

# Holography of BGG-Solutions

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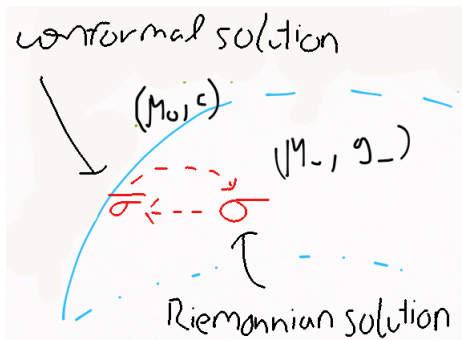
Joint work with Travis Willse (University of Vienna)

# Introductory picture: Holography of solutions

Let  $M = M_0 \cup M_-$  be a (real-analytic) *Poincaré-Einstein* manifold:

- $(M_0, \mathbf{c})$  is a conformal  $n$ -dimensional space.
- $(M_-, g_-)$  an Einstein metric with  $\text{Ric}(g_-) = -ng_-$ .
- $(M_0, \mathbf{c})$  is the *conformal infinity* of  $(M_-, g_-)$ .

Our goal is to relate overdetermined differential equations on  $M_-$  with corresponding equations on the conformal infinity  $M_0$ :



# Conformally invariant overdetermined Equations on $(M_0, \mathbf{c})$

Conformally invariant overdetermined equations on the boundary  $(M_0, \mathbf{c})$  are governed by *first BGG-operators*

$$\bar{\Theta}_0 : \Gamma(\bar{\mathcal{H}}_0) \rightarrow \Gamma(\bar{\mathcal{H}}_1).$$

Examples are:

- **Conformally-Einstein equation:** The equation governing the existence of an Einstein-metric in the conformal class:

$$\text{tf}(\bar{D}_a \bar{D}_b \sigma + \bar{P}_{ab} \sigma) = 0.$$

- **Conformal Killing form equation:**

$$\bar{D}_c \bar{\varphi}_{a_1 \dots a_k} = \bar{D}_{[a_0} \bar{\varphi}_{a_1 \dots a_k]} + \frac{k}{n - k + 1} \bar{g}_{c[a_1} \bar{g}^{pq} \bar{D}_{|p} \bar{\varphi}_{q|a_2 \dots a_k]}$$

- **Twistor spinor equation:**

$$\bar{D}_c \bar{\chi} + \frac{1}{n} \bar{\gamma}_c \bar{\not{D}} \bar{\chi} = 0.$$

# Projectively invariant overdetermined Equations on $(M_-, g_-)$

We regard overdetermined equations on the Riemannian structure  $(M, g_-)$  which are projectively invariant and are therefore governed by *first BGG-operators*  $\Theta_0 : \Gamma(\mathcal{H}_0) \rightarrow \Gamma(\mathcal{H}_1)$  of the projective structure  $\mathbf{p}$  spanned by the Levi-Civita covariant derivative  $D$  of  $g_-$ . Examples are:

- **Projectively-Ricci-flat equation:** The equation governing the existence of a Ricci-flat affine connection in the projective class of the Levi-Civita-connection  $D$ :

$$D_\alpha D_\beta \sigma + \sigma P_{\alpha\beta} = 0$$

- **Killing form equation:**

$$D_\gamma \varphi_{\alpha_1 \dots \alpha_k} = D_{[\alpha_0} \varphi_{\alpha_1 \dots \alpha_k]}$$

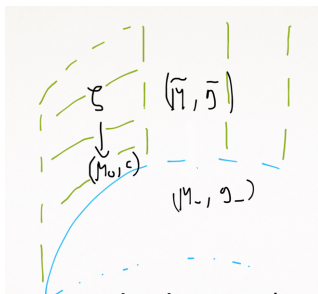
- **Killing spinor equation:**

$$D_\alpha \chi = \lambda \gamma_\alpha \chi \quad \text{for some } \lambda \in \mathbb{R}.$$

# Ambient approach to holography

The conformal structure  $\mathbf{c}$  can be understood as the ray-bundle  $\mathcal{C}$  that consists of all metrics in the given conformal class.

For  $n$  odd and  $\mathbf{c}$  real-analytic, the *Fefferman-Graham ambient metric*  $\tilde{g}$  is a Ricci-flat signature  $(n+1, 1)$  metric on an  $n+2$  dimensional *ambient space*  $\tilde{M} = \mathcal{C} \times (-1, 1)$ .



Since the ambient metric construction is natural, we obtain a (first, trivial) categorical holographic correspondence between symmetries:

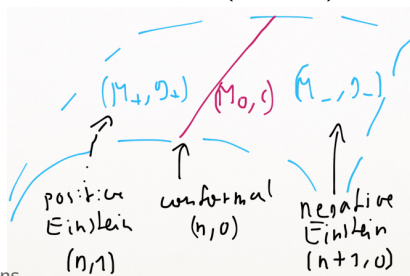
conformal Killing field  $\bar{\xi}$  on  $M_0 \Leftrightarrow$  Killing field  $\xi$  of  $g_-$ .

# Klein-Einstein

Instead of working with Poincaré-Einstein structures, it will be useful for our purposes to work with the closely related class of *Klein-Einstein* structures, which are more directly related to Fefferman-Graham ambient metrics: (Fefferman-Graham, 2011; Čap-Gover-H., 2012)

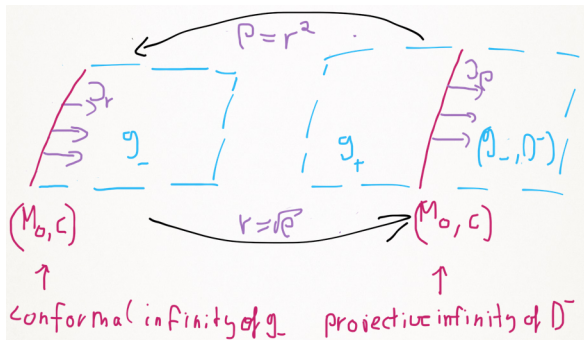
Let  $M = M_+ \cup M_0 \cup M_-$  be a (real-analytic) *Klein-Einstein* manifold:

- $(M_0, \mathbf{c})$  is a conformal  $n$ -dimensional space.
- $(M_+, g_+)$  an Einstein metric with  $\text{Ric}(g_+) = ng_+$ .
- $(M_-, g_-)$  an Einstein metric with  $\text{Ric}(g_-) = -ng_-$ .
- $(M_0, \mathbf{c})$  is the *projective infinity* of  $(M_-, g_-)$ .



# Relationship between Poincaré-Einstein and Klein-Einstein

One has that *even* Poincaré-Einstein structures correspond to Klein-Einstein structures via a “change of variables” (Fefferman-Graham, 2011).



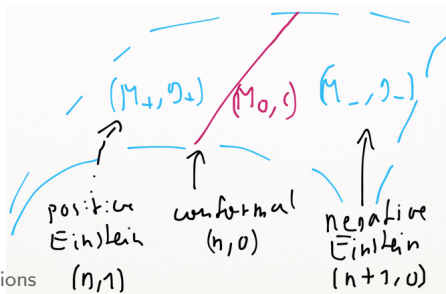
# Klein-Einstein manifolds as projective structures

Let  $M$  be an  $n+1$ -dimensional manifold endowed with a projective structure  $\mathbf{p}$  whose *normal projective tractor connection*  $\nabla^T$  preserves a signature  $(n+1, 1)$  tractor metric  $\mathbf{g}$ . Equivalently, the *projective holonomy* is reduced to

$$\text{Hol}(\mathbf{p}) \subseteq \text{SO}(n+1, 1) \subseteq \text{SL}(n+2).$$

Via the general *curved orbit decomposition* ([Čap-Gover-H], 2014) for holonomy reduced structures, one obtains (generically) a decomposition of  $M$  with properties as above into *curved*  $\text{SO}(n+1, 1)$ -orbits

$$(M_+, g_+) \cup (M_0, \mathbf{c}) \cup (M_-, g_-):$$



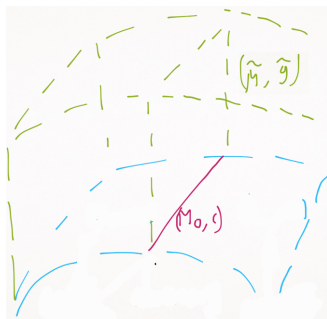


# Cone construction over a Klein-Einstein manifold

*Thomas cone construction:* The projective  $n + 1$ -dimensional structure  $(M, \mathbf{p})$  can be equivalently encoded as a canonical Ricci-flat connection  $\tilde{\nabla}$  on an  $n + 2$ -dimensional cone  $\tilde{M}$  over  $M$ .

*Metric cone:* The (sub-)cone  $\mathcal{C} \subset \tilde{M}$  over the conformal piece  $M_0 \subset M$  can be identified with the *metric cone* formed by the conformal class of metrics  $\mathbf{c}$ .

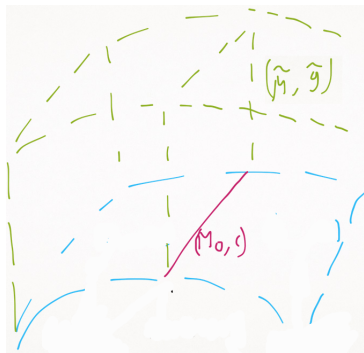
*Metric on  $\tilde{M}$ :* The signature  $(n + 1, 1)$  parallel tractor metric  $\mathbf{g}$  gives rise to a  $\tilde{\nabla}$ -parallel metric  $\tilde{g}$  of the same signature on  $\tilde{M}$ .



# Fefferman-Graham ambient space

The Thomas cone over a Klein-Einstein structure thus provides a concrete realisation of:

- A Fefferman-Graham ambient space  $\tilde{M}$  with
- Ricci-flat,  $\tilde{\nabla}$ -parallel Fefferman-Graham ambient metric  $\tilde{g}$ .



The Klein-Einstein structure  $(M, \mathbf{p})$  is an equivalent description of an (abstract) ambient space  $(\tilde{M}, \tilde{g})$ .

# Holonomies

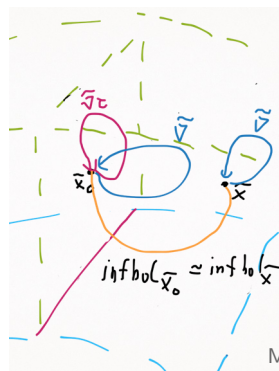
It was shown in (Čap-Gover-Graham-H. 2015) that 'infinitesimal ambient holonomy = infinitesimal conformal holonomy', which can be expressed in our (projective Klein-Einstein, 'abstract ambient' - setting) as:

$$\text{inf.hol.}_{x_0}^{M_0}(\bar{\nabla}^{conf}) = \text{inf.hol.}_{x_0}^M(\nabla^{proj}) \subset \mathfrak{so}(n+1, 1)$$

at each point  $x_0 \in M_0$ . This infinitesimal holonomy subalgebra is naturally isomorphic to infinitesimal ambient holonomy  $\text{inf.hol.}_{x_0}^{\tilde{M}}(\tilde{\nabla})$ .

Since our structure is real analytic:

- Infinitesimal holonomy is already the full (Lie algebraic) holonomy, which generates (at least locally at every point) the full holonomy.
- In particular, infinitesimal holonomy is the same (or naturally isomorphic) at every point.



## Relationships between BGG-Solutions on $M_0, M_-$

Let  $V$  be an (irreducible)  $SL(n+2)$ -representation and  $\bar{V} \subset V$  a  $SO(n+1, 1)$ -submodule. Examples are:

- $\bar{V} = V = \mathbb{R}^{n+2}$
- $\bar{V} = \mathfrak{so}(n+1, 1), V = \mathfrak{sl}(n+2)$
- $\bar{V} = V = \Lambda^{k+1}\mathbb{R}^{n+2}$
- $\bar{V} = S_0^2\mathbb{R}^{n+2}, V = S^2\mathbb{R}^{n+2}$ .

By general tractor machinery and holonomy principles:

- $\bar{V}$  gives rise to a conformal tractor bundle  $\bar{\mathcal{V}}$  endowed with normal tractor connection  $\bar{\nabla}^{\bar{\mathcal{V}}}$  for  $(M_0, \mathbf{c})$ , defined along  $M_0$ .
- $V$  gives rise to a projective tractor bundle  $\mathcal{V}$  endowed with normal tractor connection  $\nabla^{\mathcal{V}}$  on  $(M, \mathbf{p})$
- Since  $(M, \mathbf{p})$  has holonomy reduced to  $SO(n+1, 1)$ ,  $\bar{\mathcal{V}}$  is globally defined on  $M$ , and then respective projective/conformal tractor connections on this bundle agree along  $M_0$  (Čap-Gover-H., 2012).

Via the general principles of the BGG-machinery (Čap-Slovak-Souček, 2001):

- $(\bar{\mathcal{V}}, \bar{\nabla}^{\bar{\mathcal{V}}})$  gives rise to a conformally invariant (first) BGG-operator  $\bar{\Theta}_0$ .
- $(\mathcal{V}, \nabla^{\mathcal{V}})$  gives rise to a projectively invariant (first) BGG-operator  $\Theta_0$ .
- *Normal solutions* of the first BGG-operator  $\bar{\Theta}_0$  resp.  $\Theta_0$  are defined as parallel sections of  $\bar{\mathcal{V}}$  resp.  $\mathcal{V}$ :

$$\begin{array}{c} \nabla^{\mathcal{V}} - \text{parallel sections} \\ \Pi_0 \downarrow \curvearrowright L_0 \\ \{\text{normal solutions}\} \subset \ker \Theta_0^{\mathcal{V}} \end{array}$$

- For general solutions, one can employ the HSSŠ-'prolongation connection'  $\hat{\nabla}^{\mathcal{V}}$  (H.-Somberg-Souček-Šilhan, 2012):

$$\begin{array}{c} \hat{\nabla}^{\mathcal{V}} - \text{parallel sections} \\ \Pi_0 \downarrow \curvearrowright L_0 \\ \ker \Theta_0^{\mathcal{V}} \end{array}$$

# General Restriction & Extension for Normal BGG-Solutions

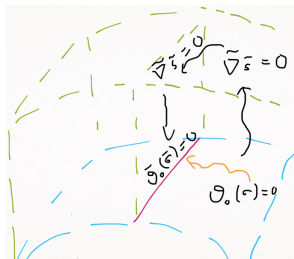
Since the connections  $\bar{\nabla}^{\bar{V}}$  resp.  $\nabla^V$  are induced from the normal conformal resp. projective tractor connections, infinitesimal holonomies agree again:

$$\text{inf.hol.}_{x_0}^{M_0}(\bar{\nabla}^{\bar{V}}) \cong \text{inf.hol.}_x^{M_-}(\nabla^V) \quad (x_0 \in M_0, x \in M_-).$$

It follows in particular that

$\{\bar{\nabla}^{\bar{V}}$ -parallel sections on  $M_0\} \cong \{\nabla^V$ -parallel sections of type  $\bar{V}$  on  $M_-\}$

and thus  $\ker_{\text{nor}} \bar{\Theta}_0^{\bar{V}} \cong \ker_{\text{nor}, \bar{V}} \Theta_0^V$ , and we have 1 : 1-correspondence between normal solutions of type  $\bar{V}$  on the interior and normal solutions on the boundary.



# Examples

For some very simple classes of BGG-equations, every solution is normal, and in particular the correspondences between solutions can be considered:

$\bar{V}$	$\tilde{M}$	$M_0$	$M_-$
$\mathbb{R}^{n+2}$	Parallel field	Einstein metric $\bar{g}_E$ in $\mathbf{c}$	Ricci-flat affine connection projectively equivalent to $D^{g_-}$
$\Delta^{n+1,1}$	Parallel spinor	twistor spinor $\bar{\chi}$	Killing spinor $\chi$

We next discuss the case  $\bar{V} = V = \Lambda^{k+1}$ , which is the first interesting case of a BGG-operator with (generally) non-normal solutions. The corresponding solutions are *conformal Killing forms* on  $M_0$  resp. Killing forms on  $M_0$ .

# Prolongations of Projective / Conformal Killing Forms

- According to (H.-Somberg-Souček-Šilhan, 2012), one has

$$\{\text{Killing forms}\} \cong \ker \Theta_0 \cong \{\hat{\nabla}\text{-parallel tractors}\}$$

where  $\hat{\nabla} = \nabla + \Psi$  is the *prolongation connection* of this BGG-equation with  $\Psi$  the *prolongation modification* of the normal tractor connection.

$\rightsquigarrow$  *Simple explicit formula available.*

- Similarly, one has

$$\{\text{Conformal Killing forms}\} \cong \ker \bar{\Theta}_0 \cong \{\hat{\bar{\nabla}}\text{-parallel tractors}\},$$

where  $\hat{\bar{\nabla}} = \bar{\nabla} + \bar{\Psi}$  is the *prolongation connection* of this BGG-equation with  $\bar{\Psi}$  the *prolongation modification* of the normal tractor connection.

$\rightsquigarrow$  *Explicit formula available.* (Gover-Šilhan 2008; HSSŠ, 2012)



# Comparison of projective and conformal prolongation connections

Since the projective prolongation connection  $\hat{\nabla}$  is defined on all of  $M$ , we can in particular restrict it to a linear connection  $\hat{\nabla}_{M_0}$  on conformal tractors in  $\Lambda^{k+1}\bar{\mathcal{T}}$  along  $M_0$ .

A conformal tractor  $\bar{s} \in \Lambda^{k+1}\bar{\mathcal{T}}$  is parallel with respect to the restricted, projective prolongation connection  $\hat{\nabla}_{M_0}$  if and only if the following conditions hold:

$$\bar{\Theta}_0(\bar{\varphi}) = 0 \quad (\text{cKf})$$

$$C_{c[a_2}{}^{pq}\bar{\sigma}_{|pq|a_3\dots a_k]} = 0, \dots \quad (\text{I})$$

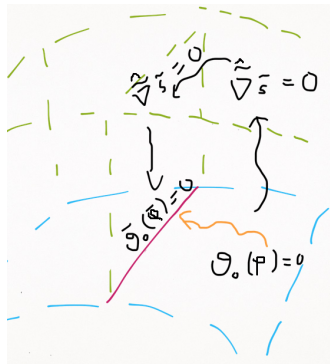
- This is the conformal Killing equation cKF together with an additional integrability condition (I).
- In particular, the restriction of the projective prolongation connection  $\hat{\nabla}$  to  $M_0$  does not coincide with the conformal prolongation connection  $\hat{\nabla}$ .

# Restriction process: $\varphi \rightsquigarrow \bar{\varphi}$

The projective prolongation connection  $\hat{\nabla} = \nabla + \Psi$  of the (projective) Killing form equation on  $M$  can be equivalently described by a suitable/corresponding 'ambient prolongation connection'  $\hat{\hat{\nabla}}$ .

Let  $\varphi_{\alpha_1 \dots \alpha_k}$  be a Killing form of  $g_-$  on  $M_-$ .

- 1 One regards the corresponding tractor as a suitably 'ambient prolongation connection'  $\hat{\hat{\nabla}}$ -parallel (multi-vector-)field  $\tilde{s}$ ,
- 2 restricts this field to  $\mathcal{C}$ ,
- 3 and obtains (via the above comparison of prolongation connections) a conformal Killing form  $\bar{\varphi}_{a_1 \dots a_k}$  on  $M_0$  satisfying additional integrability conditions (I).



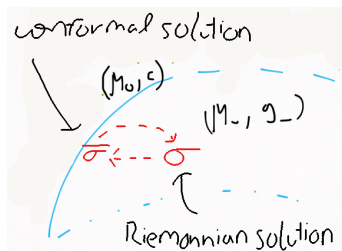
## Extension process: $\bar{\phi} \rightsquigarrow \phi$

Conversely, let  $\bar{\varphi}_{a_1 \dots a_k}$  be a conformal Killing form on  $M_0$  satisfying (I).

- 1 Let  $\tilde{s} \in \Gamma_{\mathcal{C}}(\Lambda^{k+1} T\tilde{M})$  be the corresponding ambient field defined along  $\mathcal{C}$ .
- 2 By the above comparison of prolongation connections,  $\tilde{s}$  is parallel with respect to the 'ambient prolongation connection'  $\hat{\nabla}$  along  $\mathcal{C}$  in  $\mathcal{C}$ -directions.
- 3 It is shown via an infinitesimal holonomy-computation that one obtains a global  $\hat{\nabla}$ -parallel ambient field  $\tilde{s}$ .
- 4 In particular, since  $\hat{\nabla}$  governs the (projective) Killing form equation globally on  $M$ , we obtain (via restriction), a Killing form  $\varphi_{\alpha_1 \dots \alpha_k}$  on  $M_-$ .

# Summary of BGG-Relationships

$M_0$	$M_-$
Einstein metric $\bar{g}_E$ in $\mathfrak{c}$	Ricci-flat affine connection projectively equivalent to $D^{g_-}$
conformal Killing form $\bar{\varphi}$ satisfying (I)	Killing form $\varphi$
twistor spinor $\bar{\chi}$	Killing spinor $\chi$



- For  $k = 1$  (I) is always satisfied and we recover the (categorical) relationship between conformal Killing fields  $\bar{\xi}$  on  $M_0$  and Killing fields  $\xi$  of  $g_-$  on  $M_-$ .
- All respective objects correspond to parallel fields on ambient space  $\tilde{M}$ .