

Elliptic solutions of the Yang-Baxter equation

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20 March 2017, Vienna

Based on work in collaboration with
S.É. Derkachov and V.P. Spiridonov
Comm. Math. Phys. (2016) 345

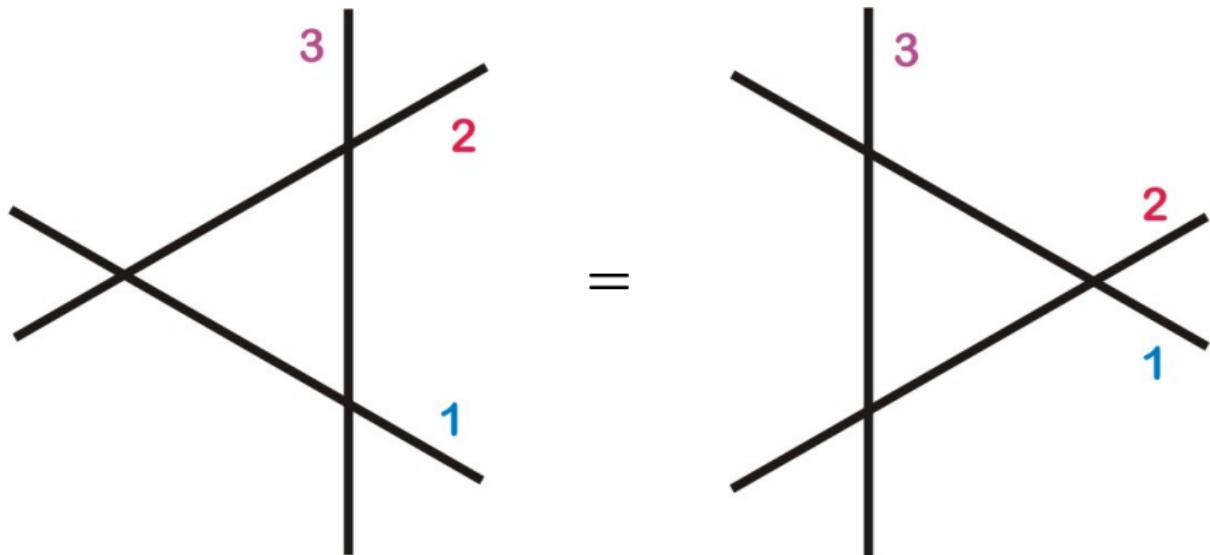
Overview

- The Yang-Baxter equation
- The master solution (integral operator with elliptic hypergeometric kernel)
- Continuous spin lattice models
- The star-triangle relation
- Factorization of the integral R-operator
- Finite-dimensional reductions

The Yang-Baxter equation

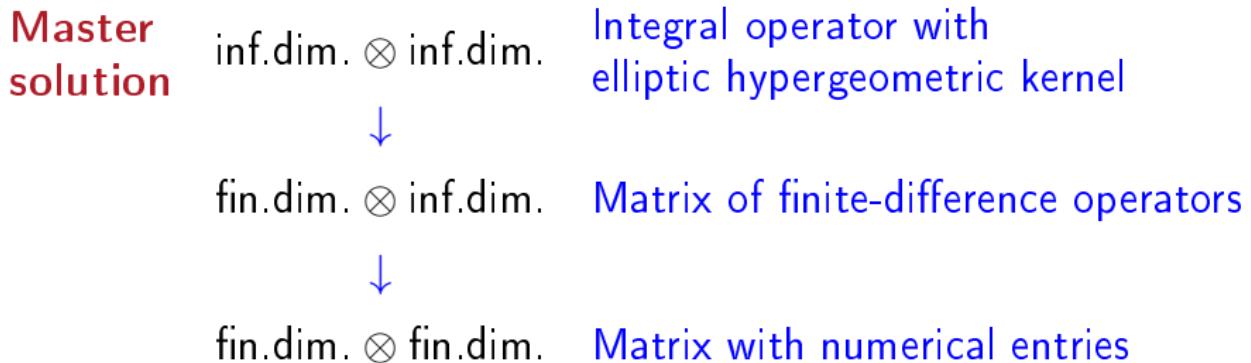
$$\mathbb{R}_{ik}(u) : V_i \otimes V_k \rightarrow V_i \otimes V_k$$

$$\mathbb{R}_{12}(u-v) \mathbb{R}_{13}(u) \mathbb{R}_{23}(v) = \mathbb{R}_{23}(v) \mathbb{R}_{13}(u) \mathbb{R}_{12}(u-v)$$



The Yang-Baxter equation

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- Fin. dim. irrep. – physical discrete spin systems on lattices (Ising model)
- Inf. dim. irrep. – noncompact spin chains, continuous spin systems

Basic definitions

The infinite q -product

$$(z; q)_\infty \equiv \prod_{k=0}^{\infty} (1 - zq^k) \quad , \quad |q| < 1$$

The shorten theta function

$$\theta(z; p) \equiv (z; p)_\infty (pz^{-1}; p)_\infty = \frac{1}{(p; p)_\infty} \sum_{k \in \mathbb{Z}} (-1)^k p^{\frac{k(k-1)}{2}} z^k$$

Quasi-periodicity

$$\theta(pz; p) = -z^{-1} \theta(z; p)$$

The elliptic gamma function

$$\Gamma(z; p, q) = \prod_{i,j=0}^{\infty} \frac{1 - z^{-1} p^{i+1} q^{j+1}}{1 - z p^i q^j} \quad , \quad |p| < 1, |q| < 1$$

[S.N.M. Ruijsenaars '97]

Two quasi periods

$$\Gamma(qz; p, q) = \theta(z; p) \Gamma(z; p, q) \quad , \quad \Gamma(pz; p, q) = \theta(z; q) \Gamma(z; p, q)$$

The elliptic beta integral

Shorthand notation $\Gamma(tz^{\pm 1}) \equiv \Gamma(tz)\Gamma(tz^{-1})$, $\Gamma(tx^{\pm 1}z^{\pm 1}) \equiv \Gamma(txz^{\pm 1})\Gamma(tx^{-1}z^{\pm 1})$

Exactly computable elliptic hypergeometric integral

$$\kappa \int_{\mathbb{T}} \frac{\prod_{j=1}^6 \Gamma(t_j z^{\pm 1}; p, q)}{\Gamma(z^{\pm 2}; p, q)} \frac{dz}{iz} = \prod_{1 \leq j < k \leq 6} \Gamma(t_j t_k; p, q)$$

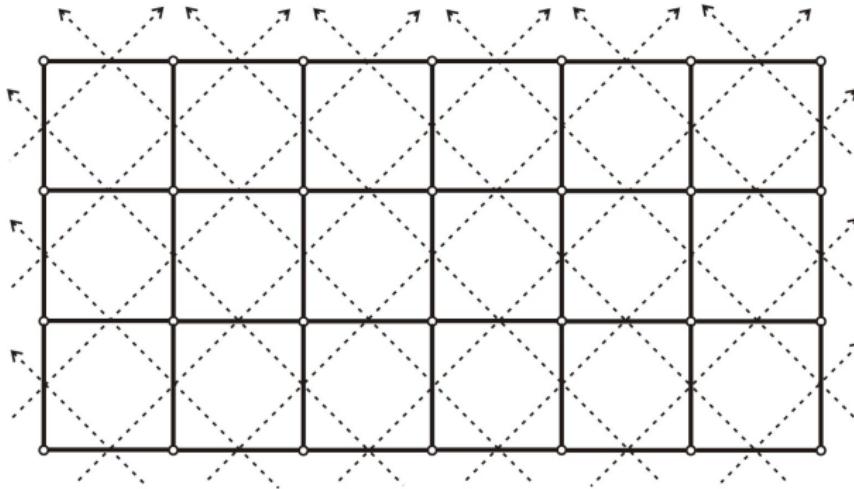
[V.Spiridonov '01]

balancing condition $\prod_{j=1}^6 t_j = pq$

$$\kappa \equiv \frac{(p; p)_{\infty} (q; q)_{\infty}}{4\pi} , \quad |t_j| < 1 , \quad j = 1, \dots, 6$$

The most general exact univariate integration formula; generalization of the Euler beta integral $\int_0^1 t^{\alpha-1} (1-t)^{\beta-1} dt = \frac{\Gamma(\alpha)\Gamma(\beta)}{\Gamma(\alpha+\beta)}$

Solvable 2d lattice models



Partition function

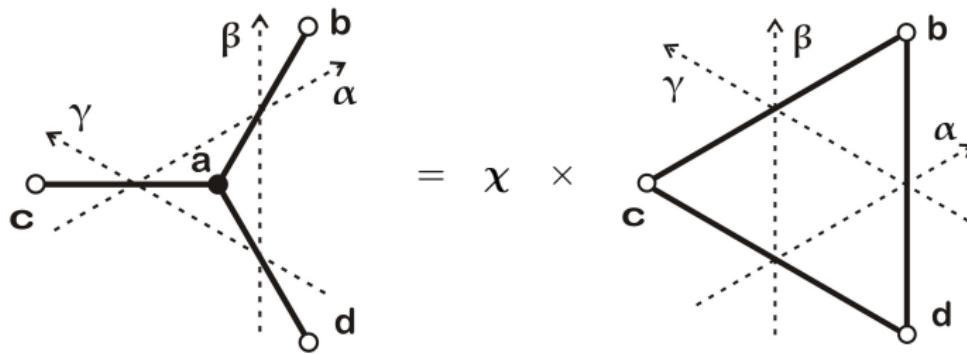
$$\mathcal{Z} = \int_{\substack{\mathbb{T} \times \dots \times \mathbb{T} \\ \# \text{ of vertex}}} \prod_{\substack{\text{horizontal} \\ \langle i,j \rangle}} W_{\alpha_i - \beta_j}(a_i, a_j) \prod_{\substack{\text{vertical} \\ \langle k,l \rangle}} \overline{W}_{\alpha_k - \beta_l}(a_k, a_l) \prod_{\substack{\text{vertices} \\ m}} \rho(a_m) da_m$$

Continuous spin variables $a_m \in \mathbb{T}$. Rapidities α_i, β_j

Star-triangle relation and Z-invariance

$$\int_0^1 \rho(a) W_{\xi-(\alpha-\beta)}(a, b) W_{\alpha-\gamma}(c, a) W_{\xi-(\beta-\gamma)}(d, a) da \\ = \chi \cdot W_{\alpha-\beta}(c, d) W_{\xi-(\alpha-\gamma)}(d, b) W_{\beta-\gamma}(c, b)$$

[V.Bazhanov, S.Sergeev '10]



Spin variables a, b, c, d, \dots Rapidity $\alpha, \beta, \gamma, \dots$

Edge Boltzmann weights $W_\alpha(a, b) = W_\alpha(b, a)$, $\overline{W}_\alpha(a, b) = W_{\xi-\alpha}(a, b)$

$$W_\alpha(a, b) = \Gamma(\sqrt{pq} e^{2\pi i(\alpha \pm a \pm b)})$$

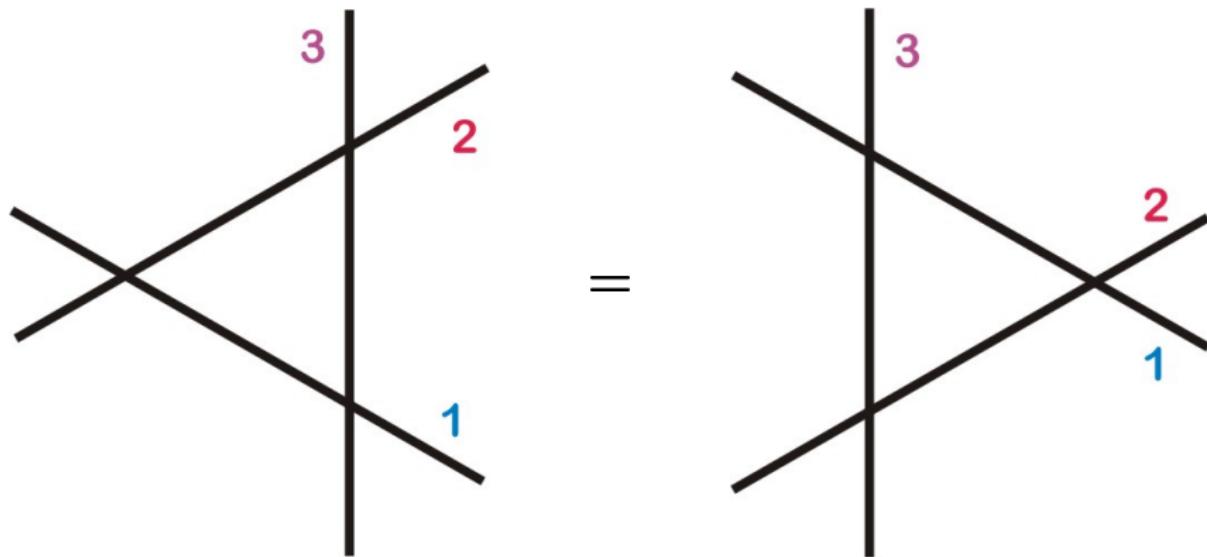
Self-interaction $\rho(a) = \kappa / \Gamma(e^{\pm 4\pi i a})$

Normalization $\chi = \chi(\alpha, \beta, \gamma)$

Crossing-parameter ξ , $e^{-2\xi} \equiv pq$

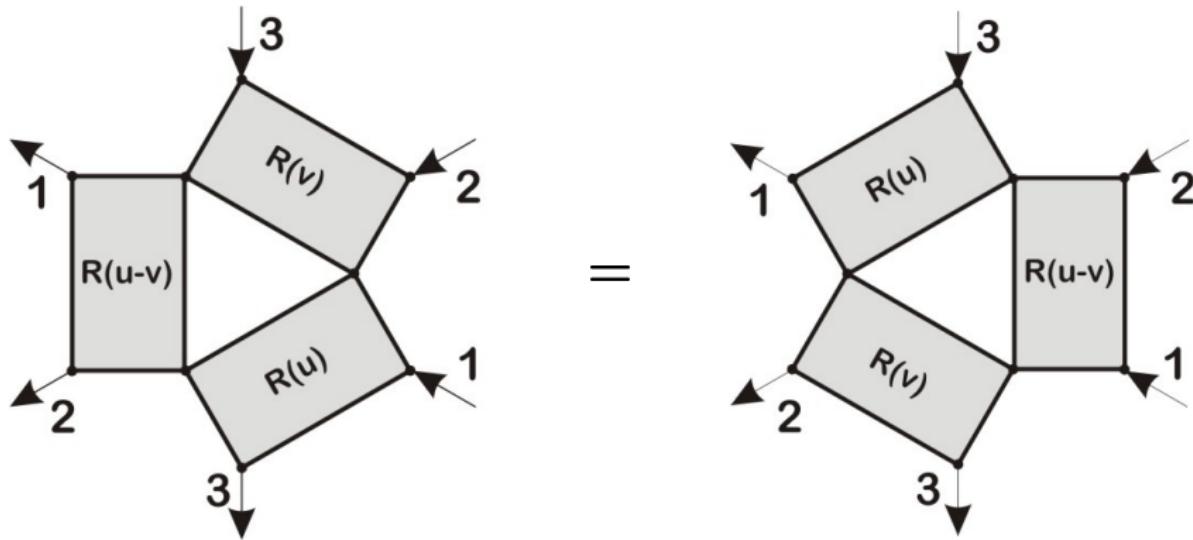
The Yang-Baxter equation

$$\mathbb{R}_{12}(u - v) \mathbb{R}_{13}(u) \mathbb{R}_{23}(v) = \mathbb{R}_{23}(v) \mathbb{R}_{13}(u) \mathbb{R}_{12}(u - v)$$



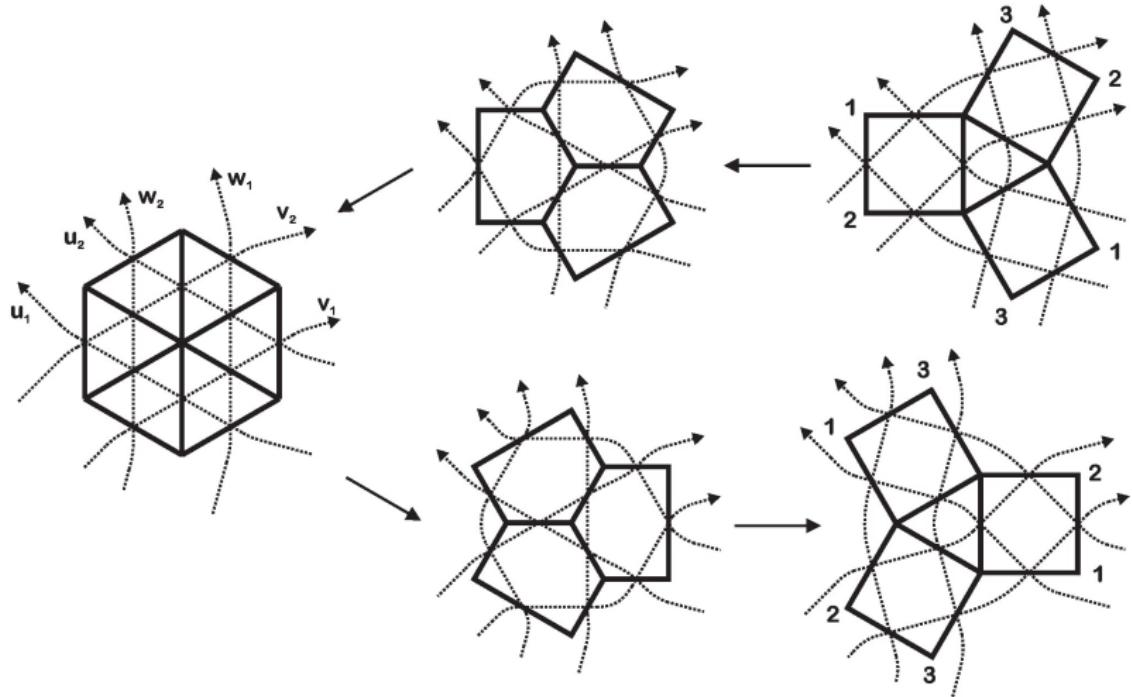
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The Yang-Baxter equation

$$\mathbb{R}_{12}(u_1, u_2 | v_1, v_2) \mathbb{R}_{13}(u_1, u_2 | w_1, w_2) \mathbb{R}_{23}(v_1, v_2 | w_1, w_2)$$
$$= \mathbb{R}_{23}(v_1, v_2 | w_1, w_2) \mathbb{R}_{13}(u_1, u_2 | w_1, w_2) \mathbb{R}_{12}(u_1, u_2 | v_1, v_2)$$



From the lattice model to the master R-operator

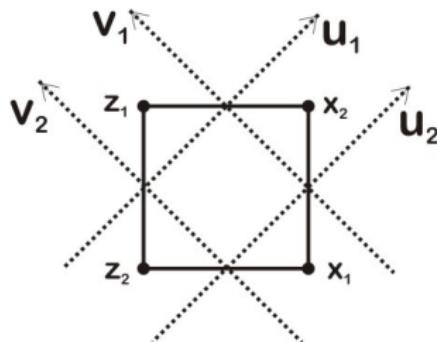
Boltzmann weights

$$W_u(x, z) = \Gamma(\sqrt{pq} e^{2\pi i(u \pm x \pm z)})$$

$$\overline{W}_u(x, z) = W_{\xi - u}(x, z)$$

Self-interaction

$$\rho(x) = \kappa / \Gamma(e^{\pm 4\pi i x})$$



$$\begin{aligned} & \left[\mathbb{R}_{12}(u - v | g_1, g_2) \Phi \right] (z_1, z_2) = W_{u_1 - v_1}(z_1, z_2) \times \\ & \times \int \overline{W}_{u_1 - v_1}(z_1, x_2) \overline{W}_{u_2 - v_2}(z_2, x_1) W_{u_2 - v_1}(x_1, x_2) \rho(x_1) \rho(x_2) \Phi(x_1, x_2) dx_1 dx_2 \end{aligned}$$

[S.Derkachov, V.Spiridonov '12]

The R-operator $\in \text{End}(V_{g_1} \otimes V_{g_2})$

Partition function
of an elementary cell

Spectral parameters u, v , g-spins g_1, g_2

Rapidities u_1, u_2, v_1, v_2

$$u_1 = \frac{u + g_1}{2}, \quad u_2 = \frac{u - g_1}{2}, \quad v_1 = \frac{v + g_2}{2}, \quad v_2 = \frac{v - g_2}{2}$$

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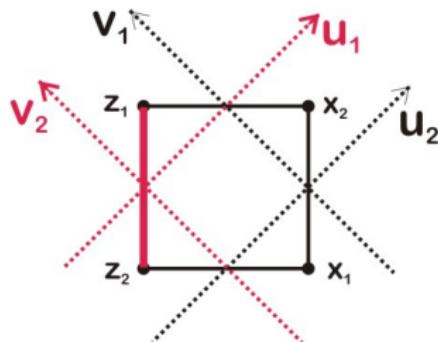
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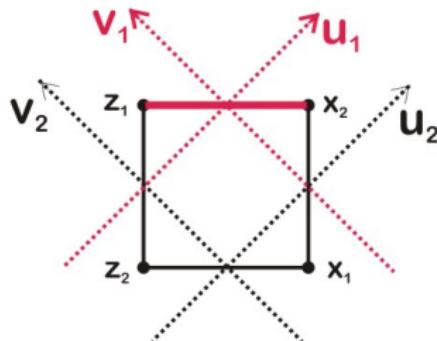
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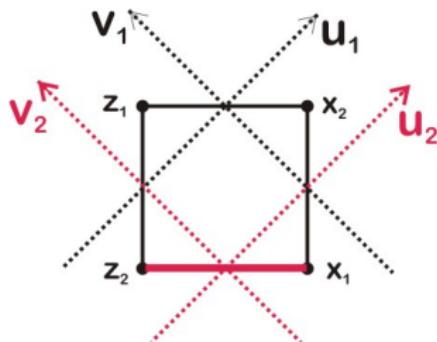
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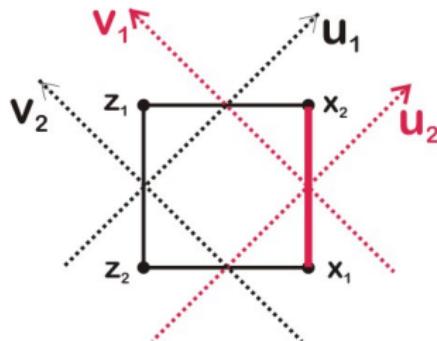
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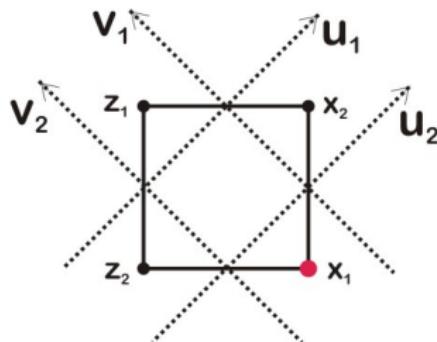
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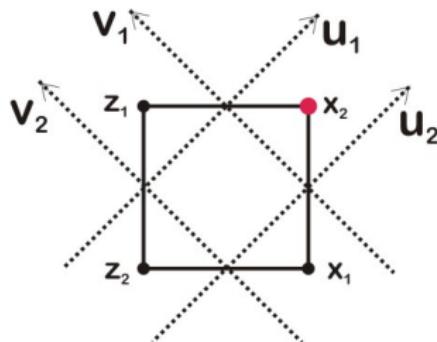
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Baxter's R-matrix and 8-vertex model

Matrix 4×4 solution of the YBE

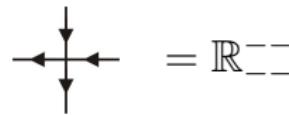
[R. Baxter '72]

$$\mathbb{R}_{12}(u) = \sum_{\alpha=0}^3 \frac{\theta_{\alpha+1}(u + \eta|\tau)}{\theta_{\alpha+1}(\eta|\tau)} \sigma_\alpha \otimes \sigma_\alpha$$

θ_α – Jacobi theta functions, σ_α – Pauli matrices, $p \equiv e^{2\pi i \tau}$, $q \equiv e^{4\pi i \eta}$



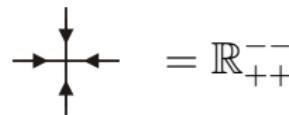
$$= \mathbb{R}_{++}^{++}$$



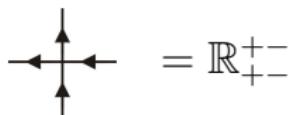
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Sklyanin algebra and XYZ spin chain

Lax operator – Matrix 2×2 solution of the YBE with operator entries

$$L(u) = \sum_{\alpha=0}^3 \frac{\theta_{\alpha+1}(u + \eta|\tau)}{\theta_{\alpha+1}(\eta|\tau)} \sigma_\alpha \otimes S^\alpha$$

Four generators S^0, S^1, S^2, S^3 respect commutation relations

$$S^\alpha S^\beta - S^\beta S^\alpha = i \cdot (S^0 S^\gamma + S^\gamma S^0)$$

$$S^0 S^\alpha - S^\alpha S^0 = i J_{\beta\gamma} \cdot (S^\beta S^\gamma + S^\gamma S^\beta)$$

[E.Sklyanin '82]

where (α, β, γ) – arbitrary cyclic permutation of $(1, 2, 3)$

$J_{\alpha\beta} = J_{\alpha\beta}(2\eta, \tau)$ – structure constants

Sklyanin algebra and XYZ spin chain

Realization of the algebra by finite-difference operators

$$\mathbf{S}^\alpha(g) = e^{\pi iz^2/\eta} \frac{i^{\delta_{\alpha,2}} \theta_{\alpha+1}(\eta|\tau)}{\theta_1(2z|\tau)} \left[\theta_{\alpha+1}(2z-g+\eta|\tau) e^{\eta\partial_z} - \theta_{\alpha+1}(-2z-g+\eta|\tau) e^{-\eta\partial_z} \right] e^{-\pi iz^2/\eta}$$

[E.Sklyanin '83]

Label g of the representation

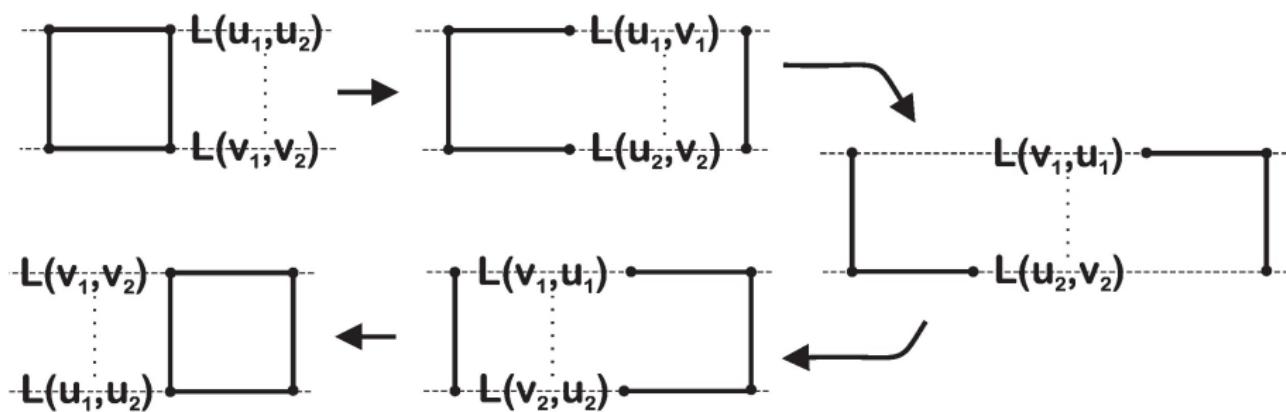
$$\text{Elliptic Modular Double} = \begin{matrix} \mathbf{S}^0, \mathbf{S}^1, \mathbf{S}^2, \mathbf{S}^3 & \tilde{\mathbf{S}}^0, \tilde{\mathbf{S}}^1, \tilde{\mathbf{S}}^2, \tilde{\mathbf{S}}^3 \\ [L.Faddeev '99; V.Spiridonov '09] & 2\eta, \tau \end{matrix} \otimes \begin{matrix} \text{Sklyanin algebra} & \text{Sklyanin algebra} \\ \tau, 2\eta \end{matrix}$$

Factorization of the master R-operator

Lax operator $L(u|g) : \mathbb{C}^2 \otimes V_g \rightarrow \mathbb{C}^2 \otimes V_g$

$$u_1 \leftrightarrow u_2 \Leftrightarrow g \leftrightarrow -g$$

$$\mathbb{R}_{12}(u-v|g_1, g_2) L_1(u|g_1) L_2(v|g_2) = L_2(v|g_2) L_1(u|g_1) \mathbb{R}_{12}(u-v|g_1, g_2)$$



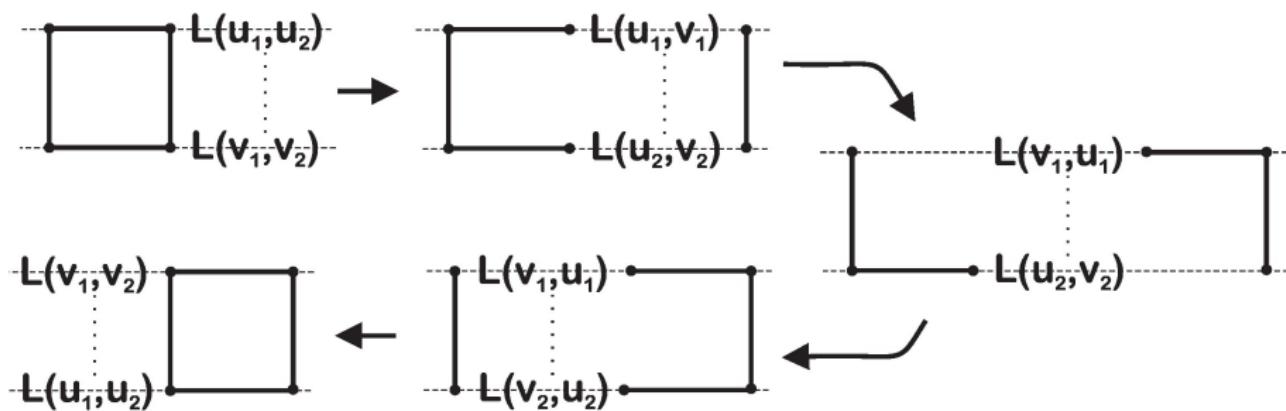
The constituents of the master R-operator form the operator representation of S_4

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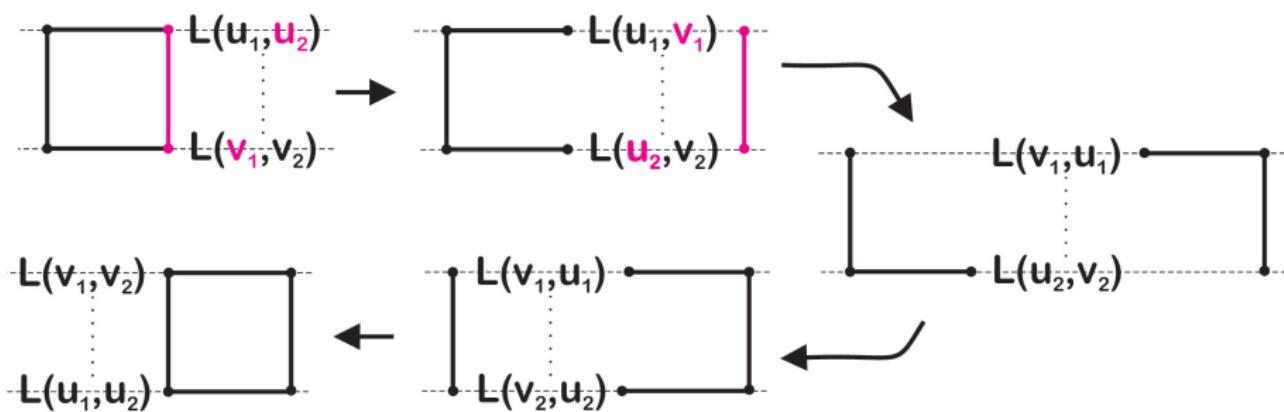
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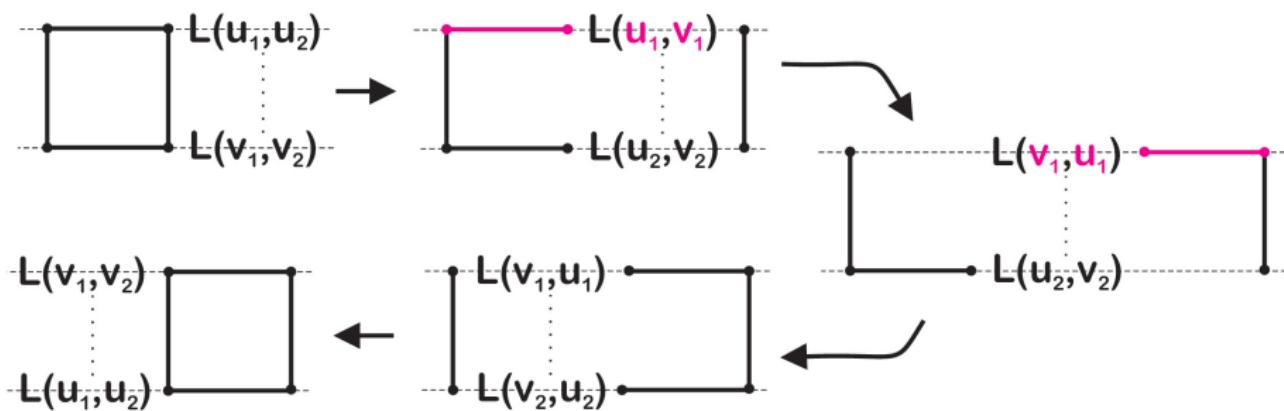
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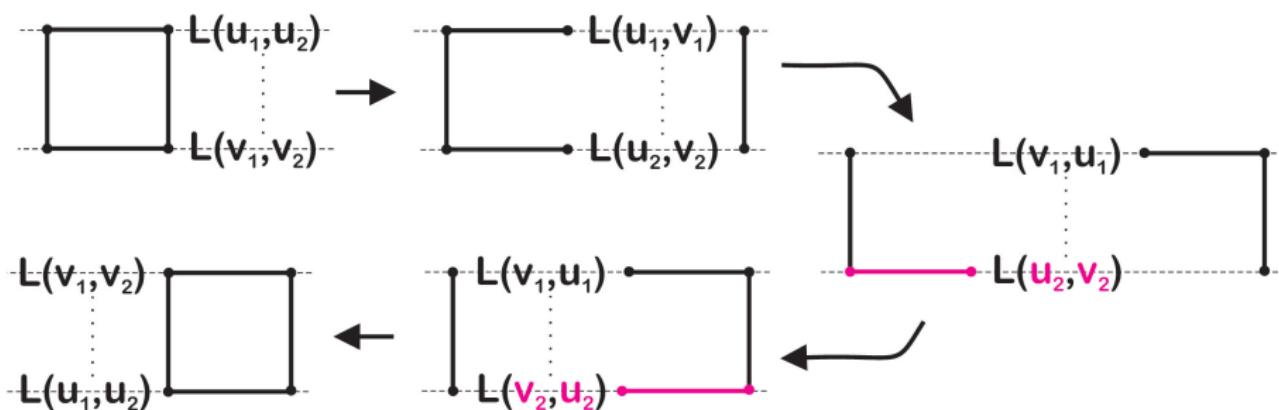
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Lax operator $L(u|g) : \mathbb{C}^2 \otimes V_g \rightarrow \mathbb{C}^2 \otimes V_g$

$$u_1 \leftrightarrows u_2 \Leftrightarrow g \leftrightarrows -g$$

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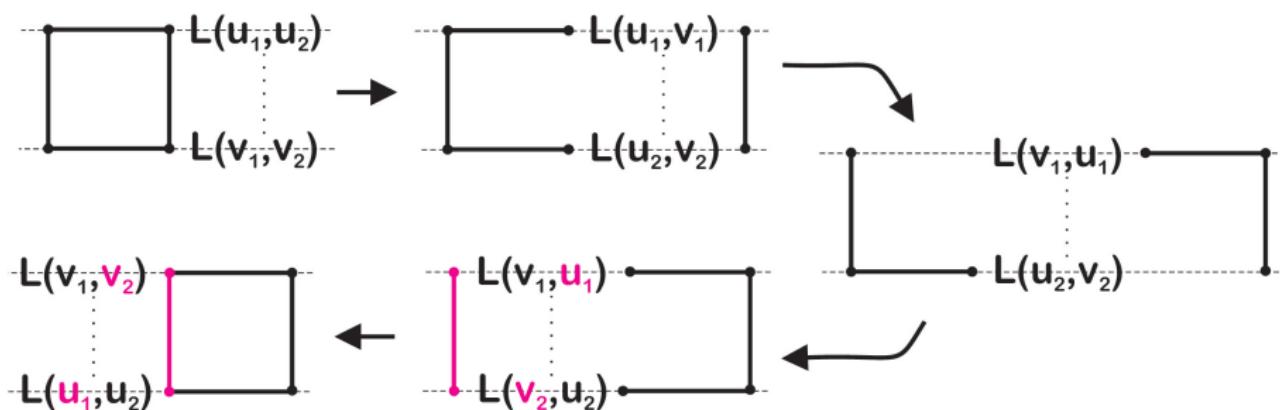
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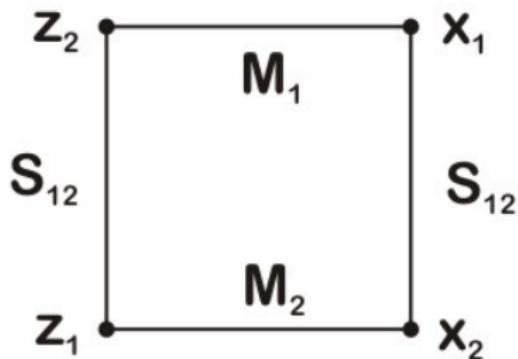
Additive notations $\Gamma(e^{2\pi i z}) \rightarrow \Gamma(z)$

- Permutation P_{12}

- Multiplication by

$$S_{12}(u) \equiv \Gamma(\pm z_1 \pm z_2 + u + \eta + \frac{\tau}{2})$$

- Elliptic Fourier transform



$$[M(g)\Phi](z) \equiv \kappa \int_0^1 \frac{\Gamma(\pm z \pm x - g)}{\Gamma(-2g)\Gamma(\pm 2x)} \Phi(x) dx$$

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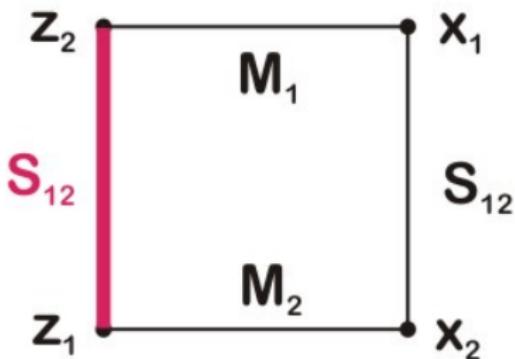
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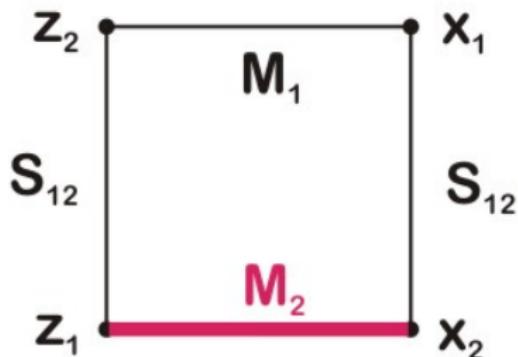
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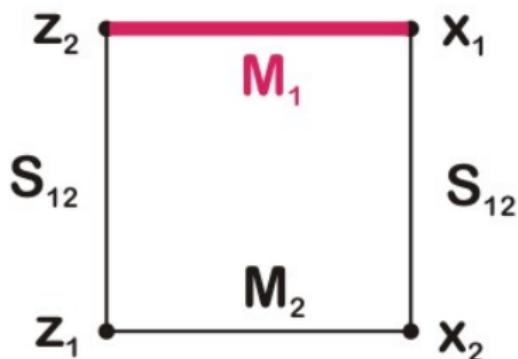
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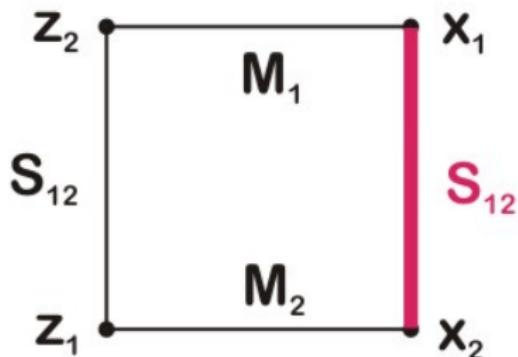
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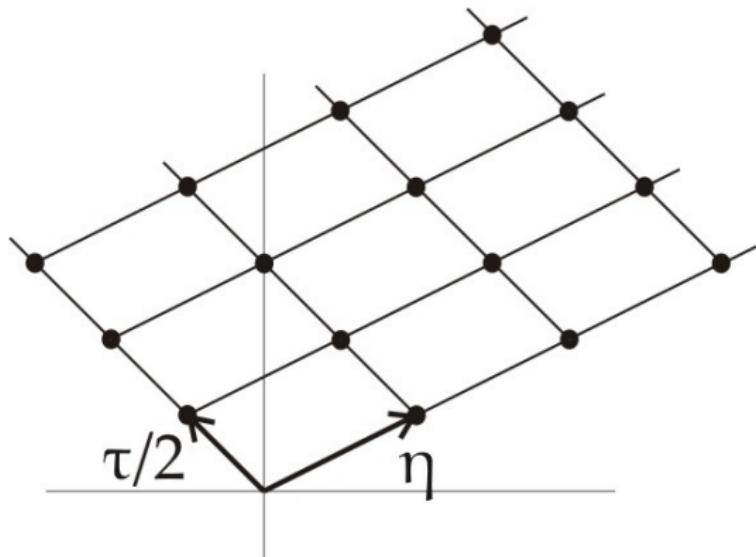
The intertwining operator for Sklyanin algebra

Discrete lattice
of g -spin

$$g = n\eta + m\frac{\tau}{2}$$

$$n, m \in \mathbb{Z}_{\geq 0}$$

Initial condition:
 $M(0) = 1$



Contiguous relations & Factorization of the intertwining operator

$$A(g) M(g) = M(g+\eta) \theta_3(z|\frac{\tau}{2}) \quad B(g) M(g) = M(g+\frac{\tau}{2}) \theta_3(z|\eta)$$



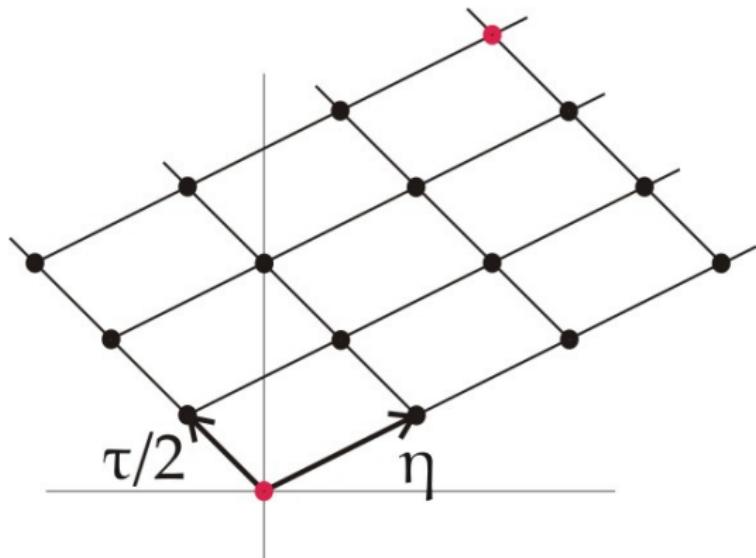
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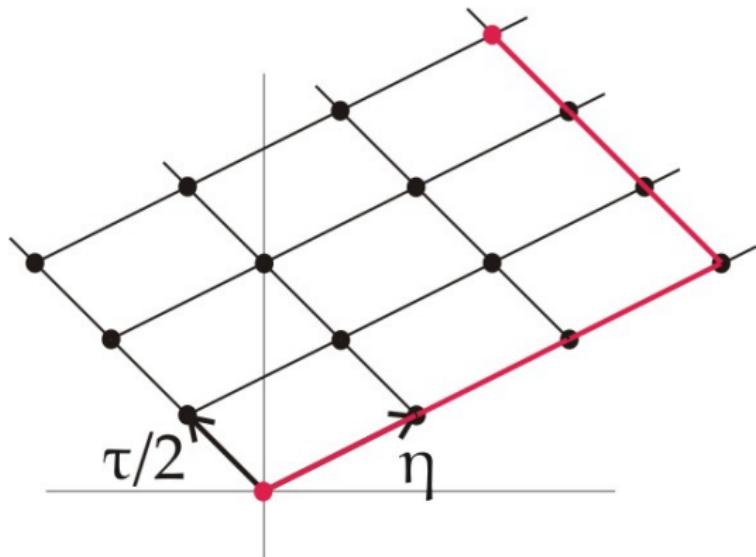
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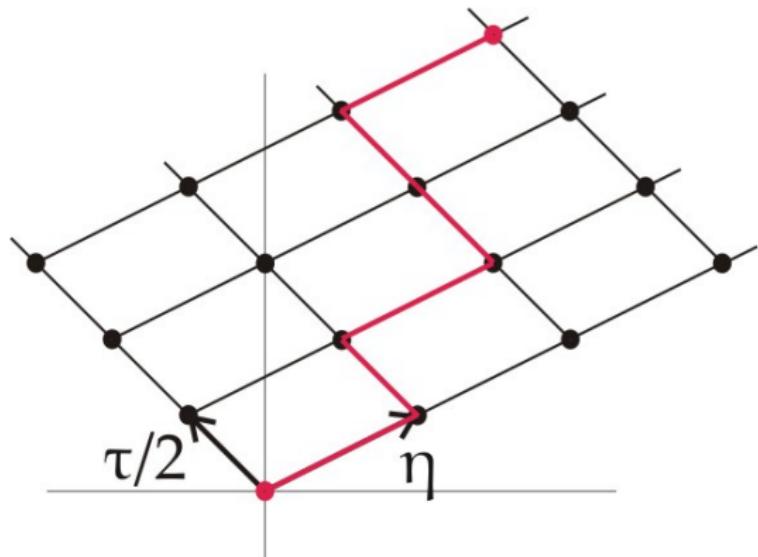
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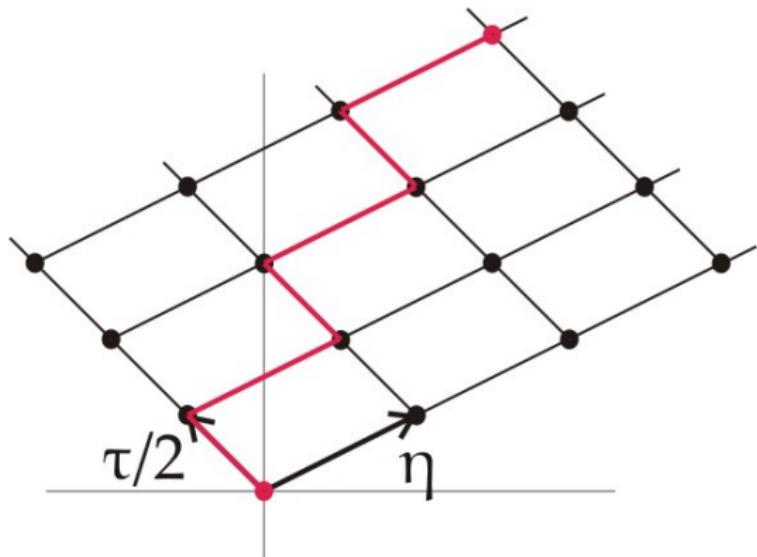
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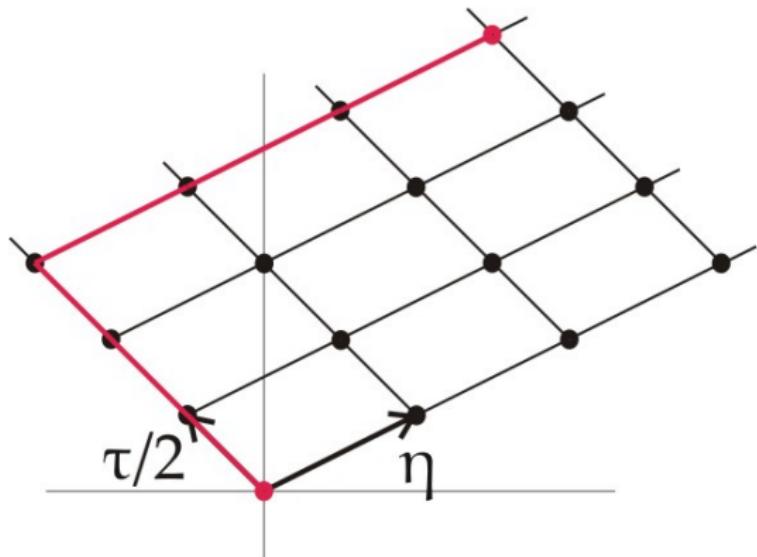
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The intertwining operator and finite-dimensional representations

Zero modes of $M(g_{n,m})$ at $g_{n,m} = (n+1)\eta + (m+1)\frac{\tau}{2}$ $n, m \in \mathbb{Z}_{\geq 0}$

Due to intertwining relations

$$M(g) S^\alpha(g) = S^\alpha(-g) M(g)$$

$$M_1(g_1) \mathbb{R}_{12}(u | g_1, g_2) = \mathbb{R}_{12}(u | -g_1, g_2) M_1(g_1)$$

the space $\text{Ker } M(g_{n,m}) \cap \text{Im } M(-g_{n,m})$ is invariant

Generating function of the finite-dimensional representation,

$$\Gamma(\pm z \pm x + \mathbf{g}_{n,m}) = \sum_{\substack{\leftarrow \text{auxiliary parameter}}} \varphi_{j,l}^{(n,m)}(z) \psi_{n-j,m-l}^{(n,m)}(x) \sim \Theta_{2n}^+(\tau) \otimes \Theta_{2m}^+(2\eta)$$

produces a pair of natural bases, $\dim = (n+1)(m+1)$,

$$\varphi_{j,l}^{(n,m)}(z) = [\theta_3(z|\frac{\tau}{2})]^j [\theta_4(z|\frac{\tau}{2})]^{n-j} [\theta_3(z|\eta)]^l [\theta_4(z|\eta)]^{m-l}$$

$$\psi_{j,l}^{(n,m)}(z) = \text{Sym} \prod_{r=0}^{n-1} \theta_{3,4} \left(z + (n-1-2r)\eta \mid \frac{\tau}{2} \right) \text{Sym} \prod_{s=0}^{m-1} \theta_{3,4} \left(z + (m-1-2s)\frac{\tau}{2} \mid \eta \right)$$

$$\quad \quad \quad j=0,1,2,\dots,n, \quad l=0,1,2,\dots,m$$



Reduction of the general R-operator

Reduce the master R-operator to finite-dimensional rep. in the first space

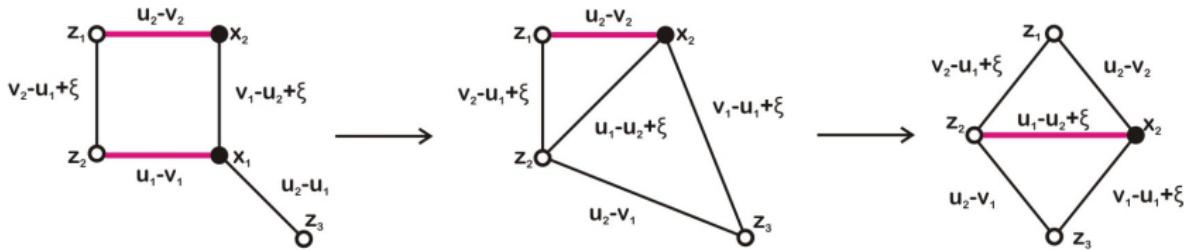
$$\mathbb{R}_{12}(u|g_{n,m}, g) : \mathbb{C}^{(n+1)(m+1)} \otimes V_g \rightarrow \mathbb{C}^{(n+1)(m+1)} \otimes V_g$$

Matrix with operator entries which are finite-difference operators of a high order

↓ auxiliary parameter

$$\mathbb{R}_{12}(u|g_{n,m}, g) \Gamma(\mp z_1 \mp z_3 + g_{n,m}) \Phi(z_2)$$

$$= \frac{\Gamma(\mp z_2 \mp z_3 + \frac{g_{n,m}+g-u}{2})}{\Gamma(\mp z_1 \mp z_2 - \frac{g_{n,m}+g+u}{2} + \eta + \frac{\tau}{2})} M_2(n\eta + m \frac{\tau}{2}) \frac{\Gamma(\mp z_1 \mp z_2 + \frac{g_{n,m}-g-u}{2}) \Phi(z_2)}{\Gamma(\mp z_2 \mp z_3 + \frac{g-g_{n,m}-u}{2} + \eta + \frac{\tau}{2})}$$



Reduced R-operator is written in the pair of bases

$$\mathbb{R}_{12}(u|g_{n,m}, g) \psi_{j,l}^{(n,m)}(\textcolor{teal}{z}_1) = \varphi_{r,s}^{(n,m)}(\textcolor{teal}{z}_1) [\mathbb{R}]_{j,l}^{r,s}$$

All fin.dim. \otimes inf.dim. solutions **factorize**

$$\mathbb{R}(u|g_{n,0}, g) = V\left(\frac{u+g}{2}, z\right) D(z, \partial) \mathbf{C} V^T\left(\frac{u-g}{2}, z\right) \mathbf{C}$$

$D(z, \partial)$ – diagonal matrix, \mathbf{C} – constant numerical matrix

Example: the Lax operator, we recover the Sklyanin algebra

$$\begin{aligned} \mathbb{R}(u|g_{1,0}, g) &= e^{\pi i z^2 / \eta} \frac{1}{\theta_1(2z|\tau)} \begin{pmatrix} \theta_3\left(z - \frac{u+g}{2} | \frac{\tau}{2}\right) & -\theta_3\left(z + \frac{u+g}{2} | \frac{\tau}{2}\right) \\ -\theta_4\left(z - \frac{u+g}{2} | \frac{\tau}{2}\right) & \theta_4\left(z + \frac{u+g}{2} | \frac{\tau}{2}\right) \end{pmatrix} \cdot \\ &\cdot \begin{pmatrix} e^{\eta \partial} & 0 \\ 0 & e^{-\eta \partial} \end{pmatrix} \begin{pmatrix} \theta_4\left(z + \frac{u-g}{2} | \frac{\tau}{2}\right) & \theta_3\left(z + \frac{u-g}{2} | \frac{\tau}{2}\right) \\ \theta_4\left(z - \frac{u-g}{2} | \frac{\tau}{2}\right) & \theta_3\left(z - \frac{u-g}{2} | \frac{\tau}{2}\right) \end{pmatrix} e^{-\pi i z^2 / \eta} \end{aligned}$$

Summary and Conclusions

- The master solution (principal series rep.) of the YBE
- Continuous spin lattice models and the star-triangle
- Finite-dimensional reductions of the master solution
- Intertwining operators
- The factorized form of the integral R-operator and its finite-dimensional reductions (Explicit concise formulae)

Rational \leftarrow Trigonometric \leftarrow Elliptic

$SL(2, \mathbb{C}) \leftarrow$ Modular double
of $U_q(sl_2)$ \leftarrow Hyperbolic
modular double \leftarrow Elliptic modular double
(Sklyanin algebra)