

Scalar Waves on a Naked Singularity Background

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¹Joint work with A. Shadi Tahvildar-Zadeh

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- ▶ There are no troublesome indices.
- ▶ People (mostly Wald² and students) suggest well-posedness of wave equations as a substitute for geodesic completeness.

²J. Math. Phys. 1980

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$$\mathbf{E}_s[u](t) = \|u(t)\|_{\dot{H}^s(\mathbf{R}^3)}^2 + \|\partial_t u(t)\|_{\dot{H}^{s-1}(\mathbf{R}^3)}^2$$

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- ▶ Dispersive Estimates (L^∞ decay):

$$\|u(t)\|_{L^\infty(\mathbf{R}^3)} \leq Ct^{-1} (\|\nabla u(0)\|_{L^1(\mathbf{R}^3)} + \|\partial_t u(0)\|_{L^1(\mathbf{R}^3)})$$

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One nice thing about this Strichartz is that it is Lorentz invariant. Strichartz estimates, like dispersive estimates, can be used to prove stability for non-linear wave equations, but Strichartz estimates are often true in contexts where dispersive estimates fail.

Scalar Wave Equation in Spherical Symmetry

Metric in isothermal coordinates:

$$g_{\mu\nu} dx^\mu dx^\nu = \alpha(r)^2 (-dt^2 + dr^2) + \rho(r)^2 (d\varphi^2 + \sin^2 \varphi d\theta^2)$$

α and ρ to be specified later.

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(Massless, Chargeless) Scalar Wave Equation:

$$g^{\mu\nu} \psi_{;\mu\nu} = 0$$

or

$$\partial_t^2 \psi - \frac{1}{\rho^2} \partial_r (\rho^2 \partial_r \psi) - \frac{\alpha^2}{\rho^2} \Delta_{\text{Sph}} \psi = 0.$$

Transferring to Minkowski Space

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$$u = \rho\psi/r$$

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Compared to the scalar wave equation in Minkowski space, there is an additional scalar potential

$$V(r) = \rho''(r)/\rho(r).$$

Reissner-Nordström

For Reissner-Nordström,

$$\rho'(r) = \alpha^2 \quad \alpha = \sqrt{1 - \frac{2m}{\rho} + \frac{e^2}{\rho^2}}$$

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- ▶ If $|e| = m$, the extremal case, then the quadratic has a double root at $\rho = m$, which is again a horizon and the metric is valid outside the horizon.
- ▶ If $|e| > m$, the super-extremal case, then there are no horizons and the metric above is valid for all $r > 0$, but is highly singular at $r = 0$.

Super-extremal Case

$$V = \frac{2m}{\rho^3} - \frac{2e^2 + 4m^2}{\rho^4} + \frac{6me^2}{\rho^5} - \frac{2e^4}{\rho^6}$$

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$$\lim_{r \rightarrow 0^+} r^2 V(r) = -\frac{2}{9}$$

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We are forced by the problem to choose

$$Au = -\frac{1}{r^2} \partial_r (r^2 \partial_r u) - \frac{\alpha^2}{\rho^2} \Delta_{\text{Sph}} u + V(r)u,$$

but with what domain?

Wait! Is the Equation Well Defined? (Continued)

For smooth functions u, v supported on compact subsets of $\mathbf{R}^3 - 0$, we have

- ▶ Positive Definiteness:

$$\langle u, Au \rangle \geq 0 \quad \langle u, Au \rangle = 0 \implies u = 0$$

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Almost all symmetric differential operators appearing in Mathematical Physics are essentially self-adjoint.

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An alternate characterisation, due to Krein, is this: If A_F is the Friedrichs extension and A_E is any other positive self-adjoint extension of A then

$$u \in \text{Dom}(A_F) \implies u \in \text{Dom}(A_E) \text{ and } \langle u, A_F u \rangle \leq \langle u, A_E u \rangle.$$

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Now that we have a well defined problem, we can use a theorem we proved earlier with Burq and Planchon:⁴

Let $V \in C^1(\mathbf{R}^+)$ satisfy

- ▶ $\sup_{r \in \mathbf{R}^+} r^2 V(r) < \infty$
- ▶ $\inf_{r \in \mathbf{R}^+} r^2 V(r) > -1/4$
- ▶ $\sup_{r \in \mathbf{R}^+} r^2 \frac{d}{dr}(rV(r)) < 1/4,$

let $P = -\Delta + V$, and let P_F be the Friedrichs extension of P . Then there exists a C such that if

$$\partial_t^2 u + P_F u = 0$$

then

$$\|u\|_{L^4(\mathbf{R}^{1+3})} \leq C \mathbf{E}_{1/2}[u].$$

⁴Indiana J of Math, 2004, but use Arxiv instead!

What needs to be done?

There are two things to be checked, before we can apply the theorem:

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- ▶ We need to check the three hypotheses on our V .

Checking V

Recall:

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$$r_0 = m \log \left(\frac{e^2}{e^2 - m^2} \right) - \frac{2m^2 - e^2}{\sqrt{e^2 - m^2}} \arctan \frac{m}{\sqrt{e^2 - m^2}}$$

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We still have nasty transcendental functions of the two variables r/e and m/e whose zeroes we need to find.

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Idea: Forget the transcendental relation between the two and treat r/e and ρ/e as independent. That gives us an algebraic problem in the three variables ρ/e , r/e and m/e .

Checking V

Recall:

$$V = \frac{2m}{\rho^3} - \frac{2e^2 + 4m^2}{\rho^4} + \frac{6me^2}{\rho^5} - \frac{2e^4}{\rho^6}$$

$$r = \rho - r_0 + m \log \left(\frac{\rho^2 - 2m\rho + e^2}{e^2 - m^2} \right) + \frac{2m^2 - e^2}{\sqrt{e^2 - m^2}} \arctan \frac{\rho - m}{\sqrt{e^2 - m^2}},$$

$$r_0 = m \log \left(\frac{e^2}{e^2 - m^2} \right) - \frac{2m^2 - e^2}{\sqrt{e^2 - m^2}} \arctan \frac{m}{\sqrt{e^2 - m^2}}$$

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Doesn't quite work as written, but something very similar does, at least for $|e| \geq 2m$.

Algebra

Eventually, after much algebra, the problem reduces to one of real algebraic geometry:

Is the curve

$$\begin{aligned} &72x^{14} - (576 + 432y^2)x^{13} + (1947 + 3552y^2 + 1152y^4)x^{12} \\ &- (3504 + 11988y^2 + 10464y^4 + 1440y^6)x^{11} \\ &+ (3452 + 20360y^2 + 38762y^4 + 15384y^6 + 720y^8)x^{10} \\ &- (1536 + 16456y^2 + 71800y^4 + 66316y^6 + 10536y^8)x^9 \\ &+ (2040y^2 + 62966y^4 + 143492y^6 + 57803y^8 + 2160y^{10})x^8 \\ &- (-4608y^2 + 8608y^4 + 153832y^6 + 154672y^8 + 21648y^{10})x^7 \\ &+ (-20100y^4 + 48272y^6 + 208760y^8 + 83120y^{10} + 2760y^{12})x^6 \\ &- (-36120y^6 + 104440y^8 + 151552y^{10} + 20824y^{12})x^5 \\ &+ (-33769y^8 + 109100y^{10} + 58958y^{12} + 1908y^{14})x^4 \\ &- (-17900y^{10} + 62848y^{12} + 11680y^{14})x^3 \\ &+ (-5530y^{12} + 20912y^{14} + 944y^{16})x^2 \\ &- (-972y^{14} + 3888y^{16})x + (-81y^{16} + 324y^{18}) = 0 \end{aligned}$$

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Algebra (Continued)

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Note that there is a classical algorithm for checking the existence of real zeroes of polynomials in one variable, the *Sturm test*.

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So now what?

For the scalar wave equation in super-extremal
Reissner-Nordström,

- ▶ energy estimates are trivial, once you figure out how to get a well-defined problem,
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There are some interesting new non-linear stability results for MBI, by Speck.

Questions?