteq \lim_{n \rightarrow \infty} \sigma(A_n)$.*

If A_n converges to A in generalized norm resolvent sense, we have $\sigma(A) = \lim_{n \rightarrow \infty} \sigma(A_n)$. Moreover, for every $\lambda \in \rho(A)$ there is a neighborhood which is eventually contained in $\rho(A_n)$.

Proof. First, we claim $\sigma(A) \subseteq \lim_{n \rightarrow \infty} \sigma(A_n)$. If not, then we could find a $\lambda \in \sigma(A)$ and some $\varepsilon > 0$ such that

$$\sigma(A_n) \cap (\lambda - 2\varepsilon, \lambda + 2\varepsilon) = \emptyset.$$

Choose a nonnegative bounded continuous function f , which takes the value 1 on $(\lambda - \varepsilon, \lambda + \varepsilon)$ and vanishes outside $(\lambda - 2\varepsilon, \lambda + 2\varepsilon)$. Then $f(A_n) = 0$. Combining with Theorem 2.3, we have

$$J_n^* f(A_n) J_n \psi - f(A) \psi \rightarrow 0 \quad \text{for any } \psi \in \mathfrak{D}(A),$$

and thus $f(A) \psi = 0$. On the other hand, since $\lambda \in \sigma(A)$, there is a non-zero $\psi \in \text{Ran } P_A((\lambda - \varepsilon, \lambda + \varepsilon))$ implying $f(A) \psi = \psi$, which is a contradiction.

To get the converse direction, we choose $\lambda \in \rho(A)$ and $\varepsilon > 0$ such that $\sigma(A) \cap (\lambda - 2\varepsilon, \lambda + 2\varepsilon) = \emptyset$. Choose a nonnegative bounded continuous function f which takes the value 1 on $(\lambda - \varepsilon, \lambda + \varepsilon)$ and vanishes outside $(\lambda - 2\varepsilon, \lambda + 2\varepsilon)$. Then $f(A) = 0$ and by $J_n^* f(A_n) J_n \rightarrow 0$ we also have $f(A_n) \rightarrow 0$ and

$$f(A_n) P_{A_n}((\lambda - \varepsilon, \lambda + \varepsilon)) = P_{A_n}((\lambda - \varepsilon, \lambda + \varepsilon)) \rightarrow 0.$$

This implies $P_{A_n}((\lambda - \varepsilon, \lambda + \varepsilon)) = 0$ eventually and establishes the claim. $\square$

In addition, we have the following results for the convergence of spectral projections.

Lemma 2.11. *Suppose A_n converges in generalized strong resolvent sense to A . If $P_A(\{\lambda\}) = 0$, then*

$$\text{s-lim}_{n \rightarrow \infty} J_n^* P_{A_n}((-\infty, \lambda)) J_n = \text{s-lim}_{n \rightarrow \infty} J_n^* P_{A_n}((-\infty, \lambda]) J_n = P_A((-\infty, \lambda)) = P_A((-\infty, \lambda]). \quad (2.13)$$

Proof. By Theorem 2.3, the spectral measures $\mu_{n, J_n \psi}$ corresponding to A_n converge vaguely to the spectral measure μ_ψ of A . Hence $\|P_{A_n}(\Omega) J_n \psi\|^2 = \mu_{n, J_n \psi}(\Omega)$ together with [12, Lem. A.34] implies the claim. $\square$

Using $P((\lambda_0, \lambda_1)) = P((-\infty, \lambda_1)) - P((-\infty, \lambda_0])$, we also obtain

Lemma 2.12. *Suppose A_n converges in the generalized strong resolvent sense to A . If $P_A(\{\lambda_0\}) = P_A(\{\lambda_1\}) = 0$, then*

$$\text{s-lim}_{n \rightarrow \infty} J_n^* P_{A_n}((\lambda_0, \lambda_1)) J_n = \text{s-lim}_{n \rightarrow \infty} J_n^* P_{A_n}([\lambda_0, \lambda_1]) J_n = P_A((\lambda_0, \lambda_1)) = P_A([\lambda_0, \lambda_1]). \quad (2.14)$$

Moreover,

$$\begin{aligned} \dim \text{Ran } P_A((\lambda_0, \lambda_1)) &\leq \liminf_{n \rightarrow \infty} \dim \text{Ran } J_n^* P_{A_n}((\lambda_0, \lambda_1)) J_n \\ &\leq \limsup_{n \rightarrow \infty} \|J_n\|^2 \liminf_{n \rightarrow \infty} \dim \text{Ran } P_{A_n}((\lambda_0, \lambda_1)). \end{aligned} \quad (2.15)$$

Proof. It remains to show the inequalities. Since we can write

$$\dim \operatorname{Ran} P_{A_{n_k}}((\lambda_0, \lambda_1)) = \operatorname{tr}(P_{A_{n_k}}((\lambda_0, \lambda_1)))$$

the first follows from strong convergence of the spectral projectors together with Fatou's lemma. The second inequality follows since for every positive operator P of rank N we have $\operatorname{tr}(P) \leq N\|P\|$. $\square$

Corollary 2.13. *Suppose A_n converges in the generalized strong resolvent sense to A . If $\lambda \in \sigma_{\text{ess}}(A)$, then for every $\varepsilon > 0$ we have*

$$\lim_{n \rightarrow \infty} \dim \operatorname{Ran} P_{A_n}((\lambda - \varepsilon, \lambda + \varepsilon)) = \infty. \quad (2.16)$$

Proof. By decreasing ε if necessary, we can assume $P_A(\{\lambda - \varepsilon\}) = P_A(\{\lambda + \varepsilon\}) = 0$. Hence the claim follows from the previous lemma since the assumption implies $\dim \operatorname{Ran} P_A((\lambda - \varepsilon, \lambda + \varepsilon)) = \infty$ for every $\varepsilon > 0$. $\square$

For our next result recall the following well-known results: If P, Q are two projections with $\|P - Q\| < 1$, then the dimensions of their ranges agree.

Lemma 2.14. *Suppose A_n converges to A in generalized norm resolvent sense and λ is an isolated point of $\sigma(A)$. Then*

$$\lim_{n \rightarrow \infty} J_n^{-1} P_{A_n}((\lambda - \varepsilon, \lambda + \varepsilon)) J_n = P_A(\{\lambda\}), \quad (2.17)$$

for ε so small that $(\lambda - \varepsilon, \lambda + \varepsilon) \cap \sigma(A) = \{\lambda\}$. In particular, the dimensions of the ranges of both projections eventually agree.

Proof. Let f be a continuous function which is 1 on $(\lambda - \varepsilon/2, \lambda + \varepsilon/2)$ and zero outside $(\lambda - \varepsilon, \lambda + \varepsilon)$. Then, for sufficiently large n we have that $\sigma(A_n) \cap \{\alpha \mid \varepsilon/2 < |\alpha - \lambda| < \varepsilon\} = \emptyset$. Hence $f(A_n) = P_{A_n}((\lambda - \varepsilon, \lambda + \varepsilon))$ and $f(A) = P_A(\{\lambda\})$ such that $J_n^{-1} P_{A_n}((\lambda - \varepsilon, \lambda + \varepsilon)) J_n \rightarrow P_A(\{\lambda\})$ and the claim follows. $\square$

The following is the analog of [15, Satz 9.24] and [7, Theorem 4.3.4] in this context.

Theorem 2.15. *Let A_n and A be self-adjoint operators. If A_n converges to A in generalized norm resolvent sense, we have $\sigma_{\text{ess}}(A) = \lim_{n \rightarrow \infty} \sigma_{\text{ess}}(A_n)$. Here $\lim_{n \rightarrow \infty} \sigma_{\text{ess}}(A_n)$ denotes the set of all λ for which there is a sequence $\lambda_n \in \sigma_{\text{ess}}(A_n)$ converging to λ .*

Proof. If $\lambda \notin \sigma_{\text{ess}}(A)$, then λ is in the discrete spectrum (or in the resolvent set) and the claim follows from Lemma 2.14.

Conversely, let $\lambda \in \sigma_{\text{ess}}(A)$. If there were no sequence $\lambda_n \in \sigma_{\text{ess}}(A_n)$ converging to λ , we could find $\varepsilon > 0$ and a subsequence n_k such that $\dim \operatorname{Ran} P_{A_{n_k}}((\lambda - \varepsilon, \lambda + \varepsilon)) < \infty$. We pass to the subsequence for notational convenience. Moreover, we can choose a singular Weyl sequence ψ_k with $\|\psi_k\| = 1$, $\psi_k \rightharpoonup 0$, and $\|(R_A(i) -$

$(\lambda - i)^{-1})\psi_k\| \rightarrow 0$. Then

$$\begin{aligned}
\|J_n^* R_{A_n}(i)J_n - R_A(i)\| &\geq \limsup_{k \rightarrow \infty} \|(J_n^* R_{A_n}(i)J_n - R_A(i))\psi_k\| \\
&= \limsup_{k \rightarrow \infty} \|(J_n^* R_{A_n}(i)J_n - (\lambda - i)^{-1})\psi_k\| \\
&= \limsup_{k \rightarrow \infty} \|J_n^*(R_{A_n}(i) - (\lambda - i)^{-1})J_n\psi_k\| \\
&= \limsup_{k \rightarrow \infty} \|J_n^*(R_{A_n}(i) - (\lambda - i)^{-1})P_{A_n}(\mathbb{R} \setminus (\lambda - \varepsilon, \lambda + \varepsilon))J_n\psi_k\| \\
&\geq \|(J_n^*)^{-1}\|^{-1} \left(\frac{1}{\lambda^2 + 1} - \frac{1}{(|\lambda| + \varepsilon)^2 + 1} \right) \|J_n^{-1}\|^{-1}
\end{aligned}$$

contradicting norm resolvent convergence. Here we have used that $J_n^* P_{A_n}((\lambda - \varepsilon, \lambda + \varepsilon))J_n\psi_k \rightarrow 0$ by compactness. $\square$

3. APPLICATIONS TO STURM-LIOUVILLE OPERATORS

In this section, we apply the generalized norm resolvent convergence to Sturm-Liouville operators in different Hilbert spaces. Throughout this section we will make the usual assumptions on the coefficients:

Hypothesis 3.1. Assume that $I = (a, b)$ is some open interval. Let $w, p^{-1}, q \in L^1_{loc}(I, dx)$ be real-valued and $p, w > 0$ a.e. in I .

Then we consider the Sturm-Liouville operator:

$$\begin{aligned}
T : \quad \mathfrak{D}(T) &\rightarrow L^2(I, w(x)dx) \\
f &\mapsto \frac{1}{w} \left(-\frac{d}{dx} p \frac{df}{dx} + qf \right),
\end{aligned} \tag{3.1}$$

where

$$\mathfrak{D}(T) \subseteq \{f \in L^2(I, w(x)dx) : f, pf' \in AC_{loc}(I), \frac{1}{w}(-(pf')' + qf) \in L^2(I, w(x)dx)\}. \tag{3.2}$$

In order for T to be self-adjoint, one might have to impose additional boundary conditions, but we postpone this issue for now.

For our present purpose, we will need the corresponding form domain. Most results we are aware of consider the case of Schrödinger operators ($w = p = 1$), see for example [12], or the case of regular operators, see [4]. To compute the form domain in a more general situation, we will first set q equal to 0 and factorize T according to A^*A such that we get $\mathfrak{Q}(T) = \mathfrak{Q}(A)$.

To this end define

$$\begin{aligned}
A_+ : \quad \mathfrak{D}(A_+) &\rightarrow L^2(I, w(x)dx) \\
f &\mapsto \frac{1}{w} \sqrt{wp} \frac{df}{dx},
\end{aligned}$$

where $\mathfrak{D}(A_+) := \{f \in L^2(I, w(x)dx) : f \in AC_{loc}(I), \sqrt{p}f' \in L^2(I, dx)\}$ and

$$\begin{aligned}
A_- : \quad \mathfrak{D}(A_-) &\rightarrow L^2(I, w(x)dx) \\
f &\mapsto -\frac{1}{w} \frac{d}{dx} \sqrt{wp} f,
\end{aligned}$$

where $\mathfrak{D}(A_-) := \{f \in L^2(I, w(x)dx) : \sqrt{wp}f \in AC_{loc}(I), \frac{1}{\sqrt{w}}(\sqrt{wp}f)' \in L^2(I, dx)\}$.

Moreover, we will also need the associated minimal operators $A_{0,\pm}$ which are given by $A_{\pm}$ restricted to functions with compact support.

Note that for $(c, d) \subseteq (a, b)$, $f \in \mathfrak{D}(A_+)$ and $g \in \mathfrak{D}(A_-)$ integration by parts shows

$$\begin{aligned} & \int_c^d g(x) \sqrt{w(x)p(x)} f'(x) dx \\ &= (\sqrt{wpg})(d)f(d) - (\sqrt{wpg})(c)f(c) - \int_c^d (\sqrt{wpg})'(x)f(x) dx. \end{aligned} \quad (3.3)$$

Moreover, since all integrands are integrable, the limits of the boundary terms exist as $c \downarrow a$ and $d \uparrow b$.

Also, if $g \in L^2(I, w(x)dx)$ then we have that

$$f_+(x) := f_+(c) + \int_c^x \sqrt{\frac{w(y)}{p(y)}} g(y) dy \quad (3.4)$$

satisfies $f_+ \in AC_{loc}(I)$ and $\sqrt{p/w}(f_+)' = g$ and

$$f_-(x) = \frac{1}{\sqrt{p(x)w(x)}} \left((\sqrt{wpg})(c) + \int_x^c w(y)g(y) dy \right) \quad (3.5)$$

satisfies $\sqrt{wpg}f_- \in AC_{loc}(I)$ and $-1/w(\sqrt{wpg}f_-)' = g$.

The following type of result is well-known. We include a proof for the sake of completeness.

Lemma 3.2. *The operators $A_{0,\pm}$ are densely defined and their closures are given by*

$$\overline{A_{0,+}}f = A_+f, \quad \mathfrak{D}(\overline{A_{0,+}}) = \{f \in \mathfrak{D}(A_+) : \lim_{x \rightarrow a,b} (\sqrt{wpg})(x)f(x) = 0, \forall g \in \mathfrak{D}(A_-)\}$$

and

$$\overline{A_{0,-}}f = A_-f, \quad \mathfrak{D}(\overline{A_{0,-}}) = \{f \in \mathfrak{D}(A_-) : \lim_{x \rightarrow a,b} (\sqrt{wpg})(x)g(x) = 0, \forall g \in \mathfrak{D}(A_+)\}.$$

Their adjoint operators are given by

$$A_{0,\pm}^*f = A_{\mp}f, \quad \mathfrak{D}(A_{0,\pm}^*) = \mathfrak{D}(A_{\mp}).$$

Proof. By (3.3) we have $\mathfrak{D}(A_{\mp}) \subseteq \mathfrak{D}(A_{0,\pm}^*)$ and it remains to show $\mathfrak{D}(A_{0,\pm}^*) \subseteq \mathfrak{D}(A_{\mp})$. If we have

$$\langle h, A_{0,\pm}f \rangle = \langle k, f \rangle, \quad \forall f \in \mathfrak{D}(A_{0,\pm}),$$

then by (3.4), (3.5) we can find a $\tilde{h}$ which is locally in $\mathfrak{D}(A_{\mp})$ such that $A_{\mp}\tilde{h} = k$ and integration by parts shows

$$\int_a^b (h(x) - \tilde{h}(x)^*) A_{\pm}f(x) dx = 0, \quad \forall f \in \mathfrak{D}(A_{0,\pm}).$$

Now consider the linear functionals

$$\ell(g) = \int_a^b (h(x) - \tilde{h}(x)^*) g(x) w(x) dx, \quad \ell_{\pm}(g) = \int_a^b u_{\mp}(x)^* g(x) w(x) dx,$$

on $L_c^2(I, w(x)dx)$, where $u_+ = 1$ and $u_- = \frac{1}{\sqrt{wp}}$. Then $g \in \ker(\ell_\pm)$ implies

$$f(x) = \begin{cases} \int_a^x \sqrt{\frac{w(y)}{p(y)}} g(y) dy \in \mathfrak{D}(A_{0,+}) & \text{when } g \in \ker(\ell_+) \\ \frac{1}{\sqrt{p(x)w(x)}} \int_x^b w(y)g(y)dy \in \mathfrak{D}(A_{0,-}) & \text{when } g \in \ker(\ell_-) \end{cases} \quad (3.6)$$

and $A_\pm f = g$. Hence $\ell(g) = 0$, i.e., $g \in \ker(\ell)$. So by [12, Lemma 9.3] we conclude $\ell = \alpha_\mp \ell_\pm$, where $\alpha_\mp$ are constants. And thus $h(x) = \tilde{h}(x) + \alpha_\mp u_\mp(x) \in \mathfrak{D}(A_\mp)$ as required.

If $\mathfrak{D}(A_{0,\pm})$ were not dense, then for $h = 0$ we could choose some nonzero $k \in \mathfrak{D}(A_{0,\pm})^\perp$ and our calculation from above would give the contradiction $k = A_\mp \tilde{h} = A_\mp h = 0$. So we have shown $A_{0,\pm}^* = A_\mp$.

Finally, abbreviate $\mathfrak{D}_\pm := \{f \in \mathfrak{D}(A_\pm) : \lim_{x \rightarrow a,b} (\sqrt{wp}f)(x)g(x) = 0, \forall g \in \mathfrak{D}(A_\mp)\}$. Then (3.3) shows $\mathfrak{D}_\pm \subseteq \mathfrak{D}(A_{0,\pm}^{**}) = \mathfrak{D}(\overline{A_{0,\pm}})$. Conversely, since $\overline{A_{0,\pm}} = A_\mp^* \subseteq A_\pm$, we can use (3.3) and take limits to conclude

$$\sqrt{w(b)p(b)}f(b)h(b) - \sqrt{w(a)p(a)}f(a)h(a) = 0, \quad f \in \mathfrak{D}(\overline{A_{0,\pm}}), h \in \mathfrak{D}(A_\mp).$$

Now replace h by a $\tilde{h} \in \mathfrak{D}(A_\mp)$ which coincides with h near a and vanishes identically near b . Then

$$\sqrt{w(a)p(a)}f(a)h(a) = \sqrt{w(a)p(a)}f(a)\tilde{h}(a) - \sqrt{w(b)p(b)}f(b)\tilde{h}(b) = 0.$$

Finally, $\sqrt{w(b)p(b)}f(b)h(b) = \sqrt{w(a)p(a)}f(a)h(a) = 0$ shows $f \in \mathfrak{D}_\pm$. $\square$

Based on this result, we define

$$\mathfrak{Q} := \mathfrak{D}(A_+), \quad \mathfrak{Q}_0 := \{f \in \mathfrak{Q} : \lim_{x \rightarrow a,b} (\sqrt{wp}g)(x)f(x) = 0, \forall g \in \mathfrak{D}(A_-)\}. \quad (3.7)$$

Then we get two associated operators, the Friedrichs extension $A_- \overline{A_{0,+}}$ with form domain $\mathfrak{Q}_0$ and $\overline{A_{0,-}} A_+$ with form domain $\mathfrak{Q}$. If the underlying differential expression $-\frac{1}{w} \frac{d}{dx} p \frac{d}{dx}$ is limit point at both end points, then both operators will of course agree.

In order to incorporate q , we will assume $w, p \geq \delta$ for some constant $\delta > 0$ such that $L^2(I, w(x)dx) \subseteq L^2(I, dx)$ and $\mathfrak{Q} \subseteq H^1(I, dx)$, where $H^1(I, dx) := W^{1,2}(I, dx)$ is the Sobolev space.

Lemma 3.3. *Suppose w, p, q satisfy*

$$w, p \geq \delta, \quad M := \sup_{k \in \mathbb{Z}} \int_{[k, k+1) \cap I} |q(x)| dx < \infty, \quad (3.8)$$

where δ is a positive constant. If I is bounded, then the condition on q amounts to requiring that q is integrable.

Then q/w is relatively form bounded with bound 0 with respect to both $A_- \overline{A_{0,+}}$ and $\overline{A_{0,-}} A_+$, i.e., for every $\varepsilon > 0$,

$$\int_I |q(x)| |f(x)|^2 dx \leq \frac{M}{\delta} \left(\varepsilon \int_I p(x) |f'(x)|^2 dx + \left(1 + \frac{1}{\varepsilon}\right) \int_I |f(x)|^2 w(x) dx \right), \quad f \in \mathfrak{Q}.$$

Proof. This follows literally as in Lemma 9.33 from [12]. $\square$

Hence we can define either $T = A_- \overline{A_{0,+}} + q/w$ with $\mathfrak{Q}(T) = \mathfrak{Q}_0$ or $T = \overline{A_{0,-}} A_+ + q/w$ with $\mathfrak{Q}(T) = \mathfrak{Q}$ as form sums such that the associated quadratic form is

$$\int_a^b (p(x)f'(x)^*g'(x) + q(x)f(x)^*g(x))dx, \quad f, g \in \mathfrak{Q}(T). \quad (3.9)$$

In both cases T is a well-defined semibounded self-adjoint operator by the KLMN theorem.

Theorem 3.4. *Suppose w, p, q and w_n, p_n, q_n both satisfy Hypothesis 3.1 and (3.8). If*

$$w_n/w \rightarrow 1, \quad p_n/p \rightarrow 1, \quad C_n := \sup_{k \in \mathbb{Z}} \int_{[k, k+1) \cap I} |q_n(x) - q(x)| dx \rightarrow 0, \quad (3.10)$$

uniformly, then we have $T_n \rightarrow T$ in generalized norm resolvent convergence.

Proof. It is easy to verify

$$\begin{aligned} J_n : \quad L^2(I, w(x)dx) &\rightarrow L^2(I, w_n(x)dx) \\ f &\mapsto f, \end{aligned}$$

satisfies (2.1). Since

$$\int_I |f(x)|^2 w_n(x) dx = \int_I |f(x)|^2 w(x) \frac{w_n(x)}{w(x)} dx,$$

we have $J_n \mathfrak{Q}(T) \subseteq \mathfrak{Q}(T_n)$. Let $T = T^{(1)} + T^{(2)}$, where $T^{(1)} := \frac{1}{w}(-\frac{d}{dx} p \frac{d}{dx})$ and $T^{(2)} := \frac{q}{w}$ in $L^2(I, w(x)dx)$. For any $\psi \in \mathfrak{Q}(T)$,

$$\begin{aligned} |\mathbf{q}_{T^{(1)}}(\psi) - \mathbf{q}_{T_n^{(1)}}(J_n \psi)| &= \left| \int_I \left(\frac{p_n(x)}{p(x)} - 1 \right) p(x) |\psi'(x)|^2 dx \right| \\ &\leq \left\| \frac{p_n}{p} - 1 \right\|_\infty \int_I p(x) |\psi'(x)|^2 dx, \end{aligned}$$

and by Lemma 3.3,

$$\begin{aligned} |\mathbf{q}_{T^{(2)}}(\psi) - \mathbf{q}_{T_n^{(2)}}(J_n \psi)| &\leq \int_I |q(x) - q_n(x)| |\psi(x)|^2 dx \\ &\leq C_n C(p, w) \left(\varepsilon \int_I p(x) |\psi'(x)|^2 dx + \left(1 + \frac{1}{\varepsilon} \right) \int_I |\psi(x)|^2 w(x) dx \right), \end{aligned}$$

where $C(p, w)$ is a positive constant associated with p and w .

Note that $T_n, T \geq \gamma$ are bounded from below by some fixed constant γ . Therefore, using Lemma 2.7 we finally have $T_n \rightarrow T$ in generalized norm resolvent sense. $\square$

The following result generalizes Theorem 3.5 from [1].

Theorem 3.5. *Suppose $I = \mathbb{R}$, w_0, p_0, q_0 and $w_\infty, p_\infty, q_\infty$ both satisfy Hypothesis 3.1 and (3.8). If*

$$w_\infty(x)/w_0(x) \rightarrow 1, \quad p_\infty(x)/p_0(x) \rightarrow 1 \quad \text{as } |x| \rightarrow \infty,$$

and

$$\int_k^{k+1} |q_\infty(x) - q_0(x)| dx \rightarrow 0 \quad \text{as } |k| \rightarrow \infty.$$

Then $\sigma_{ess}(T_\infty) = \sigma_{ess}(T_0)$.

Proof. Define

$$w_n(x) = \begin{cases} w_\infty(x), & |x| \leq n, \\ w_0(x), & |x| > n \end{cases}$$

and similarly for p_n and q_n . Then it is well-known that $\sigma_{\text{ess}}(T_n) = \sigma_{\text{ess}}(T_0)$ since the essential spectrum does not change when the coefficients are changed on a compact set: Adding extra (e.g.) Dirichlet boundary conditions at $-n$ and n splits the operator into a direct sum of three operators which differ from the original one by a rank two perturbation. Hence by Weyl's theorem the essential spectrum of the original operator is the union of the essential spectra of the three parts. As the middle part corresponds to a regular operator it does not contribute to the essential spectrum.

Using Theorem 3.4 we have $T_n \rightarrow T_\infty$ in generalized norm resolvent sense and the claim follows. $\square$

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