# RELATIVE OSCILLATION THEORY FOR JACOBI MATRICES

KERSTIN AMMANN AND GERALD TESCHL

ABSTRACT. We develop relative oscillation theory for Jacobi matrices which, rather than counting the number of eigenvalues of one single matrix, counts the difference between the number of eigenvalues of two different matrices. This is done by replacing nodes of solutions associated with one matrix by weighted nodes of Wronskians of solutions of two different matrices.

### 1. INTRODUCTION

Oscillation theory for second-order differential and difference equations has a long tradition originating in the seminal work of Sturm from 1836 [9]. Since then the subject is continuously growing and many monographs have been devoted entirely to this subject. The most recent one being the monumental treatise by Agarwal, Bohner, Grace, and O'Regan [1]. One of the key results of classical oscillation theory is the fact, the k'th eigenfunction has precisely k-1 nodes (i.e., sign flips) and for a suitably chosen solution of the underlying difference equation, the number of nodes of this solutions equals the number of eigenvalues below a given value. Our aim is to add a new wrinkle to this classical result by showing that the number of weighted nodes of the Wronskian (also known as Casoratian) of two suitable solutions of two difference equations can be used to count the difference between the number of eigenvalues of the two associated Jacobi matrices.

That Wronskians are related to oscillation theory is indicated by an old paper of Leighton [7], who noted that if two solutions have a non-vanishing Wronskian, then their zeros must intertwine each other. However, it seems their real power was realized only later by Gesztesy, Simon, and Teschl in [3] with the corresponding extension to Jacobi operators given by Teschl [10]. For a pedagogical discussion we refer to the survey by Simon [8]. That these results are just the tip of the iceberg was discovered only recently by Krüger and Teschl [4], [5], [6]. Our result generalizes the main result for the case of Sturm–Liouville operators from [4] to the case of Jacobi matrices.

<sup>2000</sup> Mathematics Subject Classification. Primary 39A10, 47B36; Secondary 34C10, 34L05. Key words and phrases. Jacobi matrices, oscillation theory.

in Proceedings of the 14th International Conference on Difference Equations and Applications, M. Bohner (ed) et al., 105–115, Uğur–Bahçeşehir University Publishing Company, Istanbul, 2009.

Research supported by the Austrian Science Fund (FWF) under Grant No. Y330.

To set the stage, let us fix some real numbers a(j) < 0, b(j),  $j = 1, \dots, N-1$ and consider the Jacobi matrix

(1.1) 
$$H = \begin{pmatrix} b(1) & a(1) & 0 & 0 & 0 \\ a(1) & b(2) & \ddots & 0 & 0 \\ 0 & \ddots & \ddots & \ddots & 0 \\ 0 & 0 & a(N-1) & b(N-2) & a(N-2) \\ 0 & 0 & 0 & a(N-2) & b(N-1) \end{pmatrix}$$

in the Hilbert space  $\mathbb{C}^{N-1}$ . Furthermore, let  $s_{\pm}(z, n)$  be the solutions of the underlying difference equation (set a(0) = a(N-1) = a(N) = -1, b(N) = 0)

$$(1.2) \quad a(n)u(n+1) + b(n)u(n) + a(n-1)u(n-1) = zu(n), \qquad n = 1, \dots, N,$$

corresponding to the initial conditions

$$(1.3) s_{-}(z,0) = 0, s_{-}(z,1) = 1, s_{+}(z,N) = 0, s_{+}(z,N+1) = 1.$$

Note that  $s_{-}(\lambda, n)$  (resp.  $s_{+}(\lambda, n)$ ) will be an eigenvector of H corresponding to the eigenvalue  $\lambda \in \mathbb{R}$  if and only if  $s_{-}(\lambda, N) = 0$  (resp.  $s_{+}(\lambda, 0) = 0$ ). We will abbreviate  $s(z, n) = s_{-}(z, n)$ .

We call n a node of a solution u of (1.2) if either

(1.4) 
$$u(n) = 0$$
 or  $u(n)u(n+1) < 0$ .

We say that a node  $n_0$  of u lies between m and n if either

(1.5) 
$$m < n_0 < n$$
 or  $n_0 = m$  but  $u(m) \neq 0$ .

 $\#_{(m,n)}(u)$  denotes the number of nodes of u between m and n and  $\#(u) = \#_{(0,N)}(u)$ . Then we have the following classical result alluded to before (see e.g., [2], [11]):

**Theorem 1.1.** Let H be a Jacobi matrix and s(z, n) a corresponding solution of the underlying difference equation (1.2) corresponding to the initial condition s(z, 0) = 0. Then for every  $\lambda \in \mathbb{R}$  the number of nodes of  $s(\lambda, n)$  equals the number of eigenvalues of H below  $\lambda$ :

(1.6) 
$$\#(s(\lambda)) = \#\{E \in \sigma(H) | E < \lambda\}.$$

Here  $\sigma(H)$  denotes the spectrum of H, that is, the set of eigenvalues.

To generalize this result we will now consider two Jacobi matrices  $H_0$  and  $H_1$ associated with the coefficients  $a_0(n) = a_1(n) \equiv a(n)$  and  $b_0(n)$  respectively  $b_1(n)$ . The corresponding solutions will be denoted by  $s_{j,\pm}(n)$ , j = 0, 1, in obvious notation. Given two solutions  $u_j$  of the difference equations associated with  $H_j$  we denote by

(1.7) 
$$W_n(u_0, u_1) = a(n)(u_0(n)u_1(n+1) - u_0(n+1)u_1(n))$$

their Wronskian. As already anticipated we will relate the number of nodes of such Wronskians to the difference between the eigenvalues of  $H_0$  and  $H_1$ . Since this difference is a signed quantity, we will need to weight the nodes according to the sign of the difference between  $H_0$  and  $H_1$  as follows: Let  $u_0$ ,  $u_1$  be solutions of

the difference equations associated with  $H_0$ ,  $H_1$  and spectral parameters  $\lambda_0$ ,  $\lambda_1$ , respectively and set

$$(1.8) \qquad \#_n(u_0, u_1) = \begin{cases} \text{if } b_0(n+1) - \lambda_0 - b_1(n+1) + \lambda_1 > 0 \text{ and} \\ 1, & \text{either } W_n(u_0, u_1) W_{n+1}(u_0, u_1) < 0 \\ & \text{or } W_n(u_0, u_1) = 0 \text{ and } W_{n+1}(u_0, u_1) \neq 0, \end{cases}$$
$$(1.8) \qquad \#_n(u_0, u_1) = 0 \text{ and } W_{n+1}(u_0, u_1) \neq 0, \\ & \text{if } b_0(n+1) - \lambda_0 - b_1(n+1) + \lambda_1 < 0 \text{ and} \\ -1, & \text{either } W_n(u_0, u_1) W_{n+1}(u_0, u_1) < 0 \\ & \text{or } W_n(u_0, u_1) \neq 0 \text{ and } W_{n+1}(u_0, u_1) = 0, \\ 0, & \text{otherwise.} \end{cases}$$

Then we say the Wronskian has a weighted node at n if  $\#_n(u_0, u_1) \neq 0$ . The number of weighted nodes of the Wronskian between 0 and N is denoted as

(1.9) 
$$\#(u_0, u_1) = \sum_{j=0}^{N-1} \#_j(u_0, u_1) - \begin{cases} 0, & \text{if } W_0(u_0, u_1) \neq 0, \\ 1, & \text{if } W_0(u_0, u_1) = 0. \end{cases}$$

With this notation our main result reads

**Theorem 1.2.** Let  $H_0$ ,  $H_1$  be two Jacobi matrices with  $a_0 = a_1$  and  $s_{j,\pm}(z,n)$ , j = 0, 1, the corresponding solutions of the underlying difference equations. Then for  $\lambda_0, \lambda_1 \in \mathbb{R}$  the number of weighted nodes of  $W(s_{0,-}(\lambda_0), s_{1,+}(\lambda_1))$  equals the number of eigenvalues of  $H_1$  below  $\lambda_1$  minus the number of eigenvalues of  $H_0$  below or equal to  $\lambda_0$ :

(1.10) 
$$\#(s_{0,-}(\lambda_0), s_{1,+}(\lambda_1)) = \#(s_{0,+}(\lambda_0), s_{1,-}(\lambda_1)) = \\ = \#\{E \in \sigma(H_1) | E < \lambda_1\} - \#\{E \in \sigma(H_0) | E \le \lambda_0\}.$$

Here  $\sigma(H)$  denotes the spectrum of H, that is, the set of eigenvalues.

The proof is based on Prüfer angles to be investigated in Section 2. It will be given in Section 3.

An extension to Jacobi operators on  $\mathbb{N}$  respectively  $\mathbb{Z}$  is in preparation.

## 2. Prüfer angles

Since any nontrivial solution of (1.2) cannot vanish at two consecutive points we can introduce Prüfer variables  $(\rho_u(n), \theta_u(n))$  in the usual way (cf., e.g., [11, Chap. 4]) via

(2.1) 
$$u(n) = \rho_u(n)\sin(\theta_u(n)), \quad u(n+1) = \rho_u(n)\cos(\theta_u(n)).$$

Note that  $\rho_u(n) > 0$  for all  $n \in \mathbb{Z}$  and  $\theta_u(n)$  is only defined up to an additive integer multiple of  $2\pi$ , depending on n. For our further investigations it is essential to gain unique values for the Prüfer angle and therefore we fix  $\theta_u(0)$  and require

(2.2) 
$$\left\lceil \theta_u(n)/\pi \right\rceil \le \left\lceil \theta_u(n+1)/\pi \right\rceil \le \left\lceil \theta_u(n)/\pi \right\rceil + 1,$$

where  $\lceil x \rceil = \min\{n \in \mathbb{Z} \mid n \ge x\}$  denotes the usual ceiling function. Then the following easy result is well-known.

**Lemma 2.1.** Define  $k, \gamma, \Gamma$  via

 $\begin{array}{ll} (2.3) \quad \theta_u(n)=k\pi+\gamma, \quad \theta_u(n+1)=k\pi+\Gamma, \qquad \gamma\in(0,\pi], \ \Gamma\in(0,2\pi], \ k\in\mathbb{Z}. \end{array} \\ Then \end{array}$ 

(2.4) 
$$\gamma \in \begin{cases} (0, \frac{\pi}{2}] & iff \ n \ is \ not \ a \ node, \\ (\frac{\pi}{2}, \pi] & iff \ n \ is \ a \ node, \end{cases}$$

and

(2.5) 
$$\Gamma \in \begin{cases} (0,\pi] & \text{iff } n \text{ is not } a \text{ node,} \\ (\pi,2\pi) & \text{iff } n \text{ is } a \text{ node.} \end{cases}$$

Moreover,

(2.6) 
$$\theta_u(n) = k\pi + \frac{\pi}{2} \quad \Leftrightarrow \quad \theta_u(n+1) = (k+1)\pi.$$

As a consequence we obtain

Corollary 2.2. We have

(2.7) 
$$\left\lceil \frac{\theta_u(n+1)}{\pi} \right\rceil = \begin{cases} \left\lceil \frac{\theta_u(n)}{\pi} \right\rceil + 1 & \text{if } n \text{ is a node,} \\ \left\lceil \frac{\theta_u(n)}{\pi} \right\rceil & \text{otherwise.} \end{cases}$$

In particular, we obtain

(2.8) 
$$\#(u) = \left\lceil \frac{\theta_u(N)}{\pi} \right\rceil - \left\lfloor \frac{\theta_u(0)}{\pi} \right\rfloor - 1,$$

where  $\lfloor x \rfloor = \max\{n \in \mathbb{Z} \mid n \leq x\}$  is the usual floor function.

To find the analogous formula for the number of weighted nodes of a Wronskian we observe

(2.9) 
$$W_n(u_0, u_1) = -a(n)\rho_{u_0}(n)\rho_{u_1}(n)\sin(\Delta_{u_0, u_1}(n)),$$

where

(2.10) 
$$\Delta_{u_0,u_1}(n) = \theta_{u_1}(n) - \theta_{u_0}(n).$$

Since we can always absorb  $\lambda$  in b, we will assume  $\lambda_0 = \lambda_1 = 0$  from now on. Furthermore, note

$$(2.11) \quad W_{n+1}(u_0, u_1) - W_n(u_0, u_1) = (b_0(n+1) - b_1(n+1))u_0(n+1)u_1(n+1)$$

As a straightforward consequence of Lemma 2.1 we obtain

**Lemma 2.3.** Fix some n and let  $\theta_j(n) = k_j \pi + \gamma_j$  with  $\gamma_j \in (0, \pi]$  and  $\theta_j(n+1) = k_j \pi + \Gamma_j$  with  $\Gamma_j \in (0, 2\pi]$  for j = 0, 1. Then we have

(2.12)  $\Delta_{u_0,u_1}(n) = (k_1 - k_0)\pi + \gamma_1 - \gamma_0$  and  $\Delta_{u_0,u_1}(n+1) = (k_1 - k_0)\pi + \Gamma_1 - \Gamma_0$ , where

(1): either  $u_0$  and  $u_1$  have a node at n or both do not have a node at n, then

(2.13) 
$$\gamma_1 - \gamma_0 \in (-\frac{\pi}{2}, \frac{\pi}{2})$$
 and  $\Gamma_1 - \Gamma_0 \in (-\pi, \pi).$ 

(2):  $u_1$  has no node at n, but  $u_0$  has a node at n, then

(2.14) 
$$\gamma_1 - \gamma_0 \in (-\pi, 0)$$
 and  $\Gamma_1 - \Gamma_0 \in (-2\pi, 0).$ 

(3):  $u_1$  has a node at n, but  $u_0$  has no node at n, then

(2.15)  $\gamma_1 - \gamma_0 \in (0,\pi) \quad and \quad \Gamma_1 - \Gamma_0 \in (0,2\pi).$ 

Now we are able to show

**Lemma 2.4.** Fix some n. Then, if  $b_0(n+1) \ge b_1(n+1)$ , we have

(2.16) 
$$\left[\Delta_{u_0,u_1}(n)/\pi\right] \le \left[\Delta_{u_0,u_1}(n+1)/\pi\right] \le \left[\Delta_{u_0,u_1}(n)/\pi\right] + 1$$

and if  $b_0(n+1) \le b_1(n+1)$ , we have

(2.17) 
$$\lceil \Delta_{u_0,u_1}(n)/\pi \rceil - 1 \le \lceil \Delta_{u_0,u_1}(n+1)/\pi \rceil \le \lceil \Delta_{u_0,u_1}(n)/\pi \rceil.$$

*Proof.* We will use the notation from Lemma 2.3 where we assume  $k_0 = k_1 = 0$  without loss of generality. In particular, Lemma 2.3 implies

$$\left[\Delta_{u_0,u_1}(n)/\pi\right] - 1 \le \left[\Delta_{u_0,u_1}(n+1)/\pi\right] \le \left[\Delta_{u_0,u_1}(n)/\pi\right] + 1.$$

Hence, to show (2.16) there are two cases to exclude. Namely, (i)  $\Delta_{u_0,u_1}(n) \in (0, \frac{\pi}{2})$ ,  $\Delta_{u_0,u_1}(n+1) \in (-\pi, 0]$  (from case (1)) and (ii)  $\Delta_{u_0,u_1}(n) \in (-\pi, 0)$ ,  $\Delta_{u_0,u_1}(n+1) \in (-2\pi, -\pi]$  (from case (2)). But in case (i) we obtain a contradiction from (2.11):

$$\underbrace{W_{n+1}(u_0, u_1)}_{\leq 0} = \underbrace{W_n(u_0, u_1)}_{>0} + \underbrace{(b_0(n+1) - b_1(n+1))}_{\geq 0} \underbrace{u_0(n+1)u_1(n+1)}_{\geq 0}.$$

Similarly, in case (ii) equation (2.11) implies

$$\underbrace{W_{n+1}(u_0, u_1)}_{\geq 0} = \underbrace{W_n(u_0, u_1)}_{< 0} + \underbrace{(b_0(n+1) - b_1(n+1))}_{\geq 0} \underbrace{u_0(n+1)u_1(n+1)}_{\leq 0}.$$

Equation (2.17) can be established in a similar fashion.

**Lemma 2.5.** Let  $n \in \mathbb{Z}$ , then

(1): 
$$W_n(u_0, u_1) = W_{n+1}(u_0, u_1) = 0 \text{ or } W_n(u_0, u_1) W_{n+1}(u_0, u_1) > 0 \text{ implies}$$
  
(2.18)  $\left\lceil \frac{\Delta_{u_0, u_1}(n+1)}{\pi} \right\rceil = \left\lceil \frac{\Delta_{u_0, u_1}(n)}{\pi} \right\rceil.$ 

(2):  $W_n(u_0, u_1)W_{n+1}(u_0, u_1) < 0$  implies

(2.19) 
$$\left\lceil \frac{\Delta_{u_0,u_1}(n+1)}{\pi} \right\rceil = \begin{cases} \left\lceil \frac{\Delta_{u_0,u_1}(n)}{\pi} \right\rceil + 1, & \text{if } b_0(n+1) > b_1(n+1), \\ \left\lceil \frac{\Delta_{u_0,u_1}(n)}{\pi} \right\rceil - 1, & \text{if } b_0(n+1) < b_1(n+1). \end{cases}$$

(3):  $W_n(u_0, u_1) = 0$  and  $W_{n+1}(u_0, u_1) \neq 0$  implies

(2.20) 
$$\lceil \frac{\Delta_{u_0,u_1}(n+1)}{\pi} \rceil = \begin{cases} \lceil \frac{\Delta_{u_0,u_1}(n)}{\pi} \rceil + 1, & \text{if } b_0(n+1) > b_1(n+1), \\ \lceil \frac{\Delta_{u_0,u_1}(n)}{\pi} \rceil, & \text{if } b_0(n+1) < b_1(n+1). \end{cases}$$

(4):  $W_n(u_0, u_1) \neq 0$  and  $W_{n+1}(u_0, u_1) = 0$  implies

(2.21) 
$$\lceil \frac{\Delta_{u_0,u_1}(n+1)}{\pi} \rceil = \begin{cases} \lceil \frac{\Delta_{u_0,u_1}(n)}{\pi} \rceil, & \text{if } b_0(n+1) > b_1(n+1), \\ \lceil \frac{\Delta_{u_0,u_1}(n)}{\pi} \rceil - 1, & \text{if } b_0(n+1) < b_1(n+1). \end{cases}$$

Note that in the cases (2)–(4) we necessarily have  $b_0(n+1) \neq b_1(n+1)$ .

*Proof.* We will use the notation from Lemma 2.3 where we assume  $k_0 = k_1 = 0$  without loss of generality. Moreover, interchanging  $u_0$  and  $u_1$  using  $\Delta_{u_1,u_0}(n) = -\Delta_{u_0,u_1}(n)$  and

$$\lceil -x \rceil = \begin{cases} -\lceil x \rceil & \text{if } x \in \mathbb{Z}, \\ -\lceil x \rceil + 1 & \text{otherwise,} \end{cases}$$

we see that it suffices to show one case  $b_0(n+1) \ge b_1(n+1)$  or  $b_0(n+1) \le b_1(n+1)$ .

Suppose  $W_n(u_0, u_1) = W_{n+1}(u_0, u_1) = 0$  and  $W_n(u_0, u_1)W_{n+1}(u_0, u_1) > 0$  do not hold, then by (2.11) we have

$$W_{n+1}(u_0, u_1) - W_n(u_0, u_1) = (b_0(n+1) - b_1(n+1))u_0(n+1)u_1(n+1) \neq 0$$

and hence  $b_0(n+1) \neq b_1(n+1)$ .

(1) and (2). Suppose 
$$W_n(u_0, u_1) = W_{n+1}(u_0, u_1) = 0$$
, then by (2.9) we infer

$$\sin(\Delta_{u_0,u_1}(n)) = \sin(\gamma_1 - \gamma_0) = 0, \quad \sin(\Delta_{u_0,u_1}(n+1)) = \sin(\Gamma_1 - \Gamma_0) = 0,$$

where  $\gamma_0, \gamma_1 \in (0, \pi]$ . Thus  $\gamma_0 = \gamma_1$  and we have case (1) of Lemma 2.3 which implies  $\Gamma_1 - \Gamma_0 \in (-\pi, \pi)$  and we conclude  $\Gamma_1 - \Gamma_0 = 0$ . In summary,  $\Delta_{u_0, u_1}(n) = \Delta_{u_0, u_1}(n+1) = 0$  as claimed.

Next suppose  $W_n(u_0, u_1)W_{n+1}(u_0, u_1) \neq 0$ , then by (2.9) the sign of the Wronskian at *n* equals the sign of  $\sin(\Delta_{u_0,u_1}(n))$  and hence (2.16) respectively (2.17) finish the proof of case (1) and (2).

(3). By (2.9) we conclude  $\Delta_{u_0,u_1}(n) = \gamma_1 - \gamma_0 \equiv 0 \mod \pi$ , where  $\gamma_0, \gamma_1 \in (0,\pi]$ and thus  $\gamma_1 - \gamma_0 = 0$ . So we have case (1) of Lemma 2.3 and hence  $\Delta_{u_0,u_1}(n+1) = \Gamma_1 - \Gamma_0 \in (-\pi,\pi)$ . That is,

$$\lceil \Delta_{u_0,u_1}(n)/\pi \rceil \le \lceil \Delta_{u_0,u_1}(n+1)/\pi \rceil \le \lceil \Delta_{u_0,u_1}(n)/\pi \rceil + 1$$

and (2.17) finishes the proof of case (3) for  $b_0(n+1) < b_1(n+1)$ .

(4). By (2.9) we have  $\Delta_{u_0,u_1}(n+1) = \Gamma_1 - \Gamma_0 \equiv 0 \mod \pi$  and Lemma 2.3 leaves us with the following possibilities

$$\begin{array}{lll} \text{(a)} & \Delta_{u_0,u_1}(n) \in (-\frac{\pi}{2}, \frac{\pi}{2}) & \text{and} & \Delta_{u_0,u_1}(n+1) = 0, \\ \text{(b)} & \Delta_{u_0,u_1}(n) \in (-\pi, 0) & \text{and} & \Delta_{u_0,u_1}(n+1) = -\pi, \\ \text{(c)} & \Delta_{u_0,u_1}(n) \in (0, \pi) & \text{and} & \Delta_{u_0,u_1}(n+1) = \pi. \end{array}$$

and (2.16) shows (4) if  $b_0(n+1) > b_1(n+1)$ .

As a consequence we obtain the desired formula

(2.22) 
$$\#(u_0, u_1) = \lceil \Delta_{u_0, u_1}(N) / \pi \rceil - \lfloor \Delta_{u_0, u_1}(0) / \pi \rfloor - 1.$$

### 3. Proof of the main theorem

Our strategy will be to interpolate between  $H_0$  and  $H_1$  using  $H_{\varepsilon} = (1 - \varepsilon)H_0 + \varepsilon H_1$ , that is,  $a_{\varepsilon}(n) = a(n)$  and  $b_{\varepsilon}(n) = (1 - \varepsilon)b_0(n) + \varepsilon b_1(n)$ . If  $u_{\varepsilon}$  is a solution of the difference equation corresponding to  $H_{\varepsilon}$ , then the corresponding Prüfer angles satisfy

(3.1) 
$$\dot{\theta}_{\varepsilon}(n) = -\frac{W_n(u_{\varepsilon}, \dot{u}_{\varepsilon})}{a(n)\rho_{\varepsilon}^2(n)},$$

where the dot denotes a derivative with respect to  $\varepsilon$ .

Lemma 3.1. We have

(3.2) 
$$W_n(s_{\varepsilon,\pm}(z), \dot{s}_{\varepsilon,\pm}(z)) = \begin{cases} -\sum_{m=n+1}^N (b_0(m) - b_1(m)) s_{\varepsilon,\pm}(z,m)^2, \\ \sum_{m=1}^n (b_0(m) - b_1(m)) s_{\varepsilon,-}(z,m)^2. \end{cases}$$

*Proof.* Summing (2.11) we obtain

$$W_n(s_{\varepsilon,\pm}(z), s_{\tilde{\varepsilon},\pm}(z)) = (\tilde{\varepsilon} - \varepsilon) \begin{cases} -\sum_{m=n+1}^N (b_0(m) - b_1(m)) s_{\varepsilon,+}(z,m) s_{\tilde{\varepsilon},+}(z,m), \\ \sum_{m=1}^n (b_0(m) - b_1(m)) s_{\varepsilon,-}(z,m) s_{\tilde{\varepsilon},-}(z,m). \end{cases}$$

Now use this to evaluate the limit

$$\lim_{\tilde{\varepsilon}\to\varepsilon} W_n\Big(s_{\varepsilon,\pm}(z), \frac{s_{\varepsilon,\pm}(z)-s_{\tilde{\varepsilon},\pm}(z)}{\varepsilon-\tilde{\varepsilon}}\Big).$$

Denoting the Prüfer angles of  $s_{\varepsilon,\pm}(\lambda, n)$  by  $\theta_{\varepsilon,\pm}(\lambda, n)$ , this result implies for  $b_0 - b_1 \ge 0$ ,

(3.3) 
$$\dot{\theta}_{\varepsilon,+}(\lambda,n) = \frac{\sum_{m=n+1}^{N} (b_0(m) - b_1(m)) s_{\varepsilon,+}(z,m)^2}{a(n)\rho_{\varepsilon,+}(\lambda,n)^2} \le 0,$$
$$\dot{\theta}_{\varepsilon,-}(\lambda,n) = -\frac{\sum_{m=1}^{n} (b_0(m) - b_1(m)) s_{\varepsilon,-}(z,m)^2}{a(n)\rho_{\varepsilon,-}(\lambda,n)^2} \ge 0.$$

Furthermore, we have the following result from classical perturbation theory. We add a simple direct proof for convenience of the reader.

**Lemma 3.2.** Suppose  $b_0 - b_1 \ge 0$  (resp.  $b_0 - b_1 \le 0$ ). Then the eigenvalues of  $H_{\varepsilon}$  are analytic functions with respect to  $\varepsilon$  and they are decreasing (resp. increasing).

*Proof.* First of all the Prüfer angles  $\theta_{\varepsilon,\pm}(\lambda, n)$  are analytic with respect to  $\varepsilon$  since  $s_{\varepsilon,\pm}(\lambda, n)$  is a polynomial with respect to  $\varepsilon$ . Moreover,  $\lambda \in \sigma(H_{\varepsilon})$  is equivalent to  $\theta_{\varepsilon,+}(\lambda, 0) \equiv 0 \mod \pi$  (resp.  $\theta_{\varepsilon,-}(\lambda, N) \equiv 0 \mod \pi$ ) and monotonicity follows from (3.3).

In particular, this implies that  $P(H_{\varepsilon}) = \#\{E \in \sigma(H_{\varepsilon}) | E < \lambda\}$  is continuous from below (resp. above) in  $\varepsilon$  if  $b_0 - b_1 \ge 0$  (resp.  $b_0 - b_1 \le 0$ ).

Now we are ready for the

Proof of Theorem 1.2. It suffices to prove the result for  $\#(s_{0,+}(\lambda_0), s_{1,-}(\lambda_1))$ , where we can assume  $\lambda_0 = \lambda_1 = 0$  without restriction and set  $s_{\varepsilon,\pm}(n) = s_{\varepsilon,\pm}(0,n)$  for notational convenience. We split  $b_0 - b_1$  according to

$$b_0 - b_1 = b_+ - b_-, \qquad b_+, b_- \ge 0,$$

and introduce the operator  $H_{-} = H_0 - b_{-}$ . Then  $H_{-}$  is a negative perturbation of  $H_0$  and  $H_1$  is a positive perturbation of  $H_{-}$ .

Furthermore, define  $H_{\varepsilon}$  by

$$H_{\varepsilon} = \begin{cases} H_0 + 2\varepsilon (H_- - H_0), & \varepsilon \in [0, 1/2], \\ H_- + 2(\varepsilon - 1/2)(H_1 - H_-), & \varepsilon \in [1/2, 1]. \end{cases}$$

Let us look at (using (2.22))

$$Q(\varepsilon) = \#(s_{0,+}, s_{\varepsilon,-}) = \lceil \Delta_{\varepsilon}(N)/\pi \rceil - \lfloor \Delta_{\varepsilon}(0)/\pi \rfloor - 1, \quad \Delta_{\varepsilon}(n) = \Delta_{s_{0,+}, s_{\varepsilon,-}}(n)$$

and consider  $\varepsilon \in [0, 1/2]$ . At the left boundary  $\Delta_{\varepsilon}(0)$  remains constant whereas at the right boundary  $\Delta_{\varepsilon}(N)$  is increasing by (3.3). Moreover, it hits a multiple of  $\pi$ whenever  $0 \in \sigma(H_{\varepsilon})$ . So  $Q(\varepsilon)$  is a piecewise constant function which is continuous from below and jumps by one whenever  $0 \in \sigma(H_{\varepsilon})$ . By Lemma 3.2 the same is true for

$$P(\varepsilon) = \#\{E \in \sigma(H_{\varepsilon}) | E < 0\} - \#\{E \in \sigma(H_0) | E \le 0\}$$

and since we have Q(0) = P(0), we conclude  $Q(\varepsilon) = P(\varepsilon)$  for all  $\varepsilon \in [0, 1/2]$ . To see the remaining case  $\varepsilon = [1/2, 1]$ , simply replace increasing by decreasing and continuous from below by continuous from above.

Acknowledgments. We thank H. Krüger for several valuable discussions. Furthermore, G.T. would like to thank all organizers of the 14th International Conference on Difference Equations and Applications (ICDEA), Istanbul, July 2008, and especially Martin Bohner and Mehmet Ünal, for their kind invitation and the stimulating atmosphere during the meeting.

#### References

- R. P. Agarwal, M. Bohner, S. R. Grace, D. O'Regan, Discrete Oscillation Theory, Hindawi Publishing Corp., New York, 2005.
- [2] F. Atkinson, Discrete and Continuous Boundary Problems, Academic Press, New York, 1964.
- [3] F. Gesztesy, B. Simon, and G. Teschl, Zeros of the Wronskian and renormalized oscillation Theory, Am. J. Math. 118 571–594 (1996).
- [4] H. Krüger and G. Teschl, Relative oscillation theory, weighted zeros of the Wronskian, and the spectral shift function, Comm. Math. Phys. 287:2, 613–640 (2009).
- [5] H. Krüger and G. Teschl, Relative oscillation theory for Sturm-Liouville operators extended, J. Funct. Anal. 254-6, 1702–1720 (2008).
- [6] H. Krüger and G. Teschl, Effective Pr
  üfer angles and relative oscillation criteria, J. Diff. Eq. 245, 3823–3848 (2008).
- [7] W. Leighton, On self-adjoint differential equations of second order, J. London Math. Soc. 27, 37–47 (1952).
- [8] B. Simon, Sturm oscillation and comparison theorems, in Sturm-Liouville Theory: Past and Present (eds. W. Amrein, A. Hinz and D. Pearson), 29–43, Birkhäuser, Basel, 2005.
- [9] J.C.F. Sturm, Mémoire sur les équations différentielles linéaires du second ordre, J. Math. Pures Appl., 1, 106–186 (1836).
- G. Teschl, Oscillation theory and renormalized oscillation theory for Jacobi operators, J. Diff. Eqs. 129, 532–558 (1996).
- [11] G. Teschl, Jacobi Operators and Completely Integrable Nonlinear Lattices, Math. Surv. and Mon. 72, Amer. Math. Soc., Rhode Island, 2000.

Faculty of Mathematics, University of Vienna, Nordbergstrasse 15, 1090 Wien, Austria

*E-mail address*: Kerstin.Ammann@univie.ac.at *URL*: http://www.mat.univie.ac.at/~kerstin/

Faculty of Mathematics, University of Vienna, Nordbergstrasse 15, 1090 Wien, Austria, and International Erwin Schrödinger Institute for Mathematical Physics, Boltzmanngasse 9, 1090 Wien, Austria

*E-mail address*: Gerald.Teschl@univie.ac.at *URL*: http://www.mat.univie.ac.at/~gerald/