## Bailey pairs and strange identities

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## Recap

q-hypergeometric series are everywhere!

q-series identities are fundamental and beautiful.

The structure of Bailey pairs lies behind many important identities.

Once the iterative framework is set up, results like the Andrews-Gordon identities come easily.

One identity gives infinitely many identities.

Bailey pairs can be used to express Ramanujan's mock theta functions in terms of Appell-type series or indefinite theta functions, whose modular transformation properties were determined by Zwegers.



Ramanujan's mock theta functions can then be completed to non-holomorphic modular forms via an Eichler integral.

Standard applications of the Bailey chain do not produce families of q-hypergeometric mock theta functions, but a more nuanced approach does.

And the corresponding Bailey pairs have nice applications, including to colored Jones polynomials of certain torus knots and WRT invariants of 3-manifolds.

In fact, the colored Jones polynomials and their cyclotomic coefficients essentially form a Bailey pair!

### Remark

If you thought that the two "mini-courses" at SLC 87 were completely unrelated, see:

S.O. Warnaar, The Bailey lemma and Kostka polynomials, *J. Algebraic Combin.* **20** (2004), no. 2, 131–171.

(This paper concers a "higher level" Bailey lemma...)

## Outline

Zagier's strange identity and applications

② General strange identities

O New families of strange identities

Proofs using Bailey pairs

# Zagier's "identity"

Zagier (2001) recorded the "strange identity"

$$\sum_{n\geq 0} (q)_n " = " - \frac{1}{2} \sum_{n\geq 1} n\left(\frac{12}{n}\right) q^{(n^2-1)/24}.$$

Here

$$\left(\frac{12}{n}\right) = \begin{cases} 1, & \text{if } n \equiv \pm 1 \pmod{12}, \\ -1, & \text{if } n \equiv \pm 5 \pmod{12}, \\ 0, & \text{otherwise.} \end{cases}$$

Note that

$$(q)_{\infty} = \sum_{n\geq 0} \left(\frac{12}{n}\right) q^{(n^2-1)/24}$$

The right-hand side of Zagier's "identity" is

$$-rac{1}{2}\left(1-5q-7q^2+11q^5+13q^7-17q^{12}-\cdots
ight).$$

This is an analytic function for |q| < 1.

The left-hand side is not an analytic function on any open subset of  $\mathbb{C}.$ 

So, what does Zagier mean by " = "?

# Zagier's "identity"

Note that the left-hand side is a finite sum (and hence well-defined) when q is a root of unity.

In particular, if  $q^N = 1$  then

$$\sum_{n=0}^{\infty} (q)_n = \sum_{n=0}^{N-1} (q)_n.$$

Similarly, for any root of unity  $\zeta$ , the left-hand side is a well-defined power series in t when

$$q = \zeta e^{-t} = \zeta \left( \sum_{k \ge 0} \frac{(-1)^k t^k}{k!} \right)$$

# Zagier's "identity"

If we replace q by  $\zeta e^{-t}$  on the right-hand side, classical results in asymptotic analysis give that as  $t \to 0^+$ ,

$$-\frac{1}{2}\sum_{n\geq 1} n\left(\frac{12}{n}\right) \zeta^{(n^2-1)/24} e^{-t(n^2-1)/24} \sim \text{ a power series in } t.$$

The "=" in Zagier's "identity" means that at any root of unity these two power series are equal.

## Zagier's "identity"

For example, at q = 1 we have

$$\sum_{n\geq 0} (e^{-t}; e^{-t})_n = 3 + 11t + \frac{133}{2}t^2 + \frac{3389}{6}t^3 + \frac{148177}{24}t^4 + \cdots$$
$$\sim -\frac{1}{2}\sum_{n\geq 1} n\left(\frac{12}{n}\right)e^{-t(n^2-1)/24}.$$

At q = -1 we have

$$\sum_{n\geq 0} (-e^{-t}; -e^{-t})_n = 1 + t + \frac{3}{2}t^2 + \frac{19}{6}t^3 + \frac{207}{24}t^4 + \cdots$$
$$\sim -\frac{1}{2}\sum_{n\geq 1} n\left(\frac{12}{n}\right)(-1)^{(n^2-1)/24}e^{-t(n^2-1)/24}.$$



It turns out that Zagier's strange identity has many interesting applications.

Let

$$F(q)=\sum_{n\geq 0}(q)_n.$$

This is called the "Kontsevich-Zagier function".

# Applications - values of F(q)

Habiro (2000) and Lê (2003) showed that the colored Jones polynomial of the trefoil knot can be written

$$J_N(q) = q^{1-N} \sum_{n \ge 0} q^{-nN} (q^{1-N})_n.$$

Therefore

$$F(e^{2\pi i/N}) = J_N(e^{2\pi i/N})$$
 (The Kashaev invariant).

Using the strange identity, one can show that if L is such that  $\zeta^{\frac{L}{12}}=1,$  then

$$F(\zeta) = \frac{1}{4L} \sum_{m=1}^{L} m^2 \left(\frac{12}{m}\right) \zeta^{(m^2-1)/24}.$$

# Applications - quantum modularity

A quantum modular form of weight k is a function  $f : \mathbb{Q} \to \mathbb{C}$  such that for all  $\gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$  in (a subgroup of)  $SL_2(\mathbb{Z})$ , one has that

$$g_{\gamma}(x) = f\left(rac{ax+b}{cx+d}
ight) - ullet(cx+d)^k f(x)$$

is a a continuous (or smooth) function on  $\mathbb{R}\setminus\{\frac{-d}{c}\}$ .

Using the strange identity one can show that the function  $\phi:\mathbb{Q}\to\mathbb{C}$  defined by

$$\phi(x) = e^{\pi i x/12} F(e^{2\pi i x})$$

is a weight 3/2 quantum modular form.

# Applications - congruences for Fishburn numbers

The Kontsevich-Zagier series at 1 - q,

$$F(1-q) = 1 + q + 2q^2 + 5q^3 + 15q^4 + 53q^5 + \cdots$$

is a well-known combinatorial generating function.

It counts Fishburn matrices, (2 + 2)-free posets, ascent sequences, linearized chord diagrams,... (Bousquet-Mélou, Claesson, Dukes, Kitaev, 2010)

For example, a Fishburn matrix is an upper-triangular matrix with non-negative integer entries such that no row or column consists only of zeros.

## Applications - congruences for Fishburn numbers

The five Fishburn matrices for n = 3 are

$$(3), \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}, \begin{pmatrix} 2 & 0 \\ 0 & 1 \end{pmatrix}, \begin{pmatrix} 1 & 0 \\ 0 & 2 \end{pmatrix}, \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Using the strange identity, one can show that the Fishburn numbers satisfy many Ramanujan-type congruences, like

$$a(5n+r) \equiv 0 \pmod{5}, r \in \{3,4\}$$
  
 $a(7n+6) \equiv 0 \pmod{7},$   
 $a(11n+s) \equiv 0 \pmod{11}, s \in \{8,9,10\},$ 

# Strange identity - definition

Generalizing Zagier's example, a strange identity

$$\sum_{n\geq 0} (q)_n f_n(q) = \sum_{n\geq 0} n^{\nu} \chi(n) q^{(n^2-a)/b}$$

means that for  $q = \zeta e^{-t}$ , the right-hand side has an asymptotic expansion as a power series as  $t \to 0^+$  and this power series is given by the left-hand side at  $\zeta e^{-t}$ .

Here the  $f_n$  are polynomials,  $v \in \{0, 1\}$ ,  $a \ge 0$  and b > 0 are integers, and  $\chi : \mathbb{Z} \to \mathbb{C}$  is a periodic function such that:

(i) 
$$\chi(n) \neq 0$$
 if and only if  $\frac{n^2-a}{b} \in \mathbb{Z}$ ,

(ii) The function  $n \to \zeta^{(n^2-a)/b}\chi(n)$  is a periodic function with mean value zero for any root of unity  $\zeta$ .

# Strange identity - definition

This definition can be relaxed in several ways, notably by replacing  $(q)_n$  on the left-hand side by other *q*-factorials, such as  $(q; q^2)_n$  or  $\frac{(q)_n}{(-q)_n}$ , and appropriately restricting to a subset of the roots of unity.

In these two cases, for example, we only consider the odd roots of unity.

One can also let v be any positive integer.

# General results

Following Zagier's work, many applications of Zagier's strange identity were extended to general strange identities:

1) Bringmann-Rolen and Goswami-Osburn showed that *q*-hypergeometric series satisfying a strange identity are quantum modular forms.

2) Guerzhoy-Kent-Rolen and Ahlgren-Kim-L. showed that if G(q) satisfies a strange identity, then the coefficients of G(1-q) will have many Ramanujan-type congruences.

3) Goswami-Jha-Kim-Osburn found an asymptotic formula for the coefficients of G(1-q) when G(q) satisfies a strange identity.

# Hikami's strange identity

Many nice results and very few examples.

Aside from a few isolated examples, until recently the only other strange identities are due to Hikami (2006).

We need the q-binomial coefficient (or "Gaussian polynomial") defined by

$$\begin{bmatrix} n \\ k \end{bmatrix} = \begin{bmatrix} n \\ k \end{bmatrix}_q = \begin{cases} \frac{(q)_n}{(q)_{n-k}(q)_k}, & \text{if } 0 \le k \le n, \\ 0, & \text{otherwise.} \end{cases}$$

## Hikami's strange identity

For  $0 \le a \le k-1$  we have

$$\sum_{n_1,\dots,n_k\geq 0} (q)_{n_k} q^{n_1^2+\dots+n_{k-1}^2+n_{a+1}+\dots+n_{k-1}} \prod_{i=1}^{k-1} \begin{bmatrix} n_{i+1}+\delta_{i,a} \\ n_i \end{bmatrix}$$
  
$$= -\frac{1}{2} \sum_{n\geq 0} n\chi_{8k+4}^{(a)}(n) q^{\frac{n^2-(2k-2a-1)^2}{8(2k+1)}},$$

where  $\chi_{8k+4}^{(a)}(n)$  is the even periodic function modulo 8k + 4 defined by

$$\chi_{8k+4}^{(a)}(n) = \begin{cases} 1, & \text{if } n \equiv 2k - 2a - 1 \text{ or } 6k + 2a + 5 \pmod{8k+4}, \\ -1, & \text{if } n \equiv 2k + 2a + 3 \text{ or } 6k - 2a + 1 \pmod{8k+4}, \\ 0, & \text{otherwise.} \end{cases}$$

## Hikami's strange identity

The case k = 1 is Zagier's strange identity.

For k = 2 we have

$$\sum_{n_1,n_2\geq 0} (q)_{n_2} q^{n_1^2+n_1} \begin{bmatrix} n_2\\ n_1 \end{bmatrix}$$
  
$$" = " - \frac{1}{2} \sum_{n\geq 0} n\chi_{20}^{(0)}(n) q^{\frac{n^2-9}{40}}$$

and

$$\sum_{n_1,n_2\geq 0} (q)_{n_2} q^{n_1^2} \begin{bmatrix} n_2+1\\ n_1 \end{bmatrix}$$
  
"="  $-\frac{1}{2} \sum_{n\geq 0} n\chi_{20}^{(1)}(n) q^{\frac{n^2-1}{40}}.$ 

# Hikami's strange identity - Remarks

One can deduce formulas, quantum modularity, asymptotics, congruences, etc.

The case (k, 0) corresponds to the Kashaev invariant for the torus knot (2, 2k + 1).

Very recently another family of strange identities was found, corresponding to the torus knots  $(3, 2^t)$ .

Can any of this be understood in the context of Bailey pairs?

Yes!

# Results (L., 2022)

Here is a sample of what one can prove using the Bailey machinery.

Let  $\chi_{4k}(n)$  be the even periodic function modulo 4k defined by

$$\chi_{4k}(n) = \begin{cases} 1, & \text{if } n \equiv k - 1 \text{ or } 3k + 1 \pmod{4k}, \\ -1, & \text{if } n \equiv k + 1 \text{ or } 3k - 1 \pmod{4k}, \\ 0, & \text{otherwise.} \end{cases}$$

# Results (L., 2022)

#### Then

$$1) \sum_{n_1,\dots,n_k \ge 0} (q)_{n_k} \frac{q^{n_1^2 + \dots + n_{k-1}^2 + n_1 + \dots + n_{k-1}}}{(-q)_{n_1}} \prod_{i=1}^{k-1} {n_{i+1} \brack n_i}$$
  
" = " - (1 +  $\delta_{k,1}$ )  $\sum_{n \ge 0} n\chi_{4k}(n) q^{\frac{n^2 - (k-1)^2}{4k}}$ 

and

$$2) \sum_{n_1,\dots,n_k \ge 0} (q^2; q^2)_{n_k} \frac{q^{2n_1^2 + 2n_1 + \dots + 2n_{k-1}^2 + 2n_{k-1}}(q; q^2)_{n_1}}{(-q)_{2n_1+1}} \prod_{i=1}^{k-1} {n_{i+1} \choose n_i}_{q^2}$$
  
$$" = " - (1 + \delta_{k,1}) \frac{1}{2} \sum_{n \ge 0} n\chi_{8k-4}(n) q^{\frac{n^2 - (2k-2)^2}{8k-4}}.$$

# Results (L., 2022)

Let  $\chi^{(a)}_{8k}(n)$  be the even periodic function modulo 8k defined by

$$\chi_{8k}^{(a)}(n) = \begin{cases} 1, & \text{if } n \equiv 2k - 2a - 1 \text{ or } 6k + 2a + 1 \pmod{8k}, \\ -1, & \text{if } n \equiv 2k + 2a + 1 \text{ or } 6k - 2a - 1 \pmod{8k}, \\ 0, & \text{otherwise.} \end{cases}$$

Then

3) 
$$\sum_{n_1,\dots,n_k \ge 0} (q^2; q^2)_{n_k} \frac{q^{2n_1^2 + \dots + 2n_{k-1}^2 + 2n_{a+1} + \dots + 2n_{k-1}}}{(-q; q^2)_{n_1 + \delta_{a,0}}} \prod_{i=1}^{k-1} \left[ \frac{n_{i+1} + \delta_{i,a}}{n_i} \right]_{q^2}$$
  
$$" = " - \frac{1}{2} \sum_{n \ge 0} n\chi_{8k}^{(a)}(n) q^{\frac{n^2 - (2k - 2a - 1)^2}{8k}}.$$



There are many more.

One has all of the usual applications.

In addition, one obtains interesting q-series identities at roots of unity by comparing strange identities.

# Results (L., 2022)

Comparing the first two strange identities we have that

$$\sum_{n_1,\dots,n_{2k-1}\geq 0} (q)_{n_{2k-1}} \frac{q^{n_1^2+\dots+n_{2k-2}^2+n_1+\dots+n_{2k-2}}}{(-q)_{n_1}} \prod_{i=1}^{2k-2} \left[ \begin{array}{c} n_{i+1} \\ n_i \end{array} \right]$$
$$= 2 \sum_{n_1,\dots,n_k\geq 0} (q^2;q^2)_{n_k} \frac{q^{2n_1^2+2n_1+\dots+2n_{k-1}^2+2n_{k-1}}(q;q^2)_{n_1}}{(-q)_{2n_1+1}} \prod_{i=1}^{k-1} \left[ \begin{array}{c} n_{i+1} \\ n_i \end{array} \right]_{q^2}$$

at any odd root of unity (in which case the sums become finite), but of course not as functions inside the unit disk (where neither series converges, anyway).

The case k = 1 reads

$$\sum_{n\geq 0} \frac{(q)_n}{(-q)_n} = 2\sum_{n\geq 0} \frac{(q)_{2n}}{(-q)_{2n+1}}.$$

## Proof of Zagier's identity

Using *q*-difference equations, Zagier showed that for |x| < 1,

$$(1-x)\sum_{n\geq 0}(xq)_nx^n=\sum_{n\geq 0}(-1)^nx^{3n}q^{n(3n+1)/2}(1-x^2q^{2n+1}).$$

We want to take  $\frac{d}{dx}|_{x=1}$ .

Zagier added and subtracted  $(x)_{\infty}$  on the left-hand side to obtain

$$(1-x)\sum_{n\geq 0} ((xq)_n - (xq)_\infty) x^n + (xq)_\infty$$
$$= \sum_{n\geq 0} (-1)^n x^{3n} q^{n(3n+1)/2} (1-x^2 q^{2n+1}).$$

# Proof of Zagier's identity

Replacing x by  $x^2$  on both sides, multiplying by x, and taking  $\frac{d}{dx}|_{x=1}$  gives the "sum of tails" identity,

$$2\sum_{n\geq 0} ((q)_n - (q)_{\infty}) + (q)_{\infty} \left( -1 + 2\sum_{n\geq 1} \frac{q^n}{1 - q^n} \right)$$
$$= -\sum_{n\geq 1} n \left( \frac{12}{n} \right) q^{(n^2 - 1)/24}.$$

Letting q approach a root of unity, we obtain Zagier's strange identity, since  $(q)_{\infty}$  vanishes to infinite order.

## Proof of Hikami's identity

Using q-difference equations and some impressive calculations, Hikami showed that

$$1-x) \sum_{n_1,\dots,n_k \ge 0} (xq)_{n_k} q^{n_1^2+\dots+n_{k-1}^2+n_{a+1}+\dots+n_{k-1}} \\ \times x^{2n_1+\dots+2n_{k-1}+n_k} \prod_{i=1}^{k-1} \begin{bmatrix} n_{i+1}+\delta_{a,i} \\ n_i \end{bmatrix} \\ = \sum_{n\ge 0} (-1)^n x^{(2k+1)n} q^{\binom{n+1}{2}+(a+1)n^2+(k-a-1)(n^2+n)} \\ \times (1-x^{2(a+1)}q^{(a+1)(2n+1)}) \\ = \sum_{n\ge 0} \chi_{8k+4}^{(a)}(n) q^{\frac{n^2-(2k-2a-1)^2}{8(2k+1)}} x^{\frac{n-(2k-2a-1)}{2}}.$$

# Proof of Hikami's identity

Arguing along the lines of Zagier (with considerably more complications), Hikami arrived at his strange identity.

One new element he required was the "Andrews-Gordon variant"

$$\sum_{\substack{n_1, n_2, \dots, n_{k-1} \ge 0}} \frac{q^{n_1^2 + \dots + n_{k-1}^2 + n_{a+1} + \dots + n_{k-1}}}{(q)_{n_{k-1}}} \prod_{i=1}^{k-2} \begin{bmatrix} n_{i+1} + \delta_{a,i} \\ n_i \end{bmatrix}$$
$$= \frac{1}{(q)_{\infty}} \sum_{\substack{n \ge 0}} \chi_{8m+4}^{(a)}(n) q^{\frac{n^2 - (2k - 2a - 1)^2}{8(2k+1)}}$$
$$= \prod_{\substack{n \ne 0, \pm (a+1) \pmod{2k+1}}} \frac{1}{1 - q^n}.$$

How does this all fit into the framework of Bailey pairs?

Recall from classical work of Bailey that if  $(\alpha_n, \beta_n)$  is a Bailey pair relative to *a*, then

$$\sum_{n\geq 0} (b)_n (c)_n (aq/bc)^n \beta_n = \frac{(aq/b)_\infty (aq/c)_\infty}{(aq)_\infty (aq/bc)_\infty} \sum_{n\geq 0} \frac{(b)_n (c)_n (aq/bc)^n}{(aq/b)_n (aq/c)_n} \alpha_n.$$

Using  $(a, b, c) = (x^2q, xq, q)$  we have the key identity

$$(1-x)\sum_{n\geq 0}(xq)_n(q)_nx^n\beta_n=(1-x^2q)\sum_{n\geq 0}\frac{(q)_n}{(x^2q)_n}x^n\alpha_n.$$

Still using work of Bailey and Slater, we have the Bailey pair relative to  $x^2q$ ,

$$\alpha_n = \frac{(x^2q)_n(1-x^2q^{2n+1})(-1)^n x^{2n}q^{n(3n+1)/2}}{(q)_n(1-x^2q)}$$

and

$$\beta_n=\frac{1}{(q)_n}.$$

This gives

$$(1-x)\sum_{n\geq 0}(xq)_nx^n=\sum_{n\geq 0}(-1)^nx^{3n}q^{n(3n+1)/2}(1-x^2q^{2n+1}),$$

the starting point for Zagier's proof.

To get to the starting point for Hikami's strange identity requires much more work.

The result we need is that for  $k \ge 1$  and  $0 \le a \le k - 1$ , the following is a Bailey pair relative to  $x^2q$ :

$$\alpha_n = \frac{(x^2 q)_n}{(q)_n (1 - x^2 q)} (-1)^n x^{2kn} q^{\binom{n+1}{2} + (a+1)n^2 + (k-a-1)(n^2 + n)} \times (1 - x^{2(a+1)} q^{(a+1)(2n+1)})$$

and

$$\beta_{n} = \beta_{n_{k}} = \sum_{\substack{n_{1}, n_{2}, \dots, n_{k-1} \ge 0}} \frac{q^{n_{1}^{2} + \dots + n_{k-1}^{2} + n_{a+1} + \dots + n_{k-1}} x^{2n_{1} + \dots + 2n_{k-1}}}{(q)_{n_{k}}} \times \prod_{i=1}^{k-1} \begin{bmatrix} n_{i+1} + \delta_{a,i} \\ n_{i} \end{bmatrix}.$$

Using this in the key identity one obtains the starting point for Hikami's proof.

It also gives Hikami's Andrews-Gordon variant using a different application of the Bailey lemma.

Its proof requires a sequence of new Bailey-type lemmas, culminating in a key lemma, which says that if  $(\alpha_n, \beta_n)$  is a Bailey pair relative to *a*, then  $(\alpha''_n, \beta''_n)$  is a Bailey pair relative to *aq*, where

$$\alpha_n'' = \frac{1}{1 - aq} \left( \frac{1 - q^{n+1}}{1 - aq^{2n+2}} \alpha_{n+1} + \frac{q^n (1 - aq^n)}{1 - aq^{2n}} \alpha_n \right)$$

and

$$\beta_n'' = (1 - q^{n+1})\beta_{n+1}.$$

The proof of the desired Bailey pair then goes as follows:

Start with the unit Bailey pair relative to  $x^2$ , iterate a + 1 times along the Bailey chain with  $b, c \to \infty$ , apply the key lemma to obtain a Bailey pair relative to  $x^2q$ , iterate k - 1 - a times along the Bailey chain with  $b, c \to \infty$ , and simplify to obtain the result.

# Conclusion

Many other Bailey pairs can be used in the key identity to (ultimately) obtain strange identities.

The presence of the term  $(q)_n/(x^2q)_n$  does restrict the possibilities.

Behind the scenes, we've always iterated along the Bailey chain using  $b,c \to \infty$ . Other choices give, for example,

$$\sum_{n_1,\dots,n_k\geq 0} \frac{(q)_{n_k}(-q)_{n_{k-1}}q^{n_1^2+\dots+n_{k-2}^2+\binom{n_{k-1}+1}{2}+n_{a+1}+\dots+n_{k-2}}}{(-q)_{n_k}}$$
$$\times \prod_{i=1}^{k-1} \begin{bmatrix} n_{i+1}+\delta_{i,a} \\ n_i \end{bmatrix}$$
$$"="-\sum_{n\geq 0} n\chi_{4k}^{(a)}(n)q^{\frac{n^2-(k-a-1)^2}{4k}}.$$

# Conclusion

Here  $0 \le a < k - 1$ , and  $\chi_{4k}^{(a)}(n)$  is the even periodic function modulo 4k defined by

$$\chi_{4k}^{(a)}(n) = \begin{cases} 1, & \text{if } n \equiv k - a - 1 \text{ or } 3k + a + 1 \pmod{4k}, \\ -1, & \text{if } n \equiv k + a + 1 \text{ or } 3k - a - 1 \pmod{4k}, \\ 0, & \text{otherwise.} \end{cases}$$

It does not appear possible to obtain a strange identity for every periodic function.

# Conclusion

One important class we seem to miss is

$$\sum_{n\geq 0}n\chi_{2st}(n)q^{\frac{n^2-(st-s-t)^2}{4st}},$$

where  $\chi_{2st}(n)$  is the even periodic function modulo 2st defined by

$$\chi_{2st}(n) = \begin{cases} 1, & \text{if } n \equiv st - s - t \text{ or } st + s + t \pmod{2st}, \\ -1, & \text{if } n \equiv st - s + t \text{ or } st + s - t \pmod{2st}, \\ 0, & \text{otherwise.} \end{cases}$$

These are related to the torus knots  $T_{s,t}$ .



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